

Elastic knots

(elastic beam under finite rotation and self-contact)

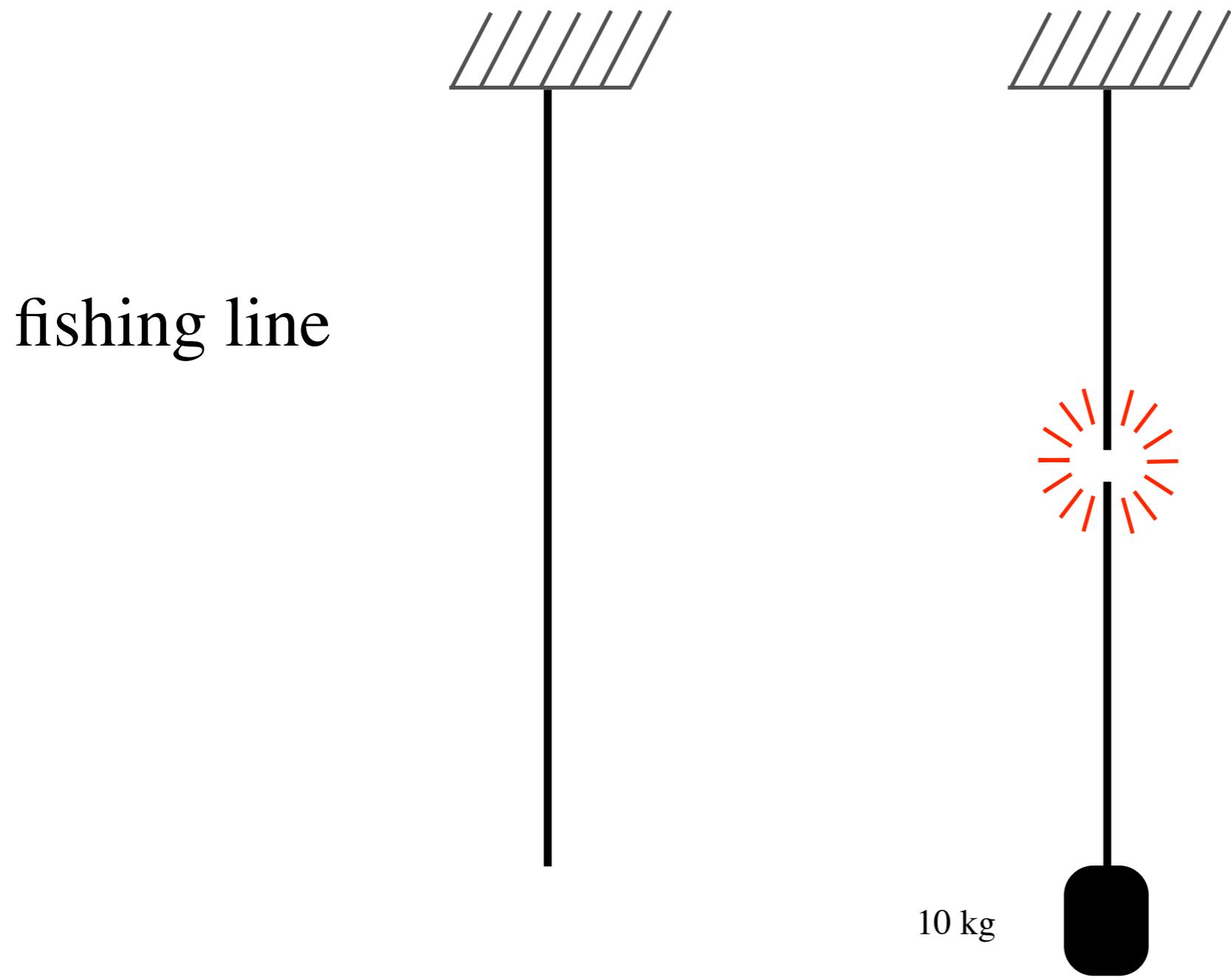
Sébastien Neukirch

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d'Alembert Institute for Mechanics

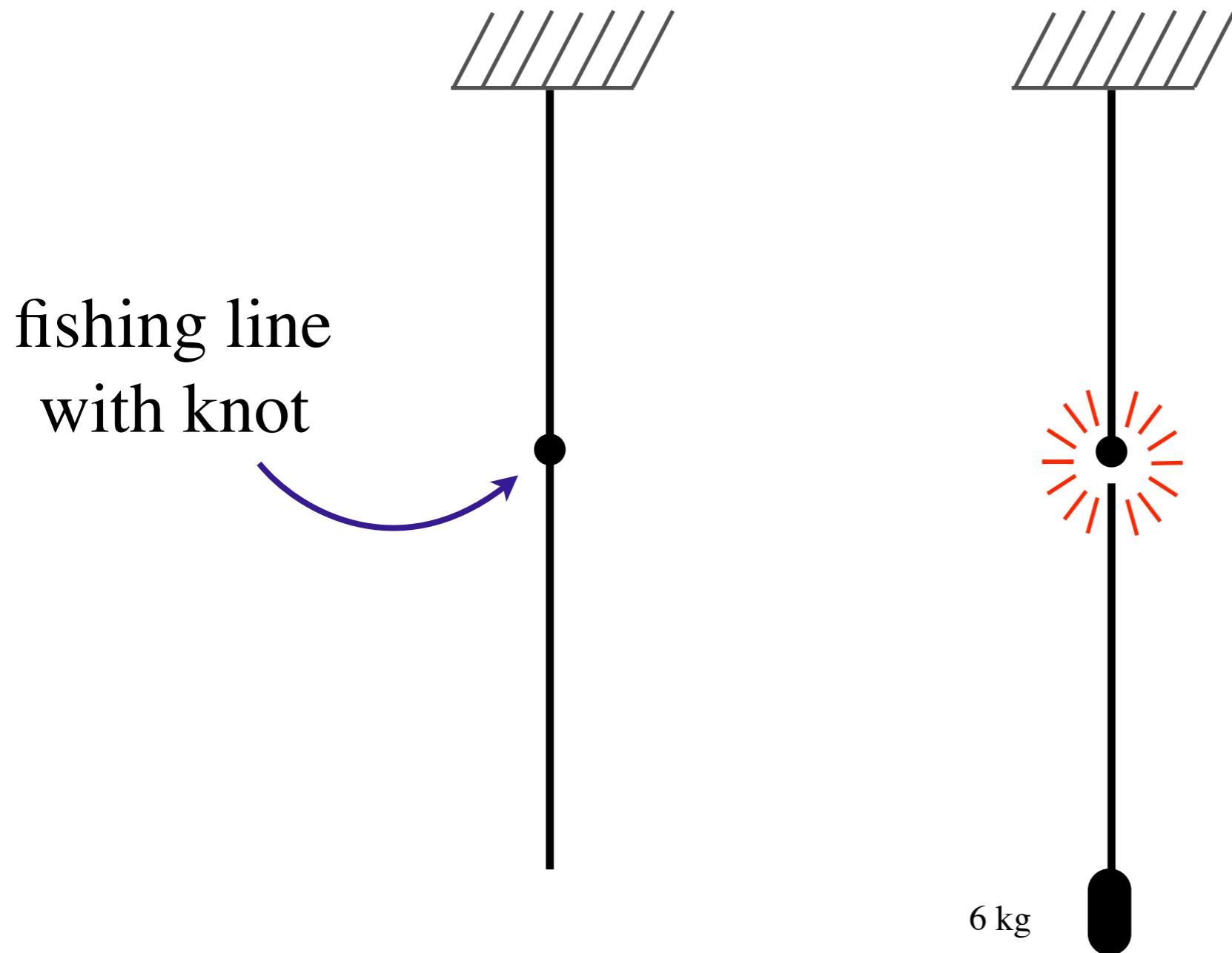
joint work with:

Nicolas Clauvelin (PhD work)
Basile Audoly

Tensile strength of a wire



Tensile strength of a wire



Stasiak et al, Science (1999)

Knots are everywhere

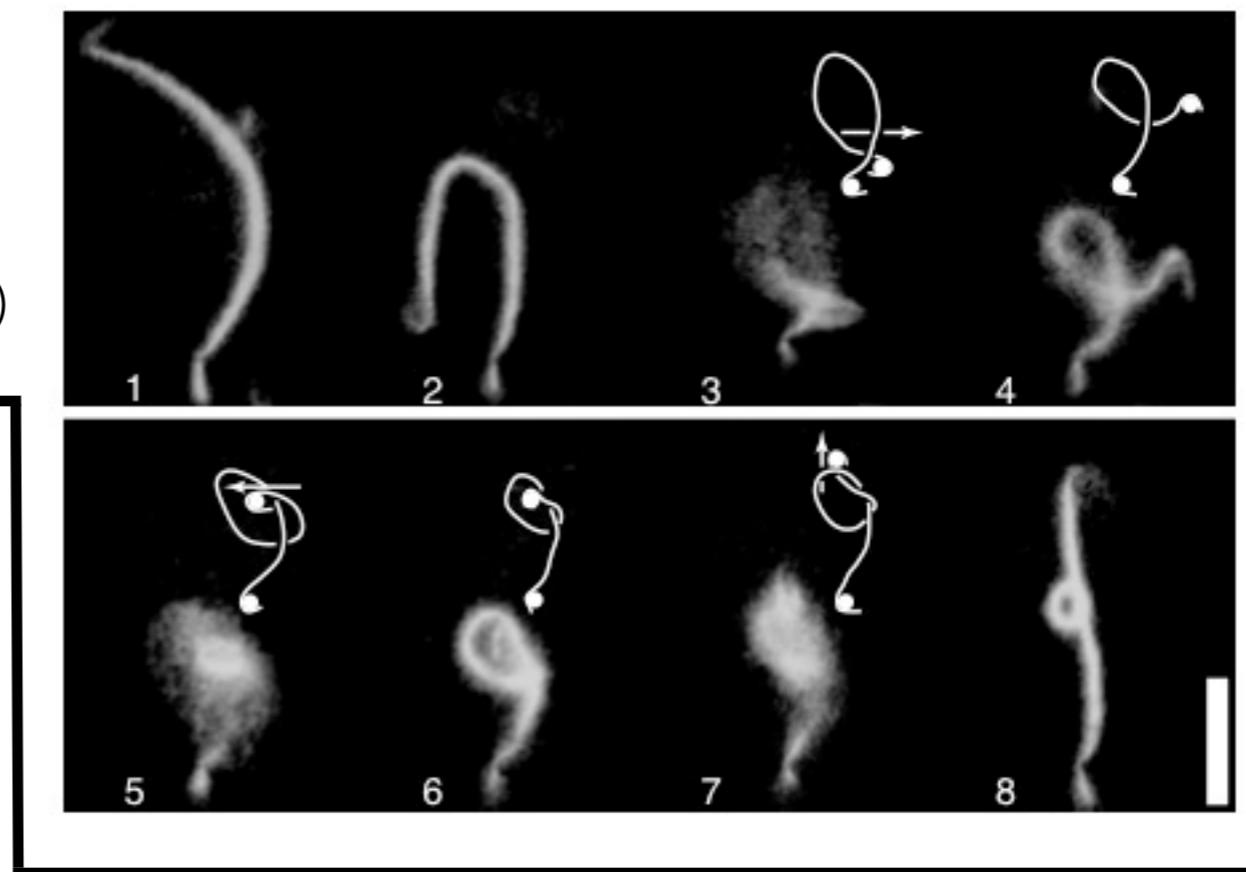
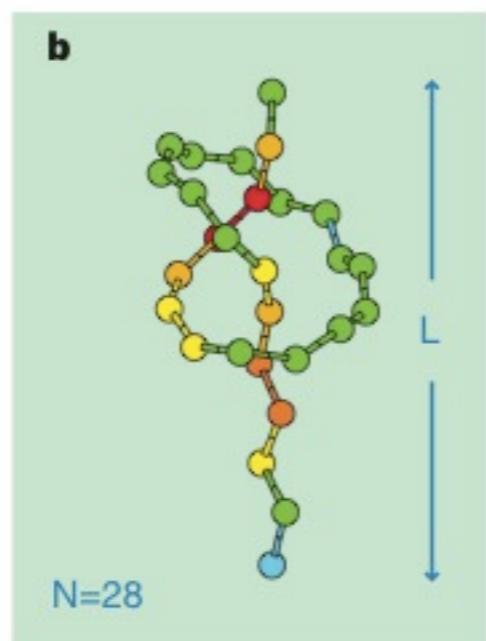
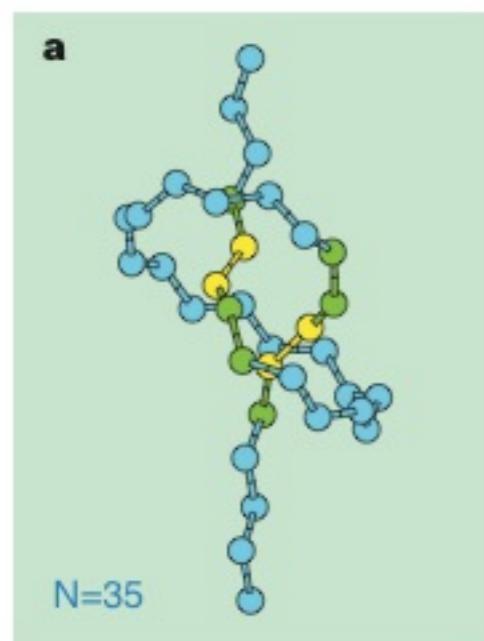
Long enough polymers are (almost) certainly knotted

Sumners+Whittington, J. Phys. A : Math. Gen. 1988

273 knotted proteins in the ProteinDataBank (1%)

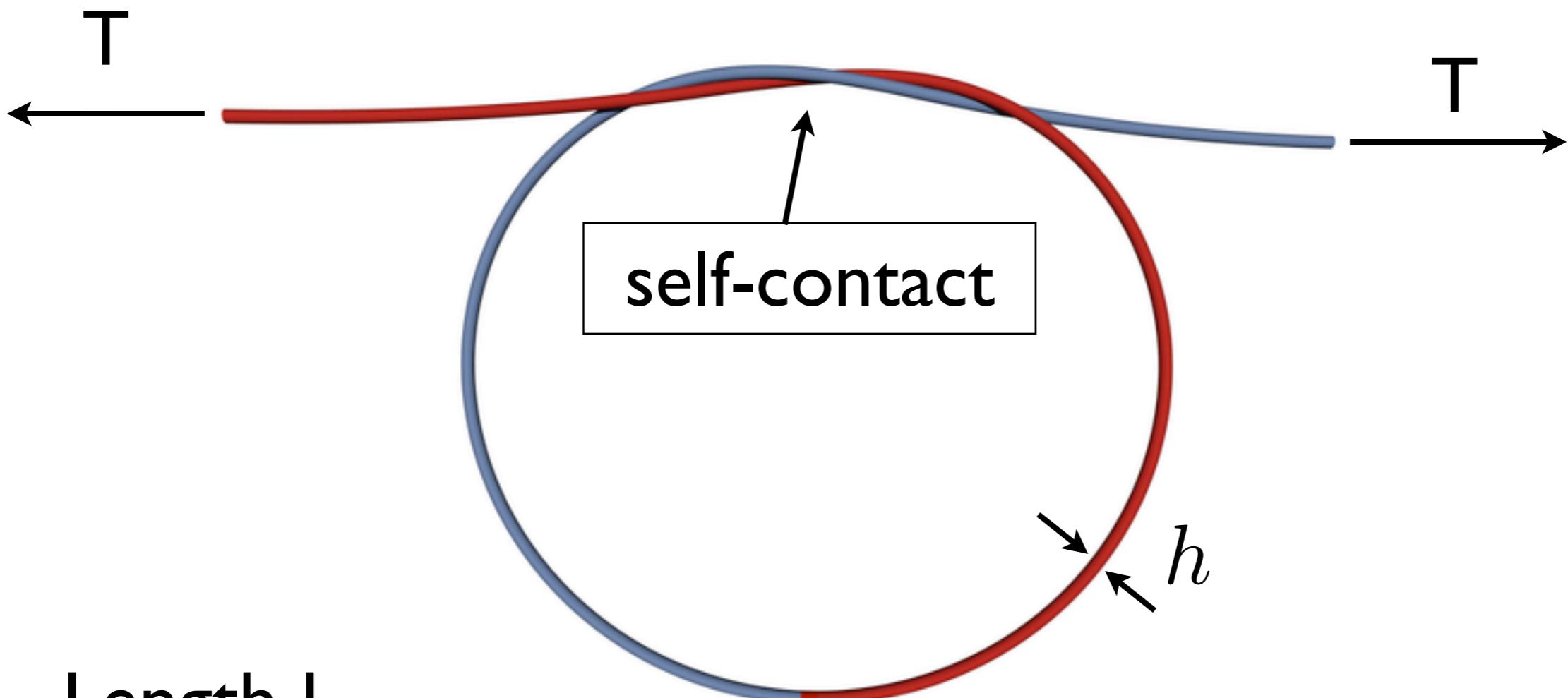
Single molecule experiment
with knotted F-Actin filaments

Arai et al, *Nature* (1999)



Ab-initio molecular simulations
for alcane molecule ($C_{10}H_{22}$)
Saitta et al, *Nature* (1999)

Elastic knots



- Length L
- Circular cross-section: radius h
- Bending rigidity : $E I$
- Twist rigidity : $G J$

E : Young's modulus

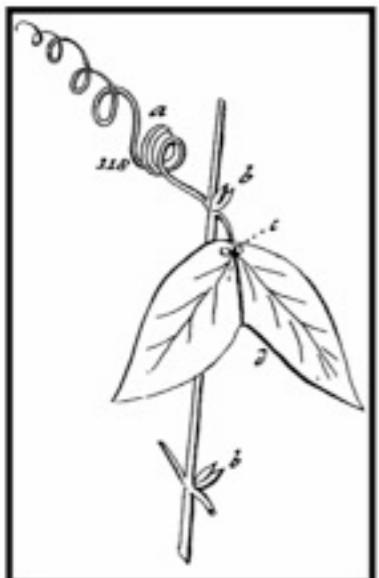
G : shear modulus

$$I = \frac{\pi h^4}{4}$$

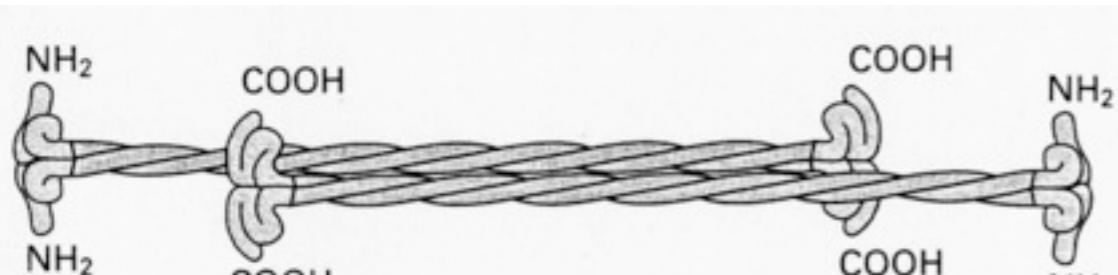
$$J = \frac{\pi h^4}{2}$$

Elastic filaments

climbing plants

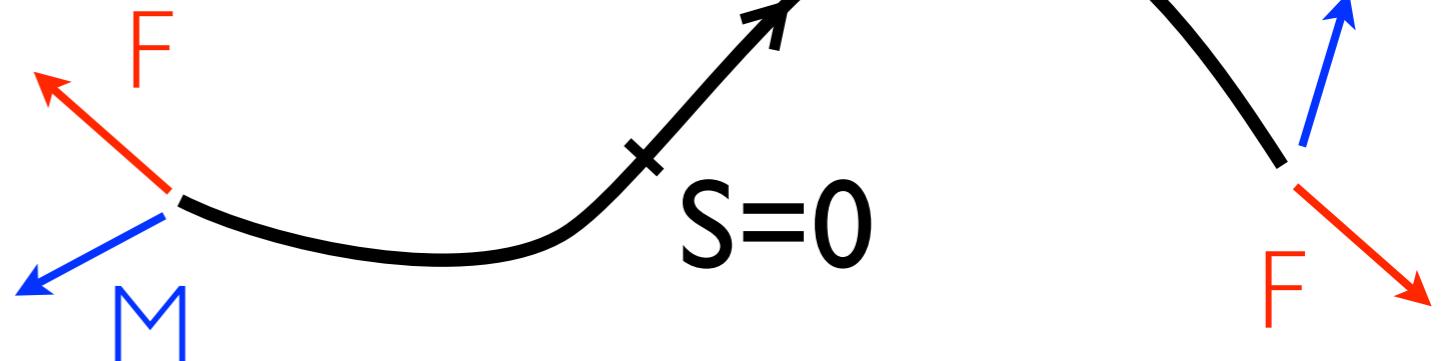


fibrous proteins



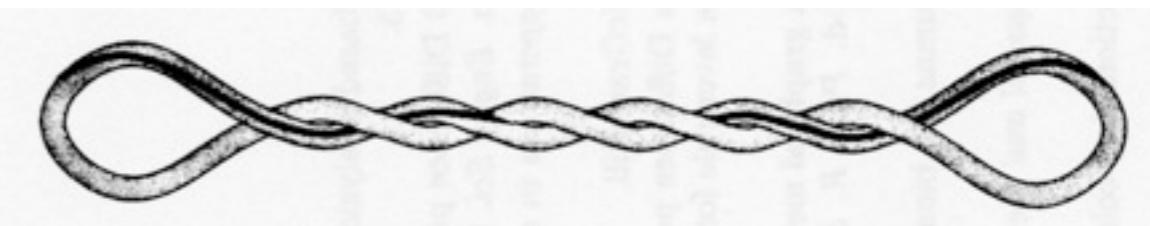
tétramère fait de deux dimères superenroulés étagés

Theory



Applications

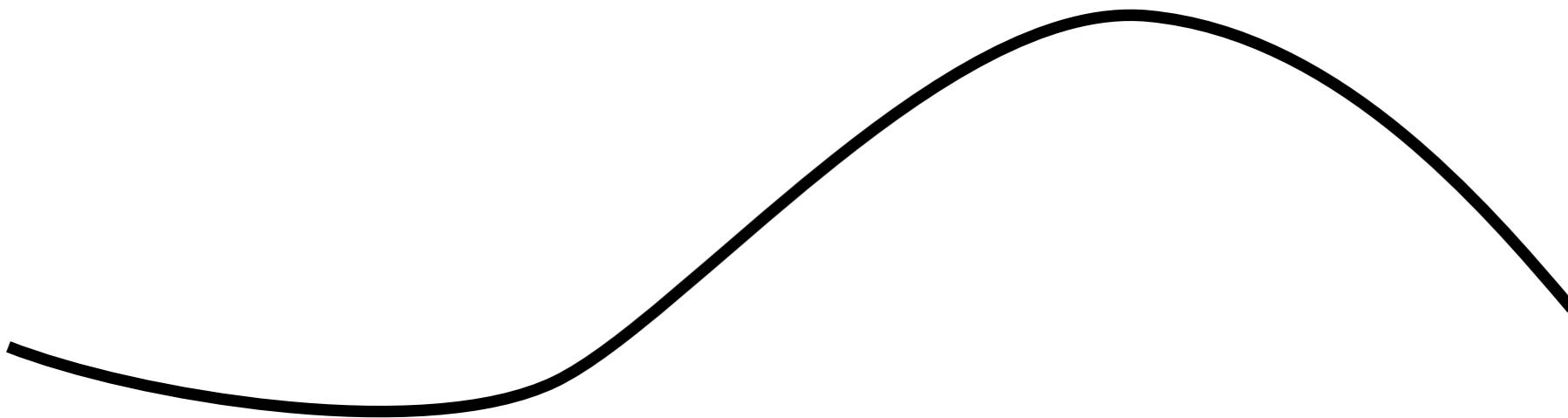
DNA supercoiling



cables



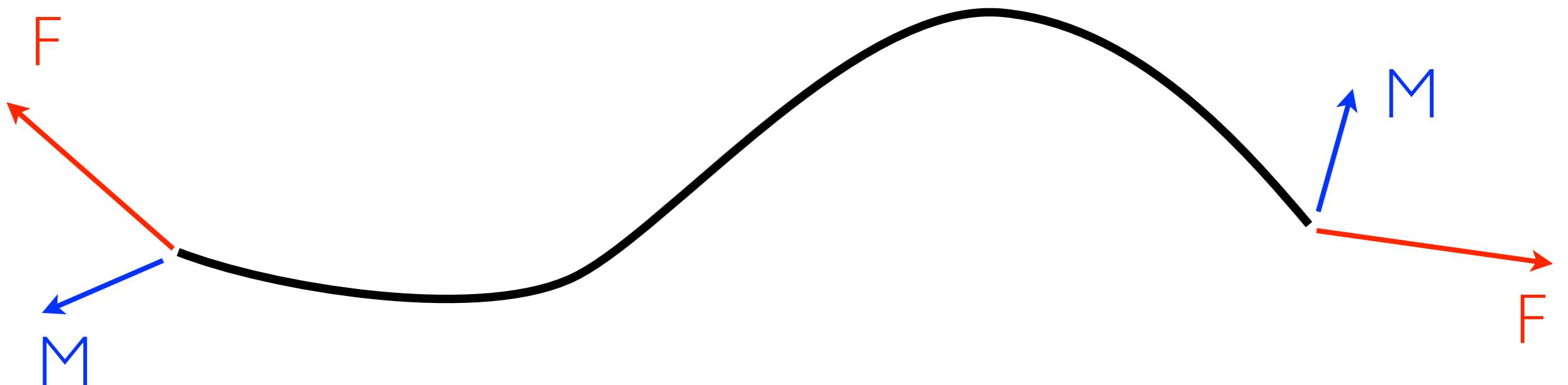
Kirchhoff equations



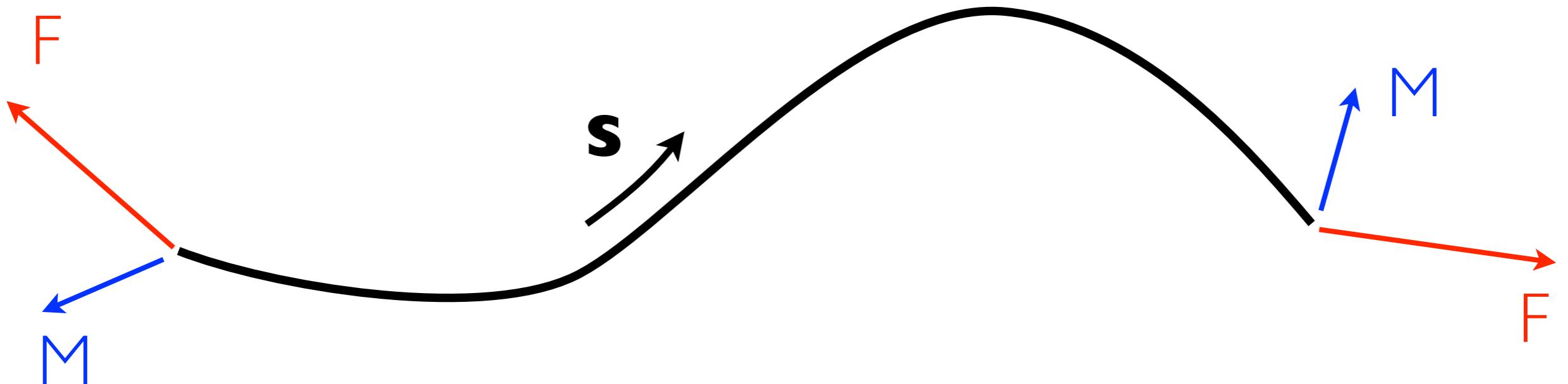
apply to :

- slender bodies
- not too bent

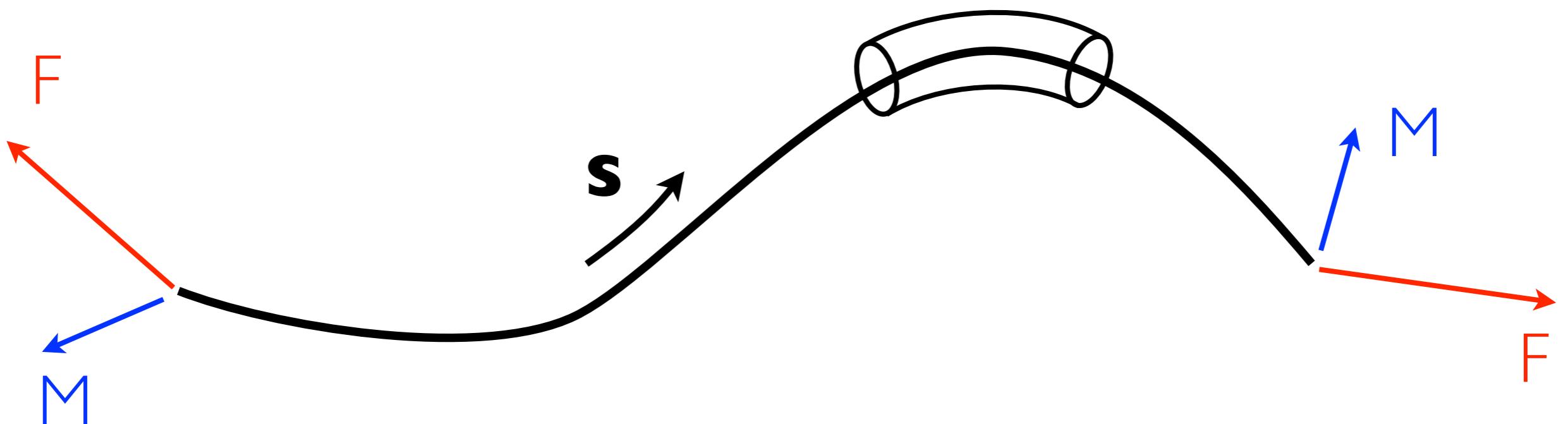
Kirchhoff equations



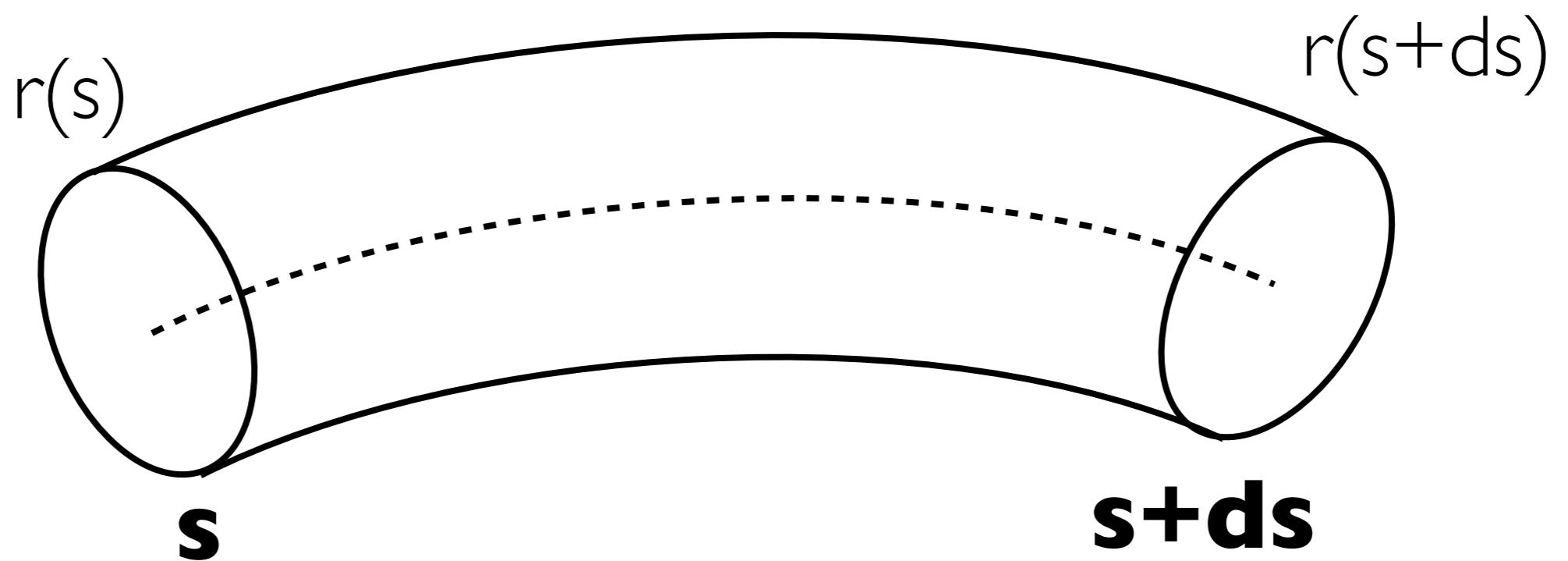
Kirchhoff equations



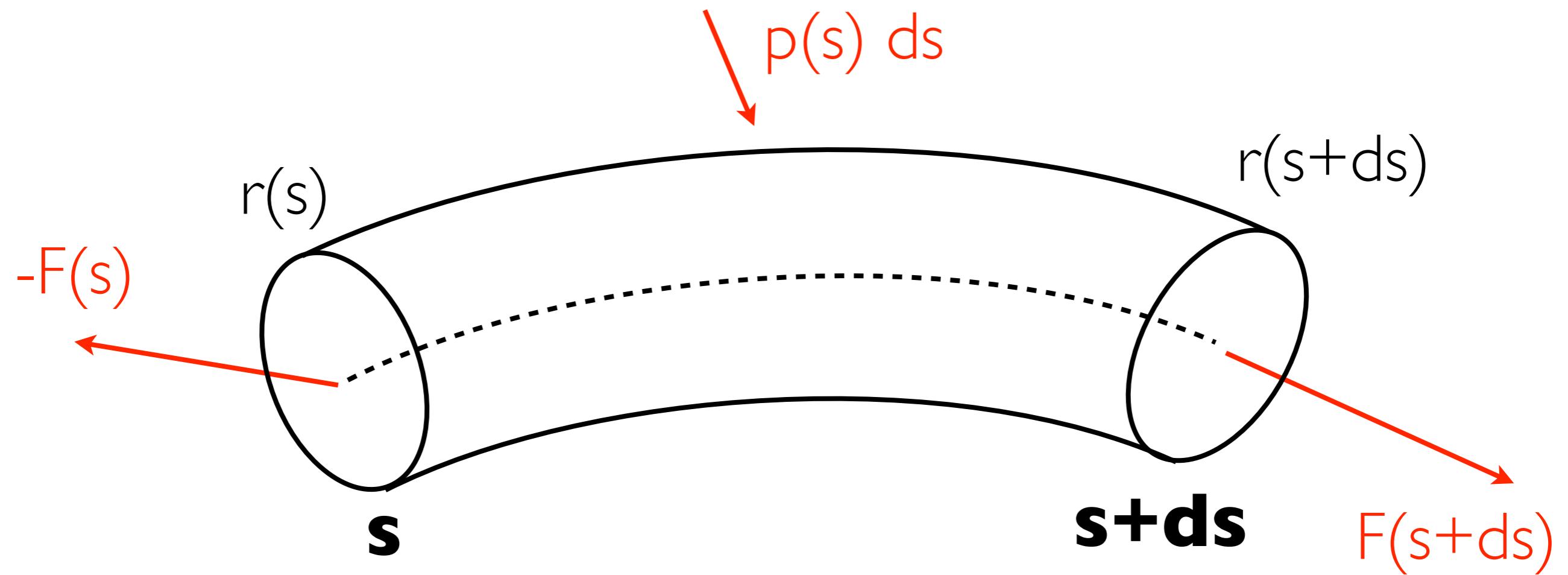
Kirchhoff equations



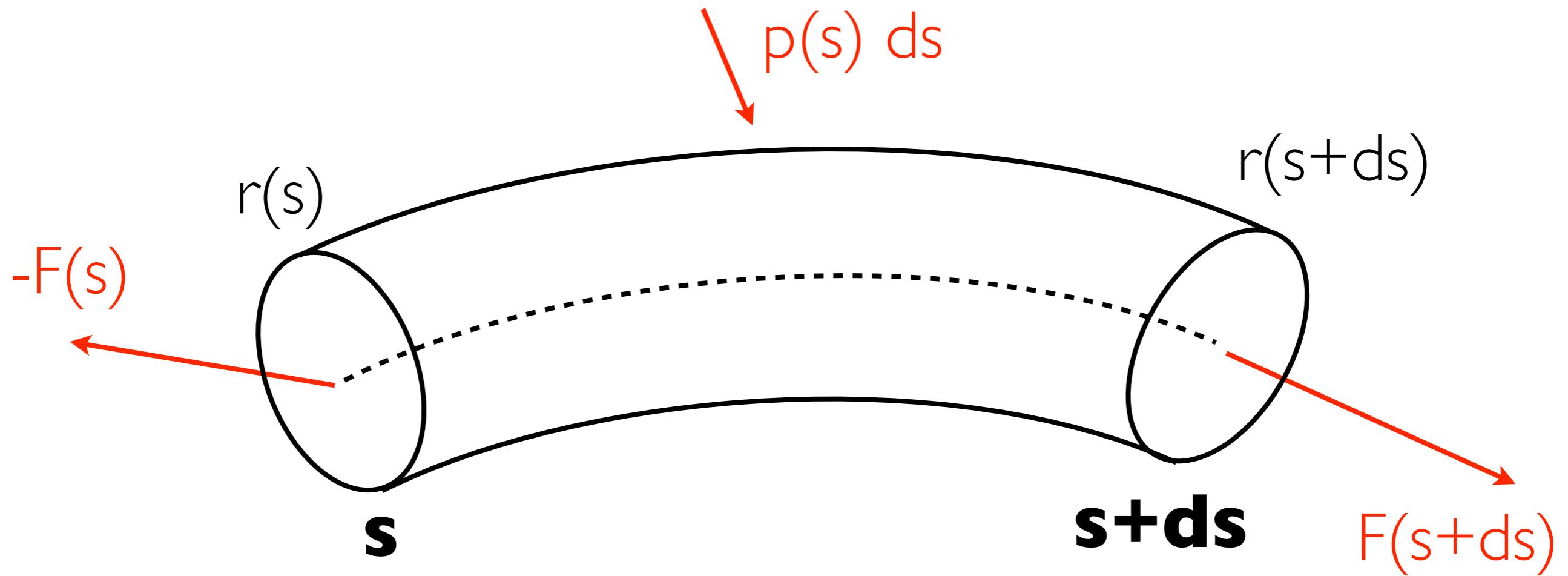
Kirchhoff equations



Kirchhoff equations



Kirchhoff equations

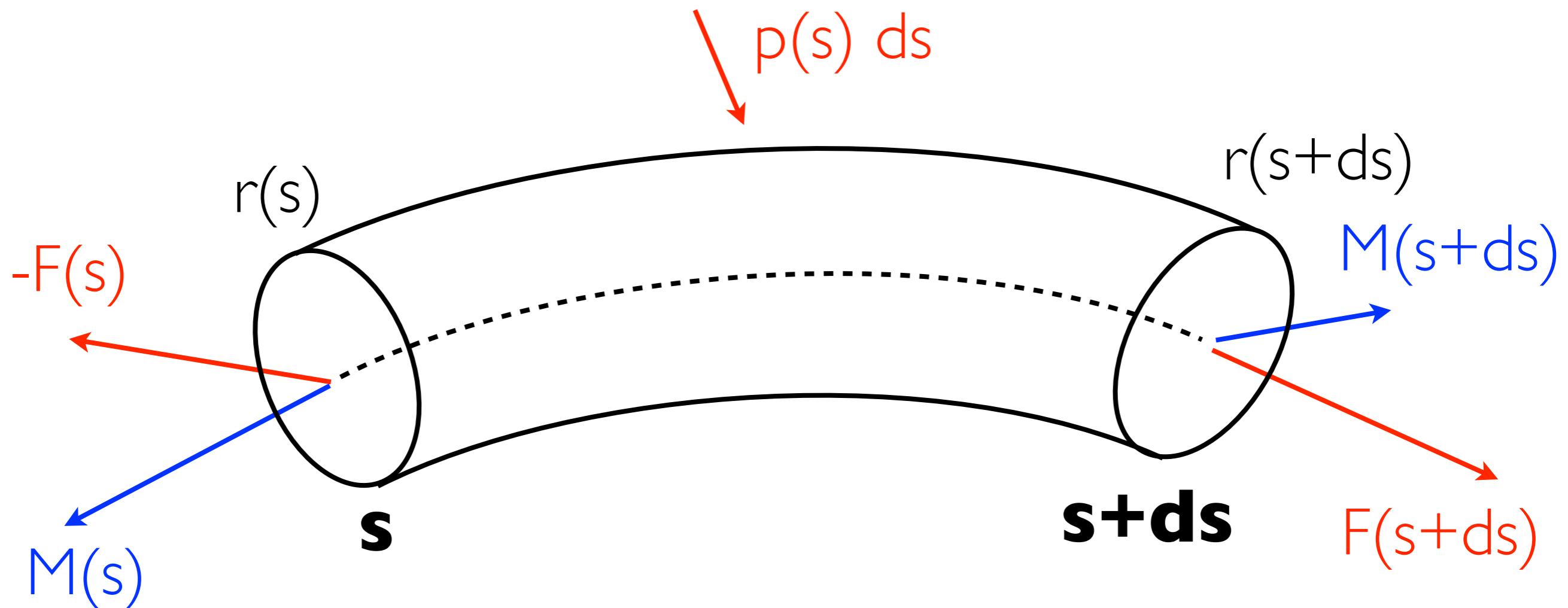


$$F(s+ds) - F(s) + p(s) ds = 0$$

Equilibrium

$$F'(s) + p(s) = 0$$

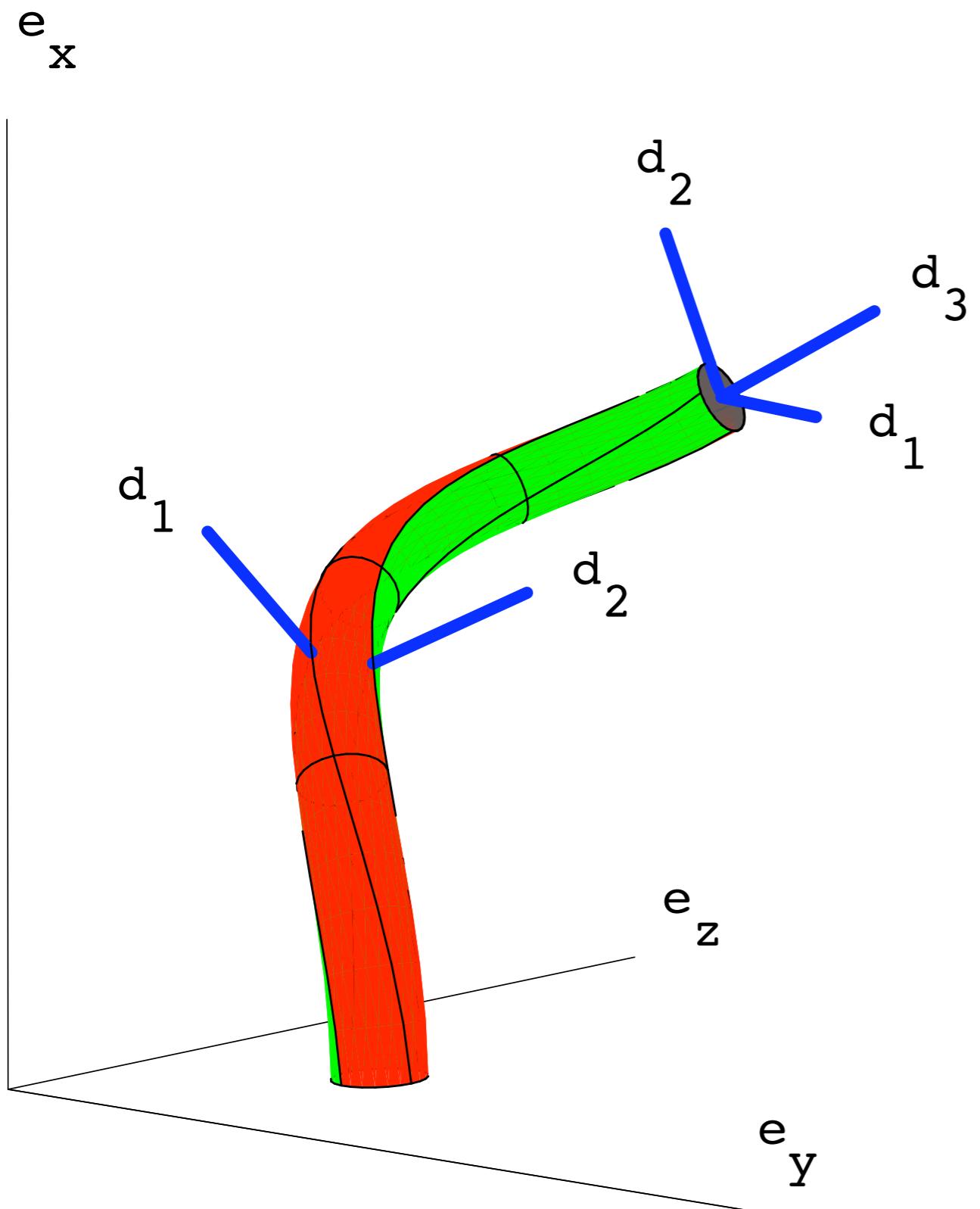
Kirchhoff equations



Equilibrium

$$M' + r' \times F = 0$$

Kirchhoff equations



Cosserat frame

$$d'_1 = u \times d_1$$

$$d'_2 = u \times d_2$$

$$d'_3 = u \times d_3$$

$$u = \{\kappa_1, \kappa_2, \tau\}_{d_i}$$

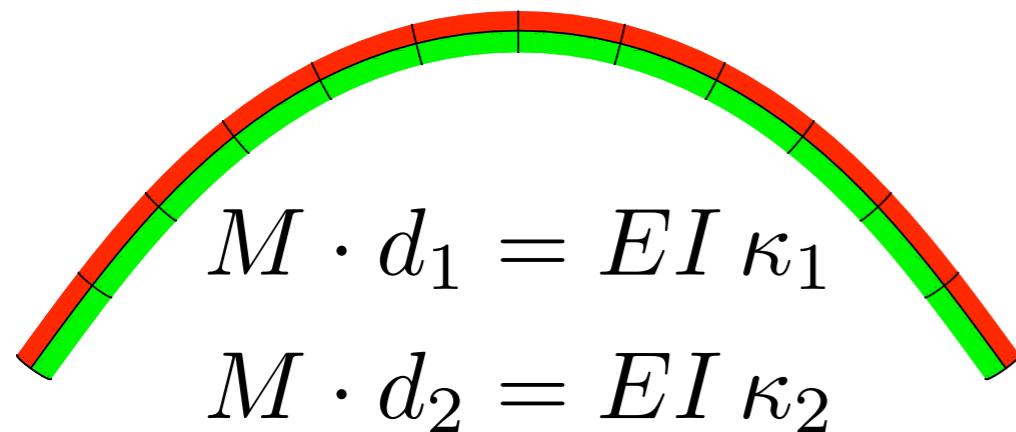
curvatures

twist

Kirchhoff equations

constitutive relations

curvature



E Young's modulus
 I second moment of area

twist



$$M \cdot d_3 = GJ \tau$$

G shear modulus
 J polar moment of area

Kirchhoff equations

21 ODEs with variable : s

ordinary differential equations

$$\frac{d}{ds} \vec{F} = \vec{p}$$

$$\frac{d}{ds} \vec{M} = \vec{F} \wedge \vec{d}_3$$

$$\frac{d}{ds} \vec{r} = \vec{d}_3$$

$$\frac{d}{ds} \vec{d}_i = \vec{u} \wedge \vec{d}_i$$

$$m_i = K_i u_i$$

21 unknowns

$$\vec{F}(s)$$

$$\vec{M}(s)$$

$$\vec{r}(s)$$

$$\vec{d}_3(s)$$

$$\vec{d}_2(s)$$

$$\vec{d}_1(s)$$

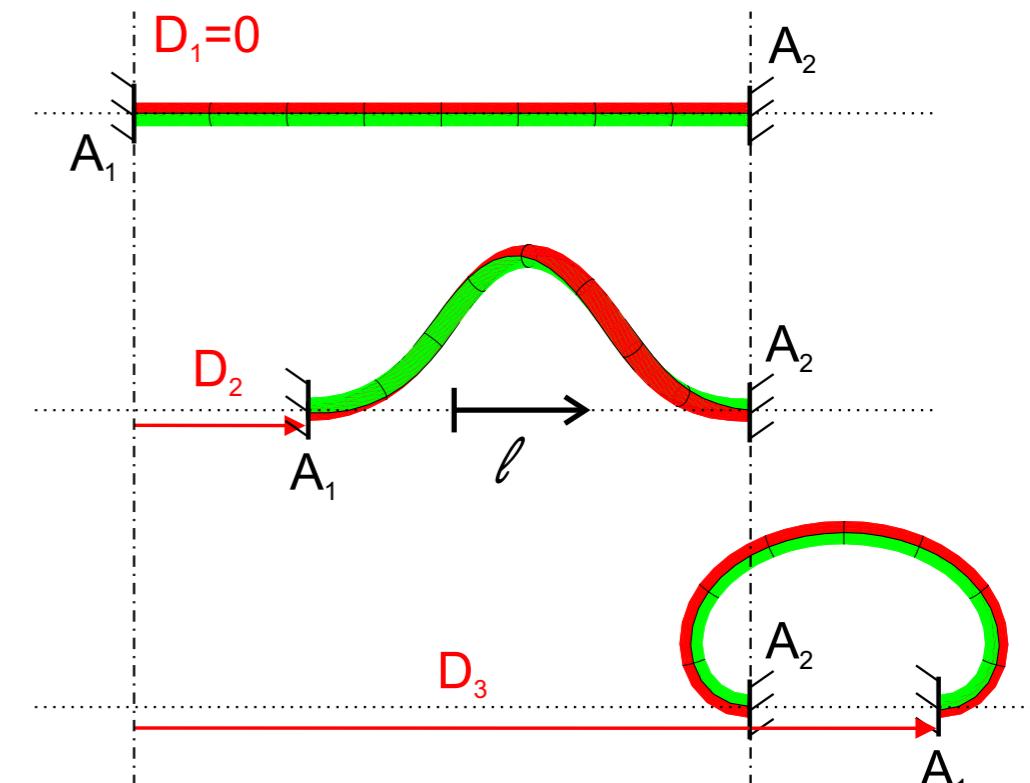
$$\vec{u}(s)$$

$$i=1,2,3$$

linear elasticity

boundary conditions

- how the rod is held
- few solutions are admissibles

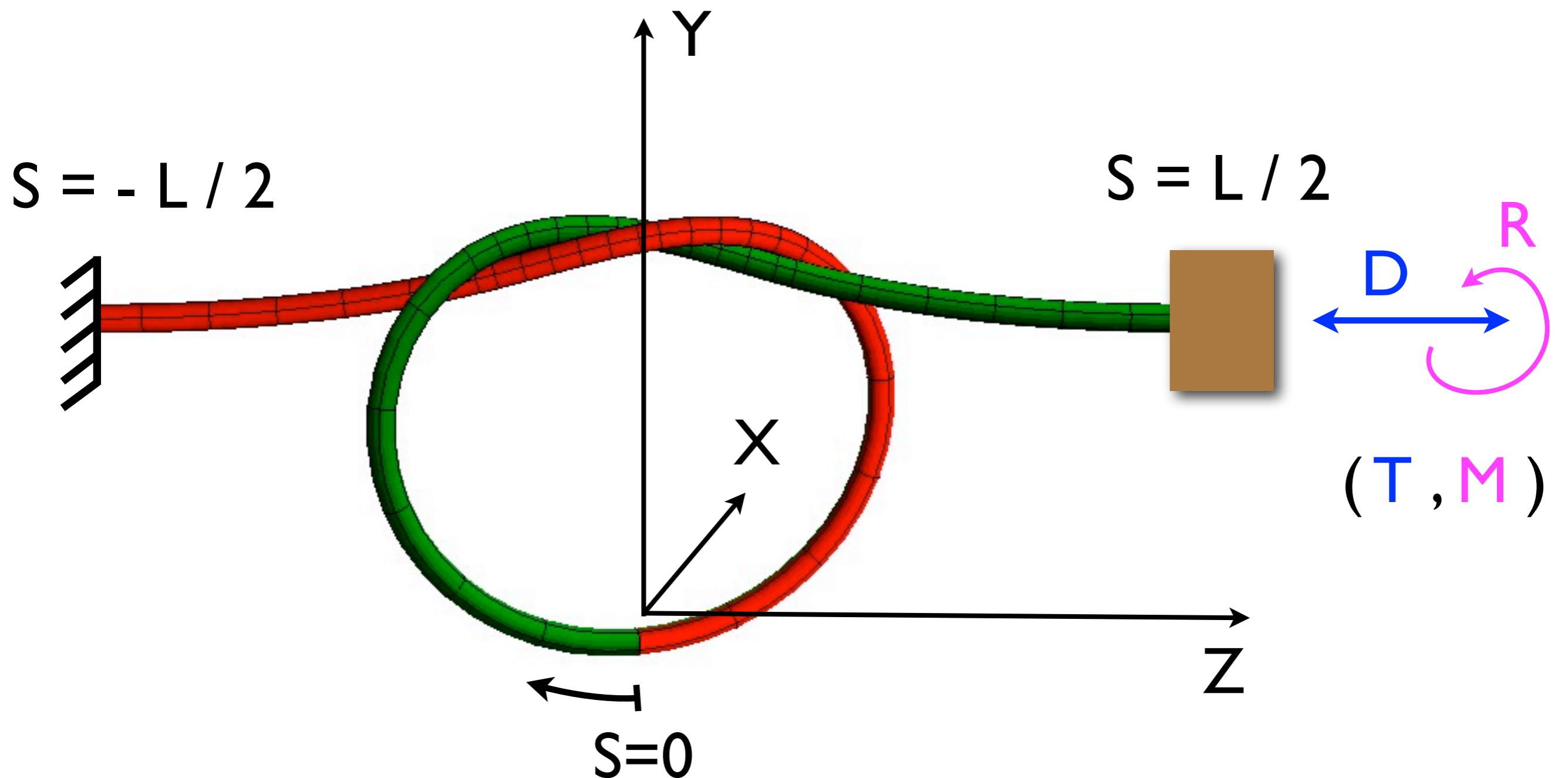


$$\vec{d}_3(A_1) = \vec{d}_3(A_2)$$

$$\vec{r}(A_2) - \vec{r}(A_1) = k \vec{d}_3(A_2)$$

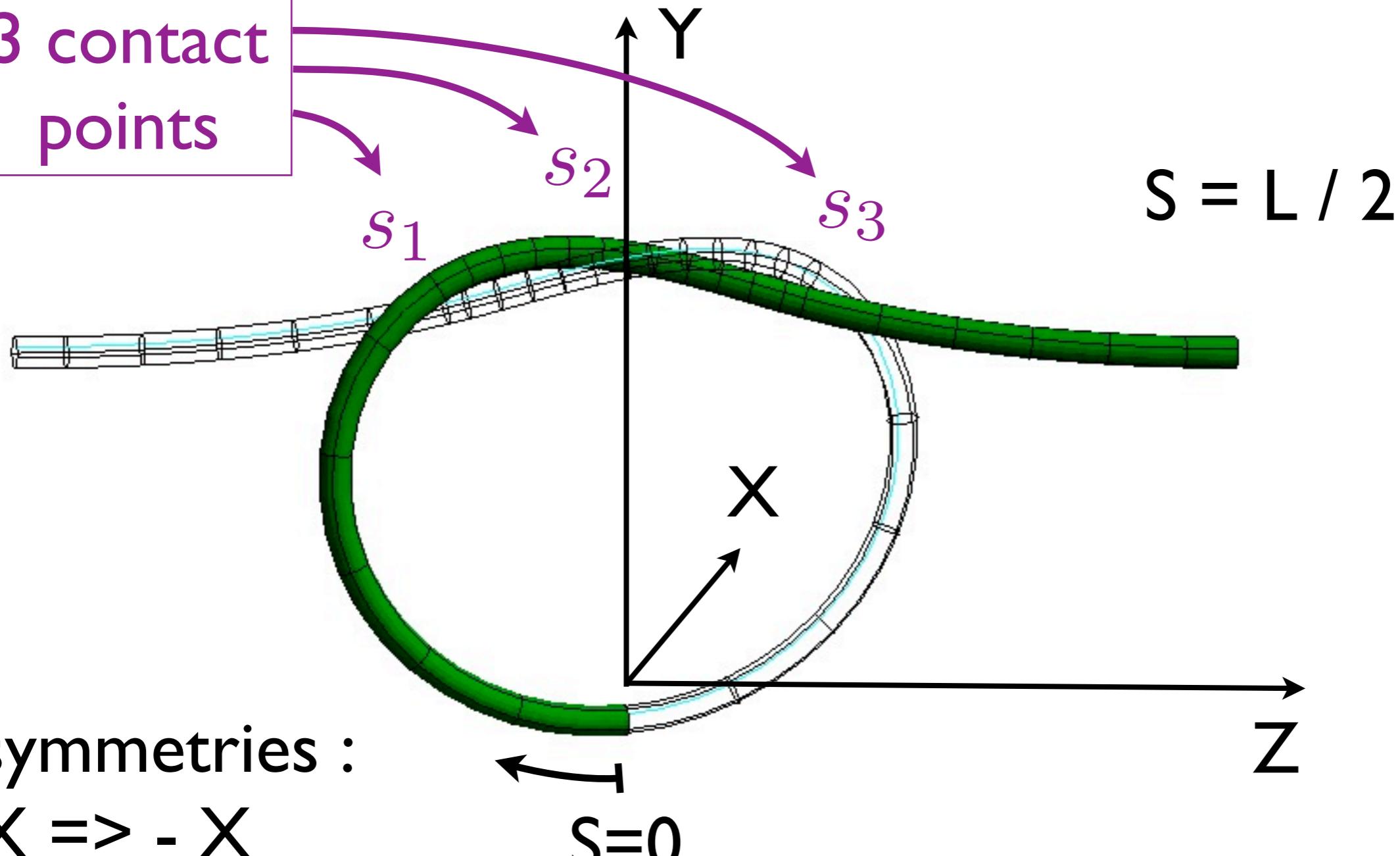
$$(D=L-k)$$

Boundary value problem



Boundary value problem

3 contact points



symmetries :

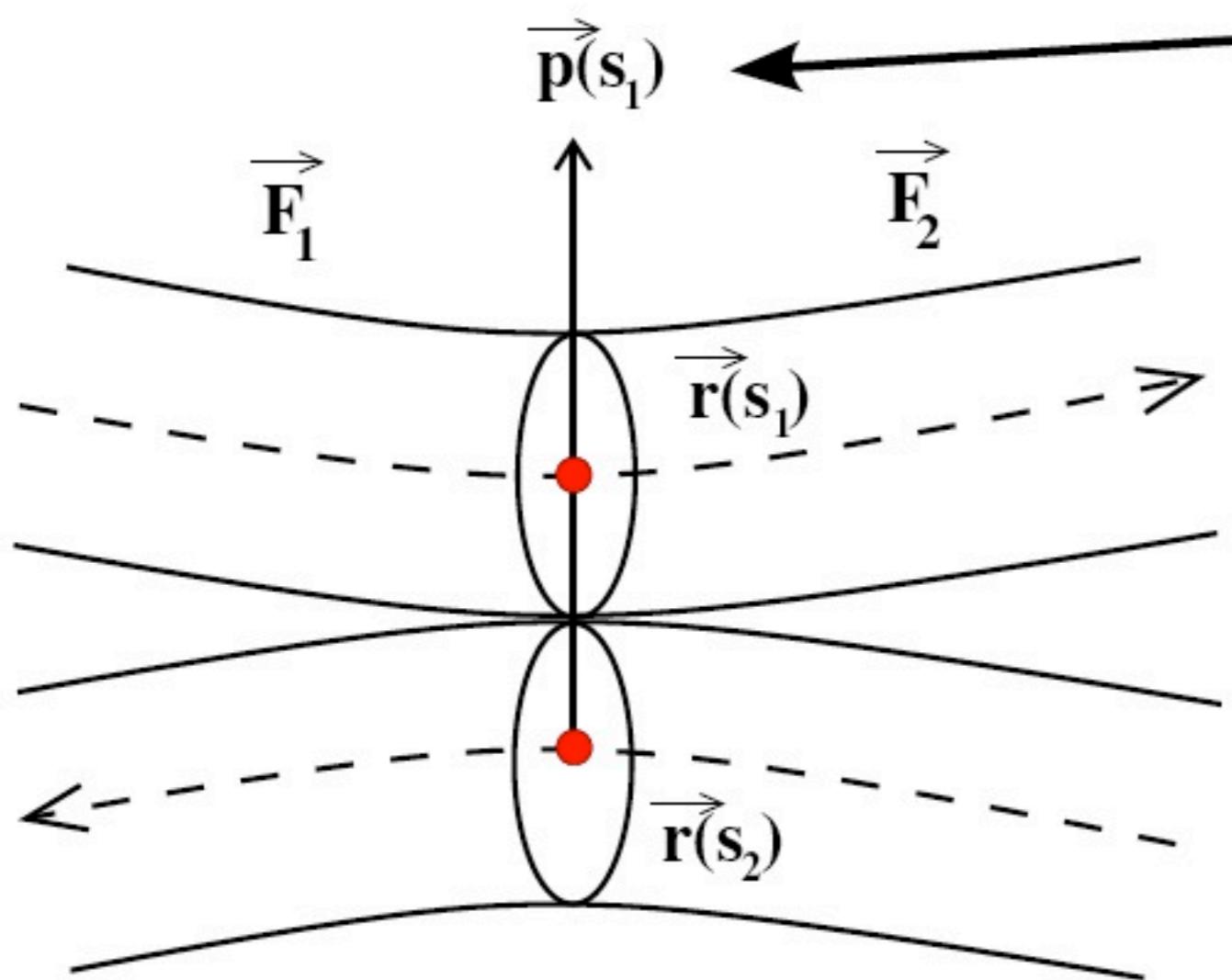
$X \Rightarrow -X$

$Y \Rightarrow Y$

$Z \Rightarrow -Z$

- Shooting method (Mathematica)
- Gauss colocation (AUTO)

Hard-wall contact, no friction



force from strand at s_2
acting on strand at s_1

$$\vec{F}_1 = \vec{p} + \vec{F}_2$$

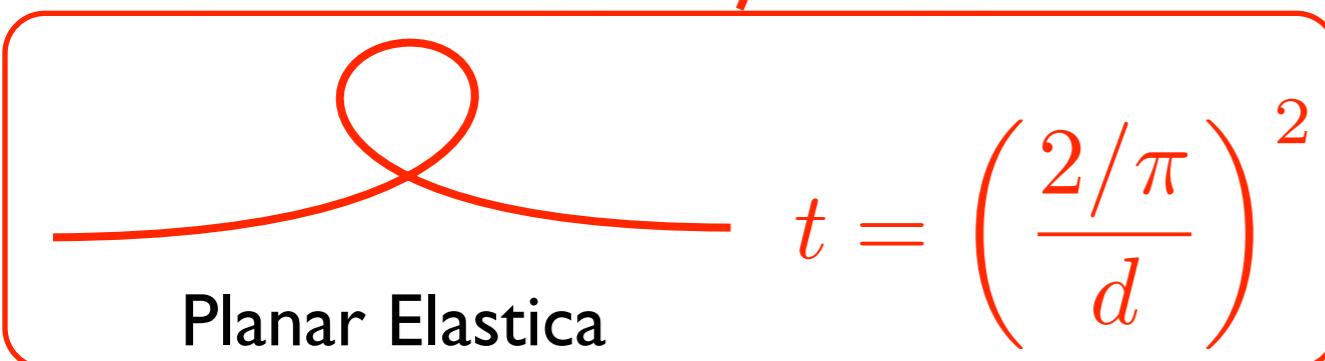
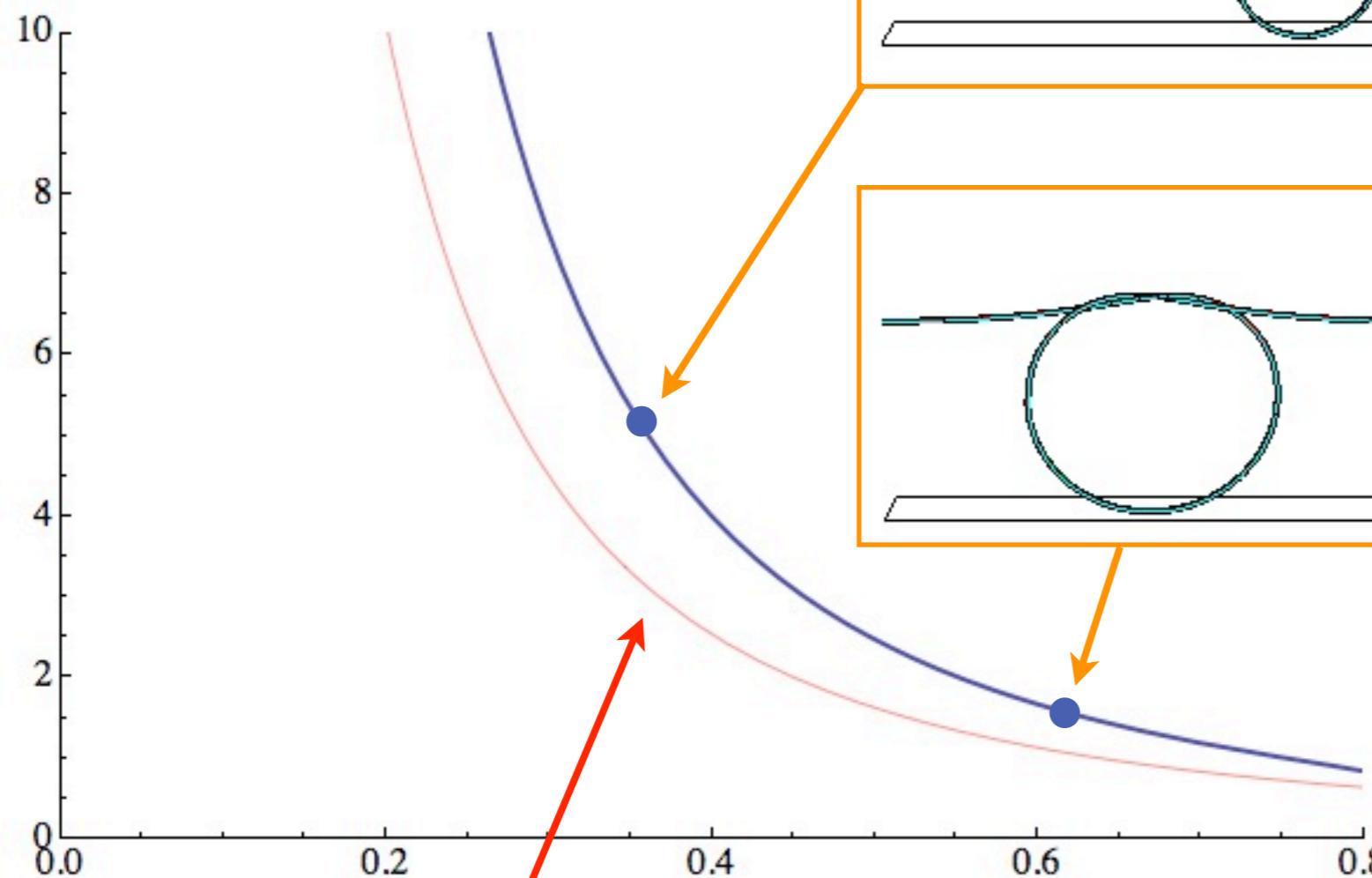
$$\vec{p} = p \frac{\vec{r}(s_1) - \vec{r}(s_2)}{|\vec{r}(s_1) - \vec{r}(s_2)|}$$

touching conditions :

$$\left\{ \begin{array}{l} |\vec{r}(s_1) - \vec{r}(s_2)| = \text{thickness} \\ (\vec{r}(s_1) - \vec{r}(s_2)) \perp \vec{d}_3(s_1) \\ (\vec{r}(s_1) - \vec{r}(s_2)) \perp \vec{d}_3(s_2) \end{array} \right.$$

Numerical Path Following : Results

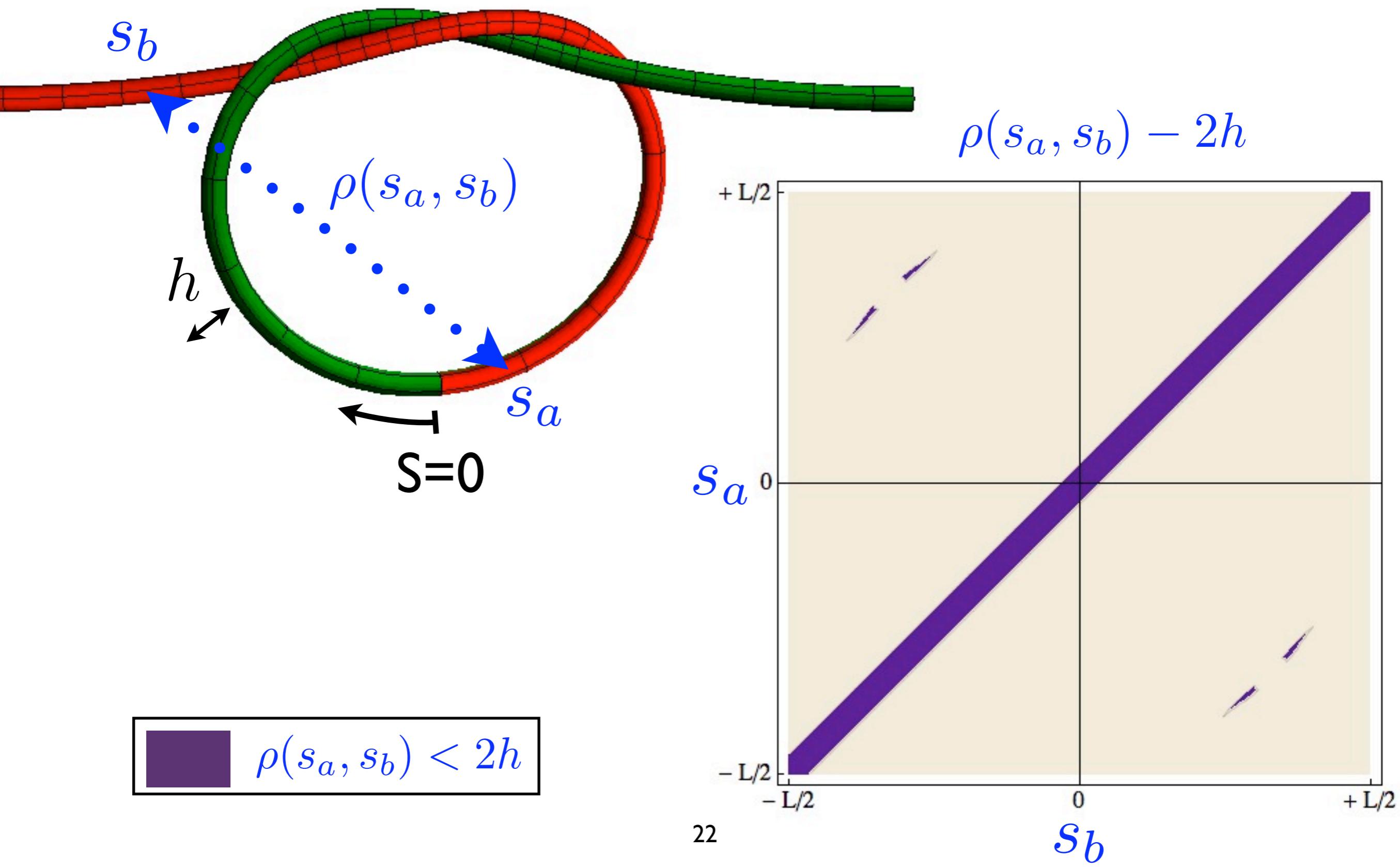
$$t = \frac{TL^2}{(2\pi)^2 EI}$$



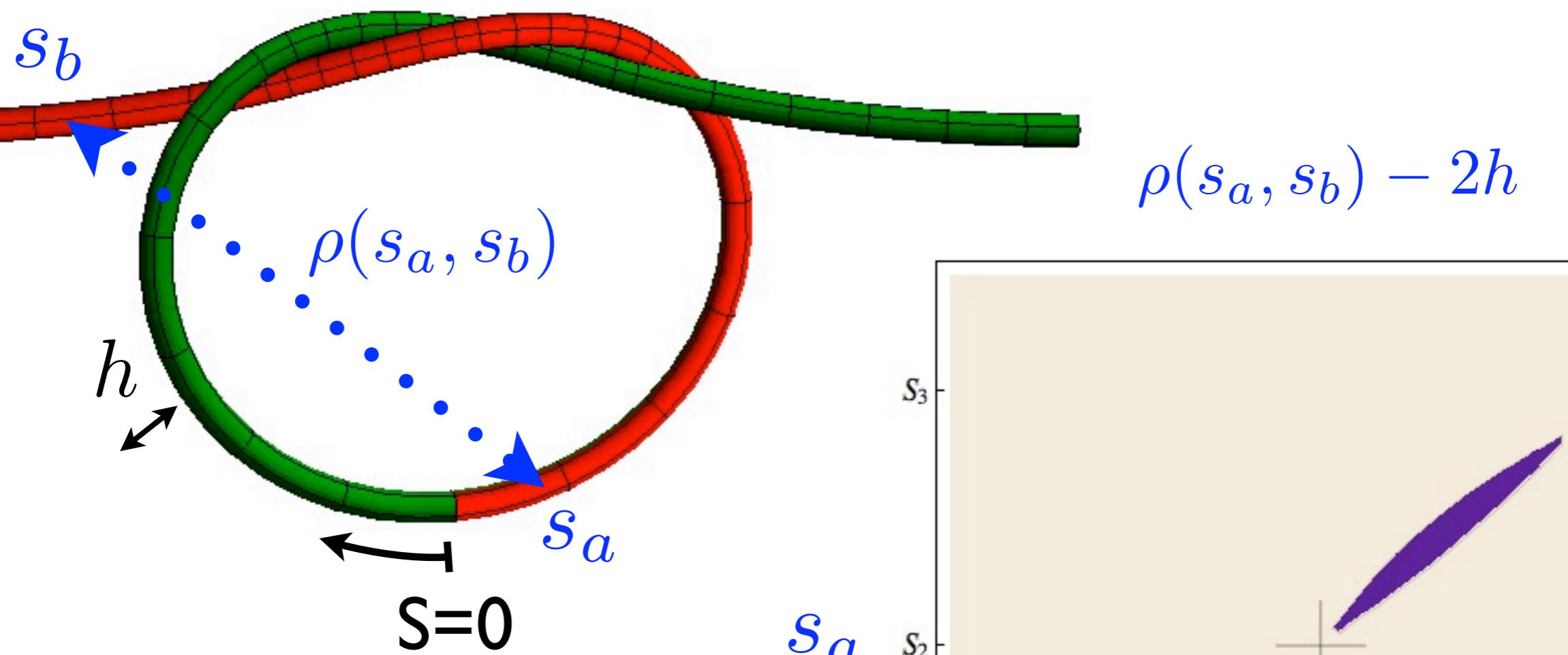
$$t = \left(\frac{2/\pi}{d} \right)^2$$

$$d = \frac{D}{L}$$

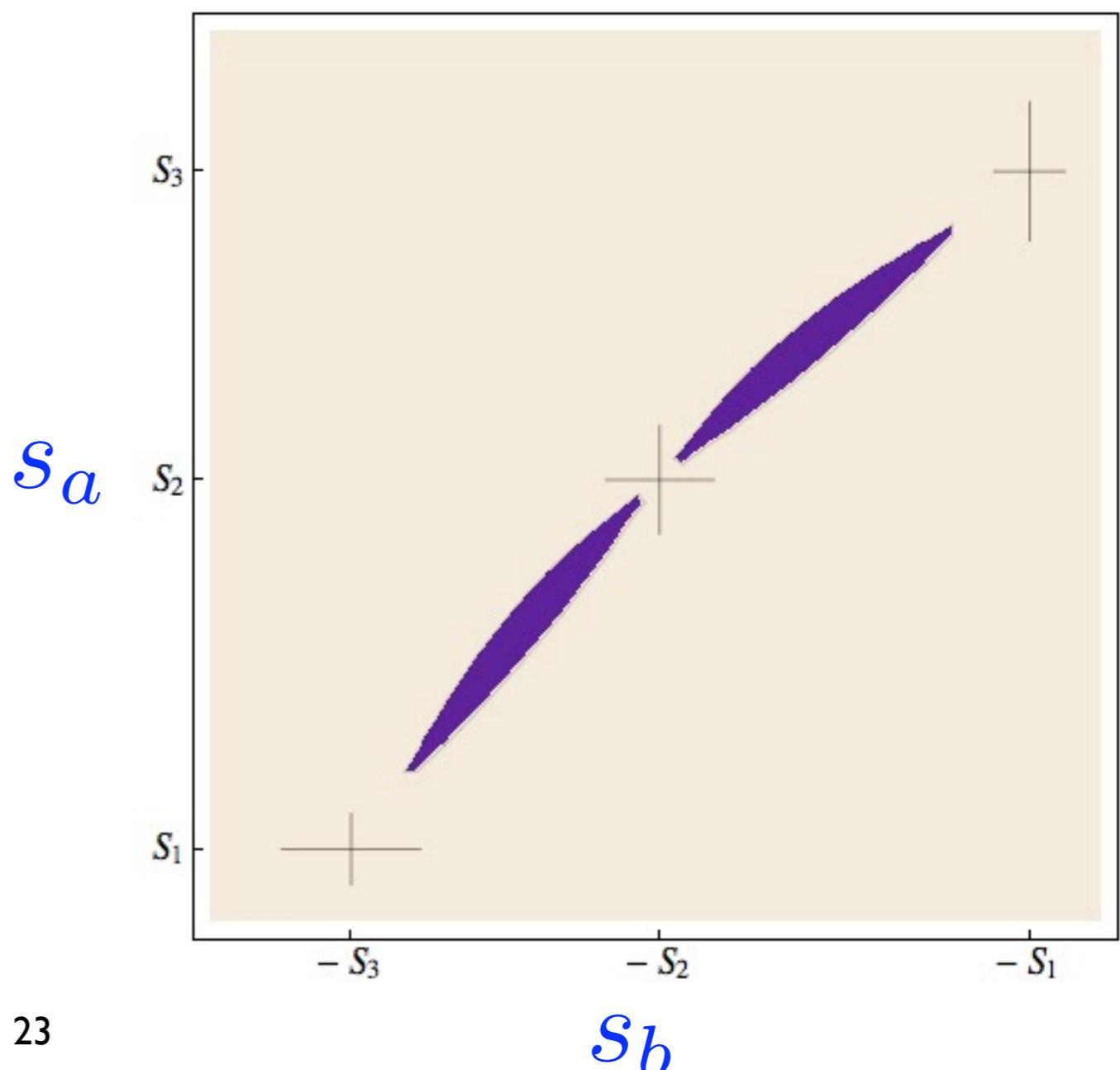
Distance of self-approach



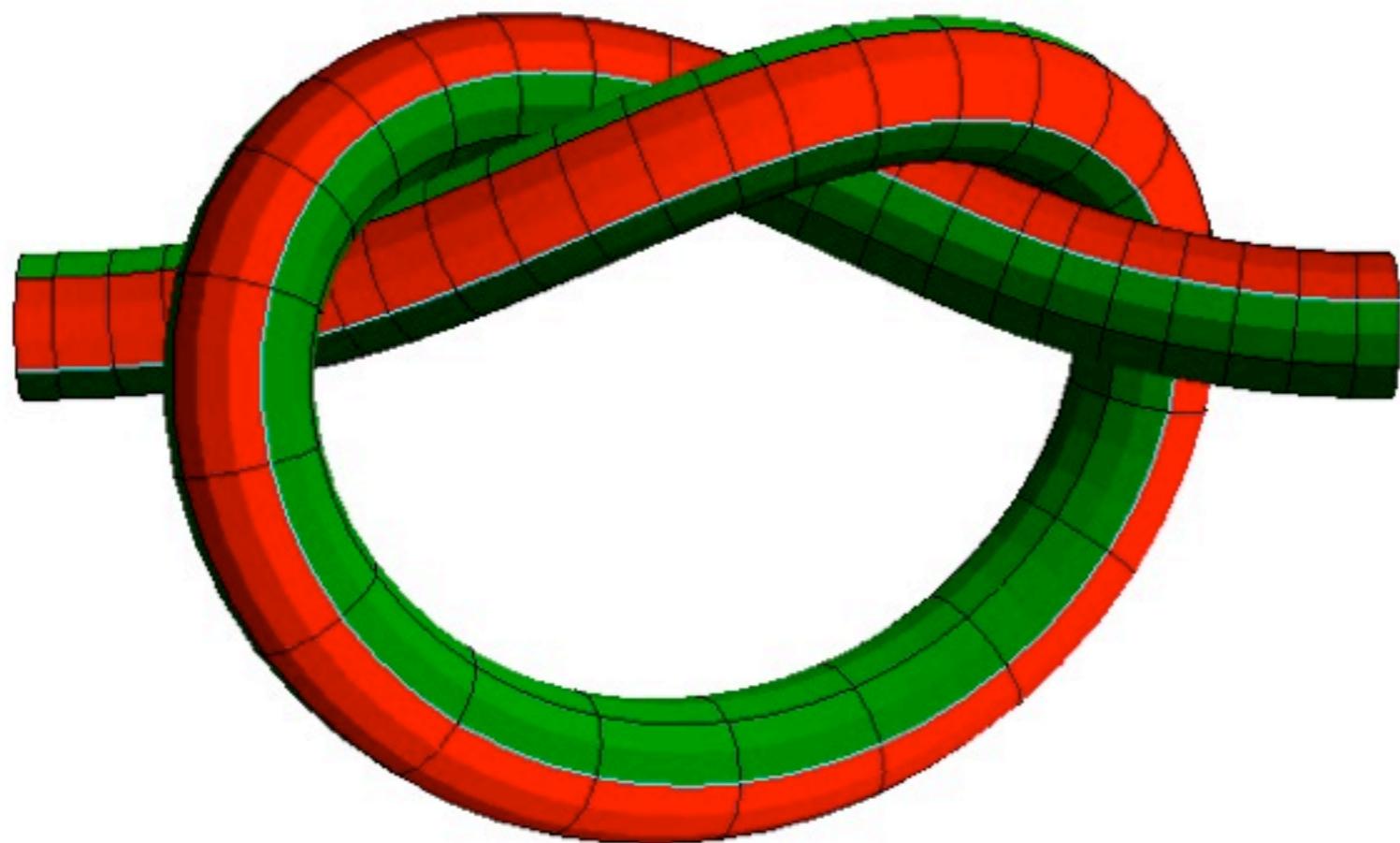
Distance of self-approach



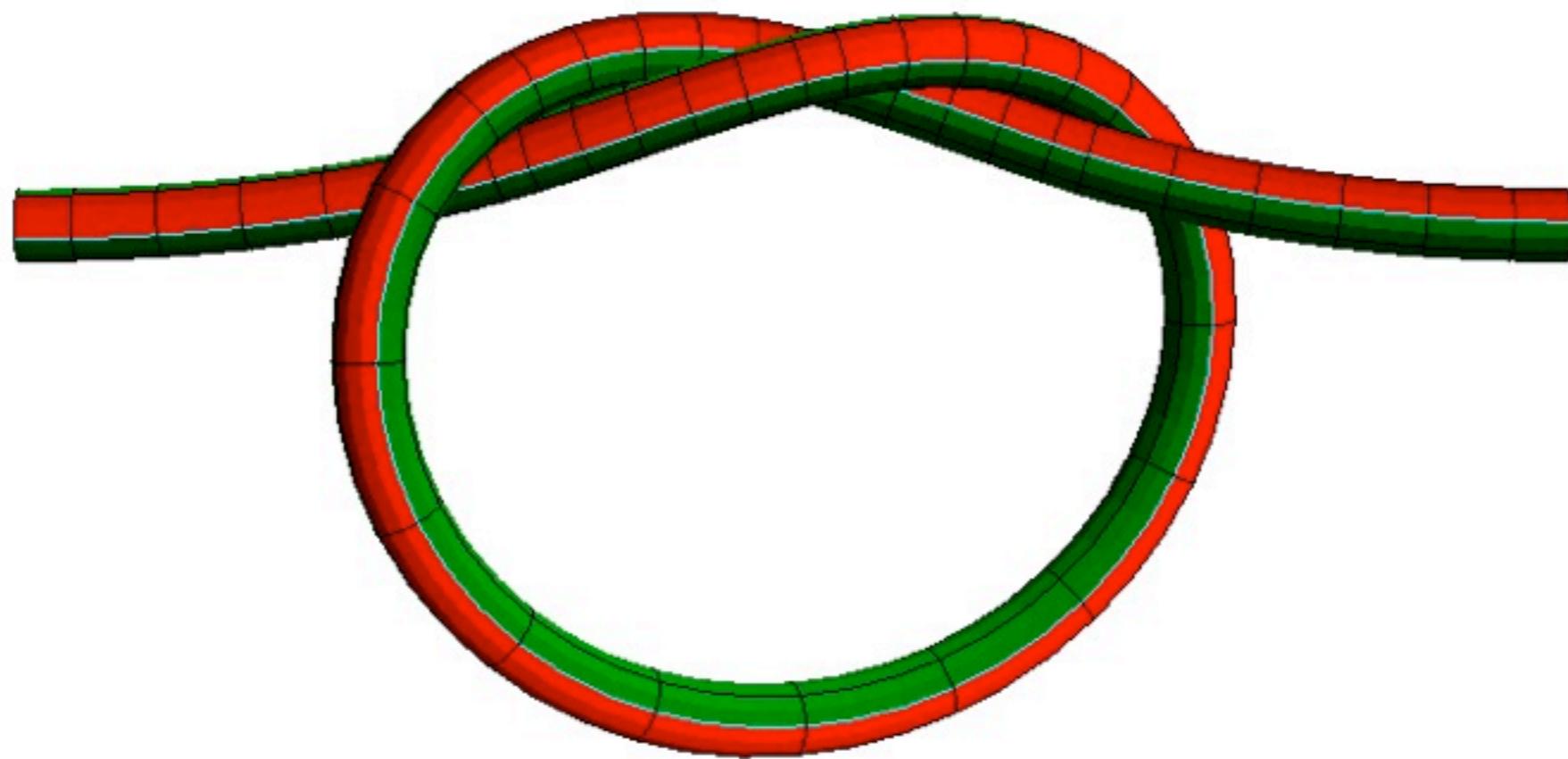
$\rho(s_a, s_b) < 2h$



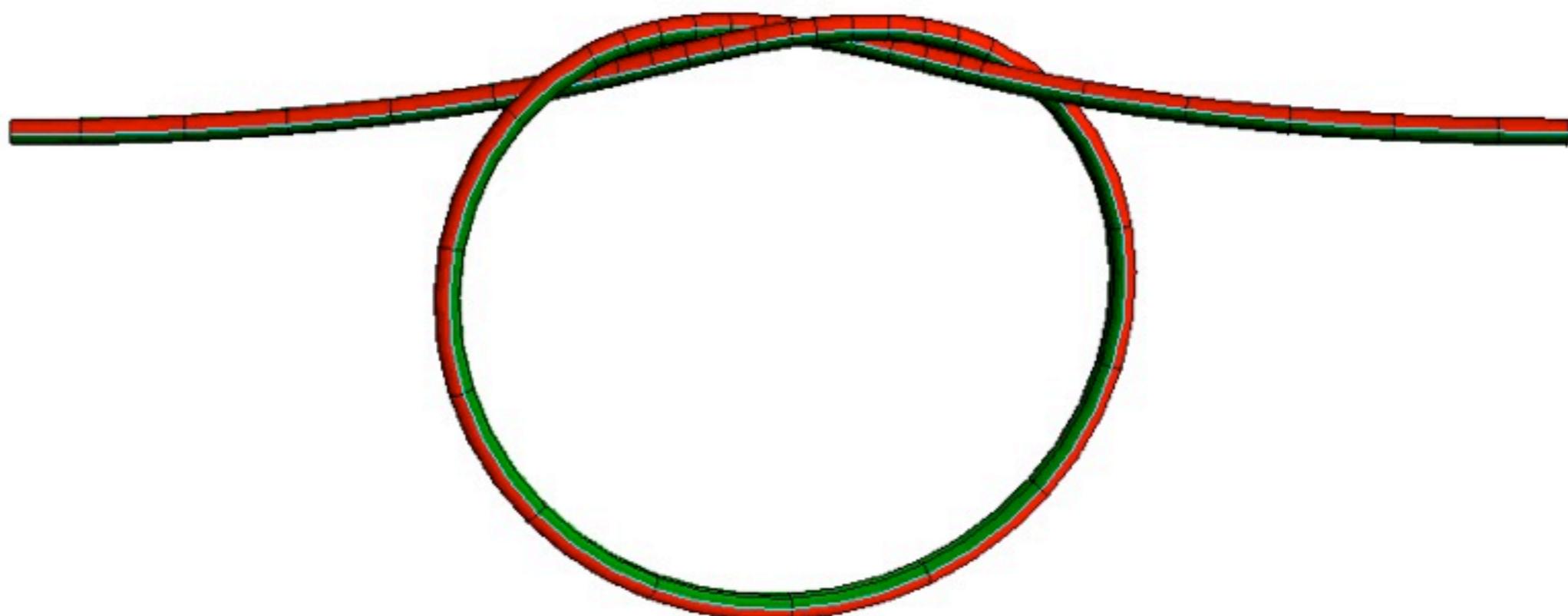
Making the rod thinner



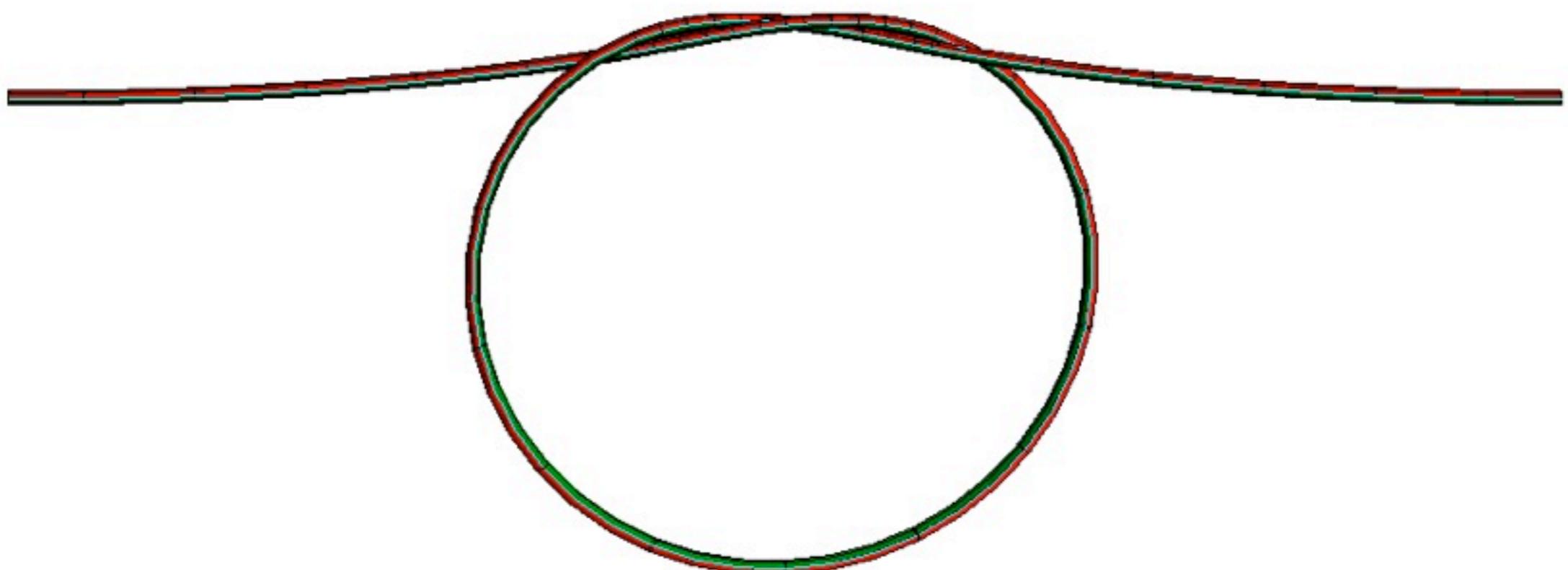
Making the rod thinner



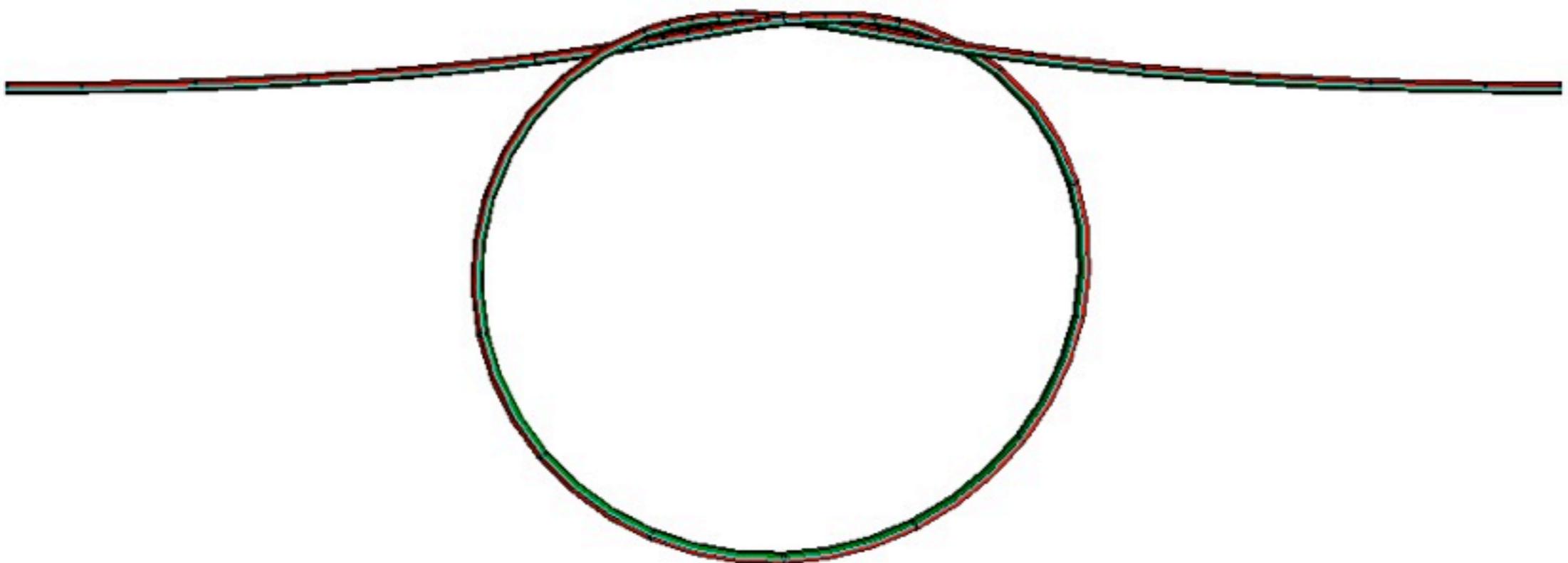
Making the rod thinner



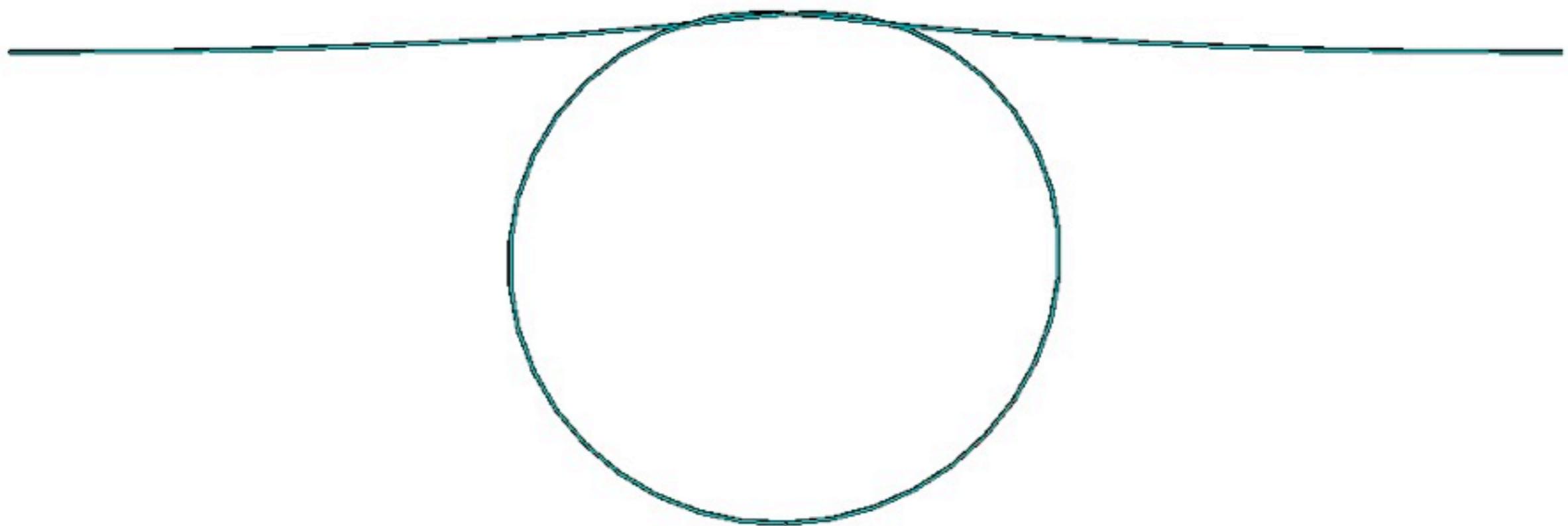
Making the rod thinner



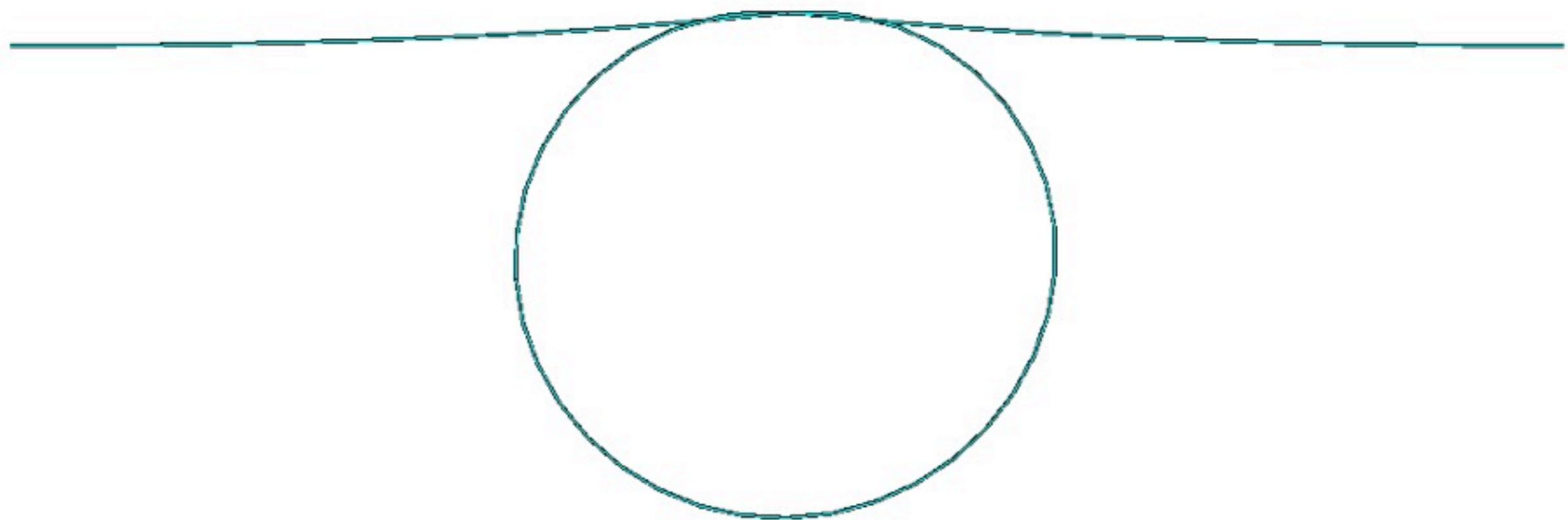
Making the rod thinner



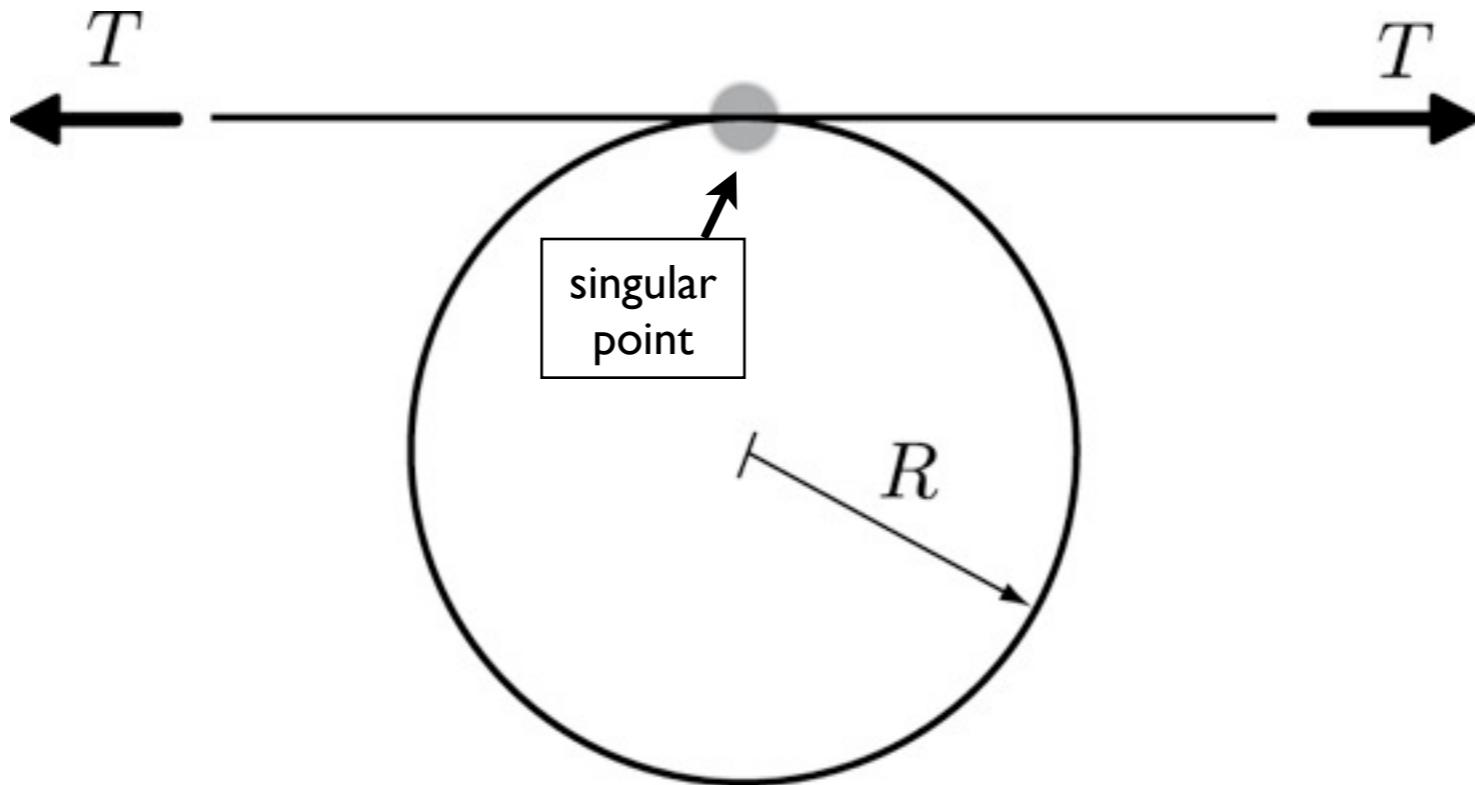
Making the rod thinner



Making the rod thinner



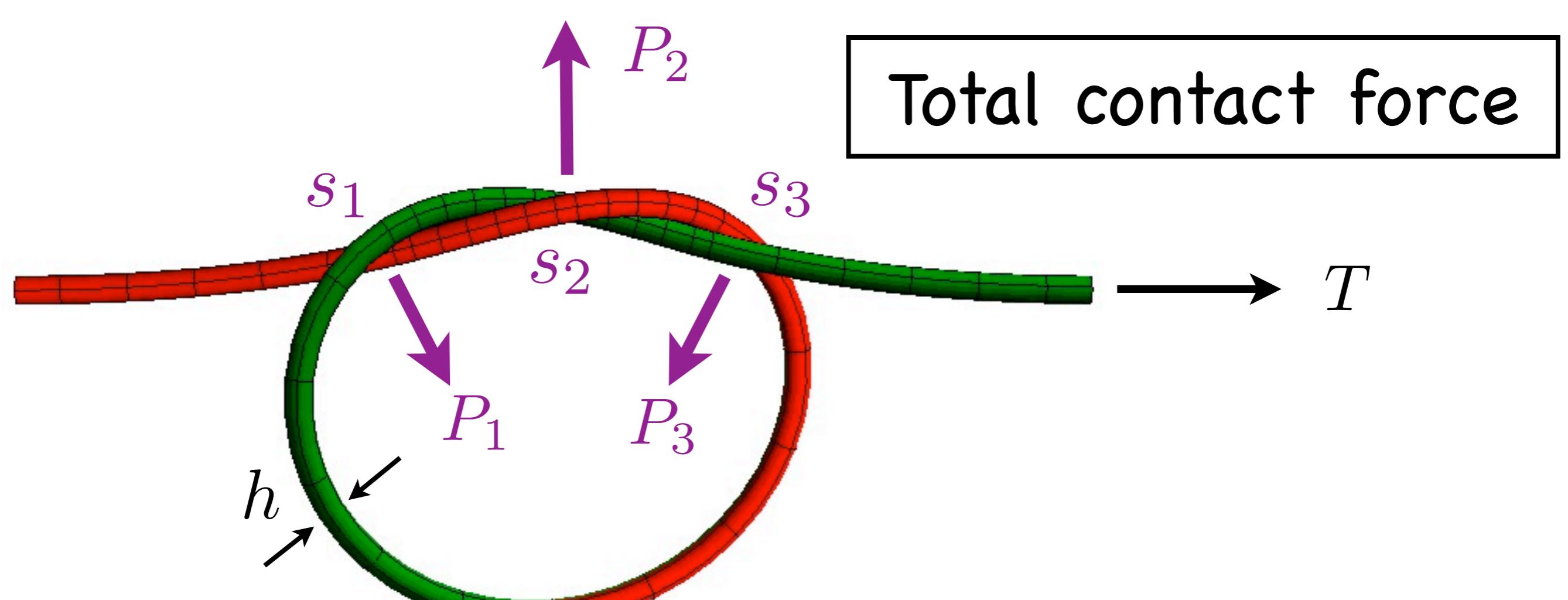
Zero thickness limit



$$\text{equilibrium : } T = \frac{EI}{2R^2}$$

Arai et al (1999)

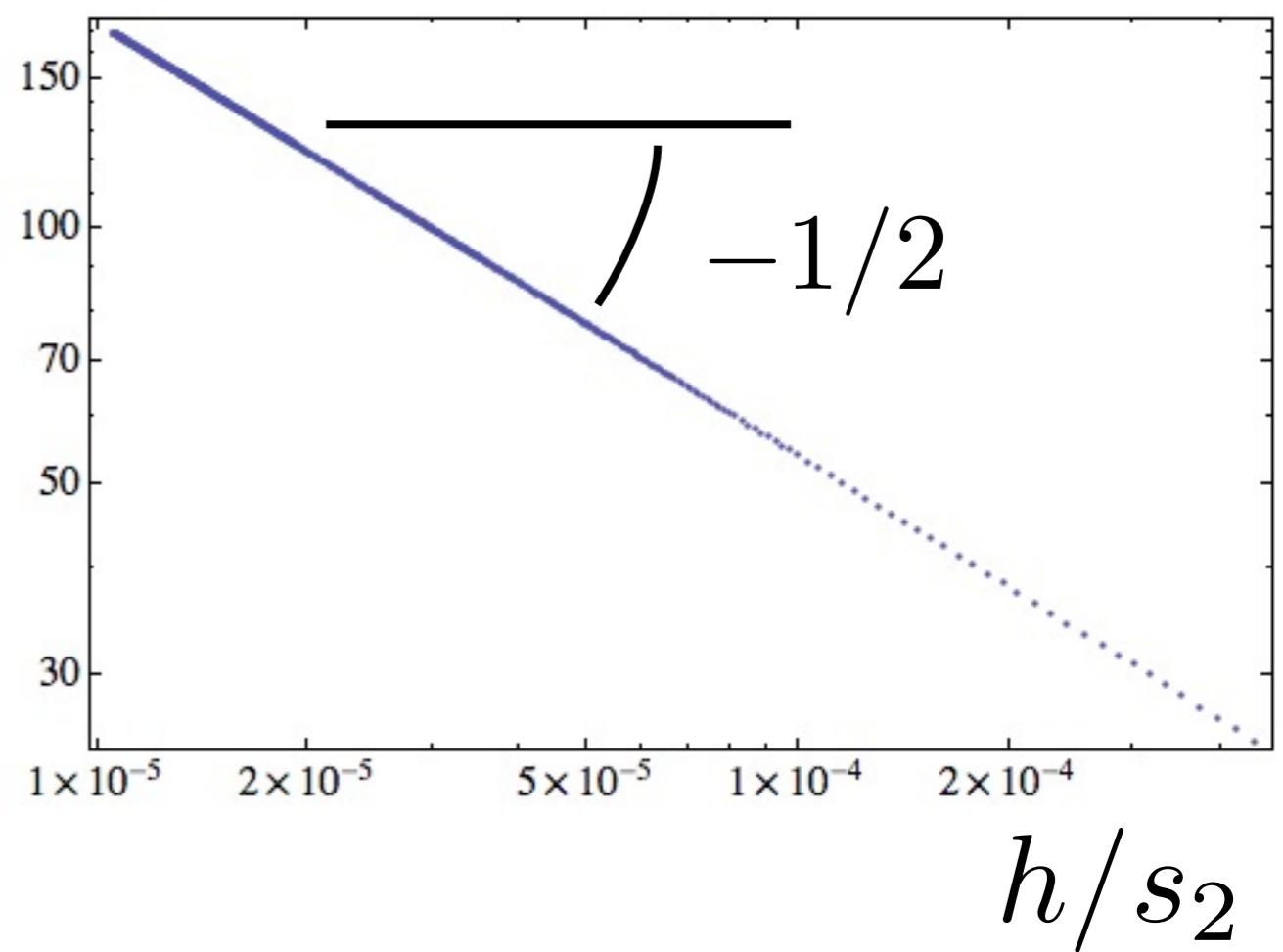
tensile force T  bending moment $\frac{EI}{R}$



Total contact force

$$\frac{1}{T} \sum_i P_i$$

$\boxed{\frac{1}{T} \sum_i P_i \simeq 0.55 (h/s_2)^{-1/2}}$



Kirchhoff Equations

$$\left\{ \begin{array}{l} \vec{F}' = -\vec{p} \\ \vec{M}' = \vec{F} \times \vec{t} \\ \vec{t}' = \frac{1}{EI} \vec{M} \times \vec{t} \\ \vec{R}' = \vec{t} \end{array} \right. \quad \begin{array}{l} \text{forces equil.} \\ \text{moments equil.} \\ \text{kinematics} \\ \text{tangent def.} \end{array}$$

$$' \equiv \frac{d}{ds}$$

$\vec{p}(s)$ ext. pressure

$\vec{F}(s)$ internal force

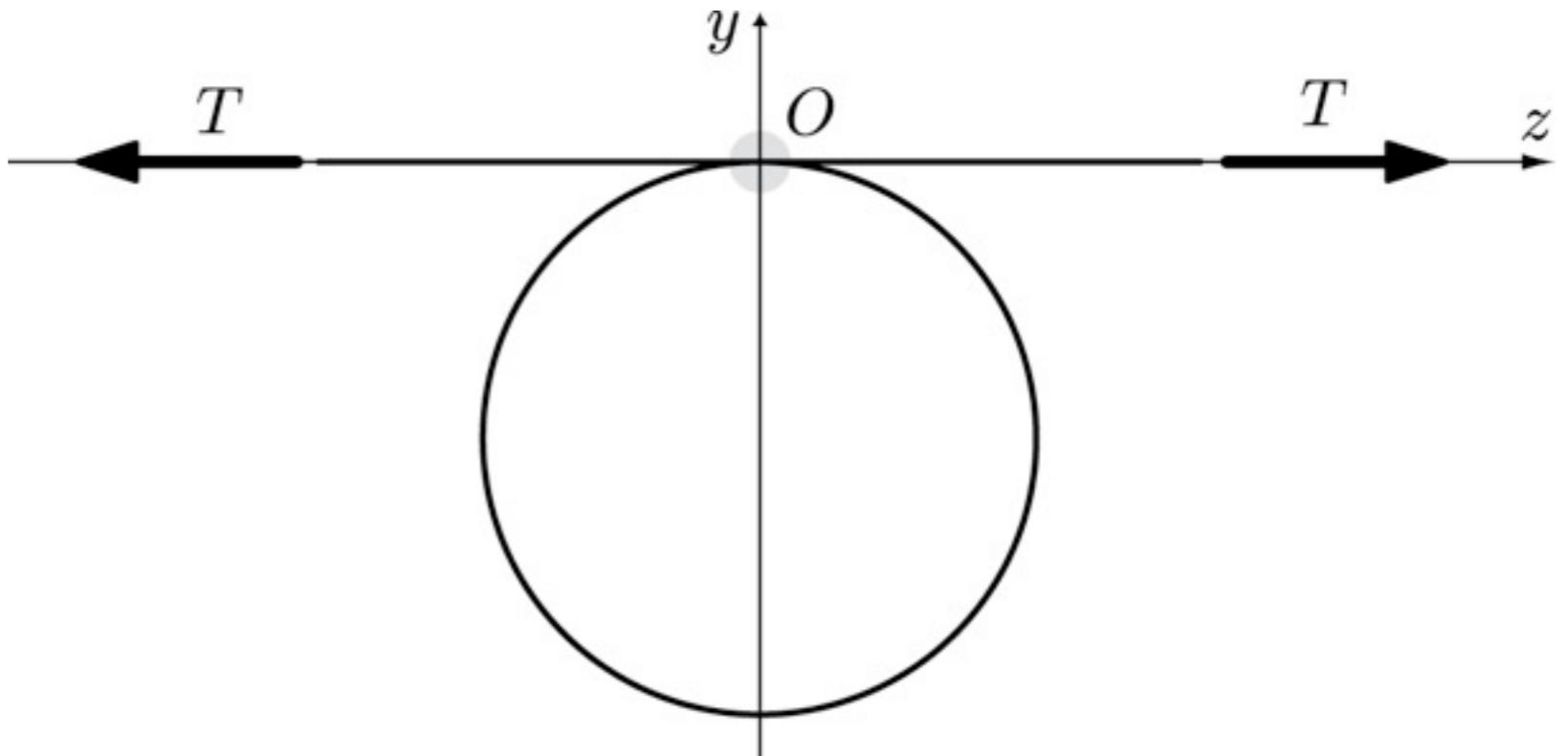
$\vec{M}(s)$ internal moment

$\vec{R}(s)$ position

$\vec{t}(s)$ tangent

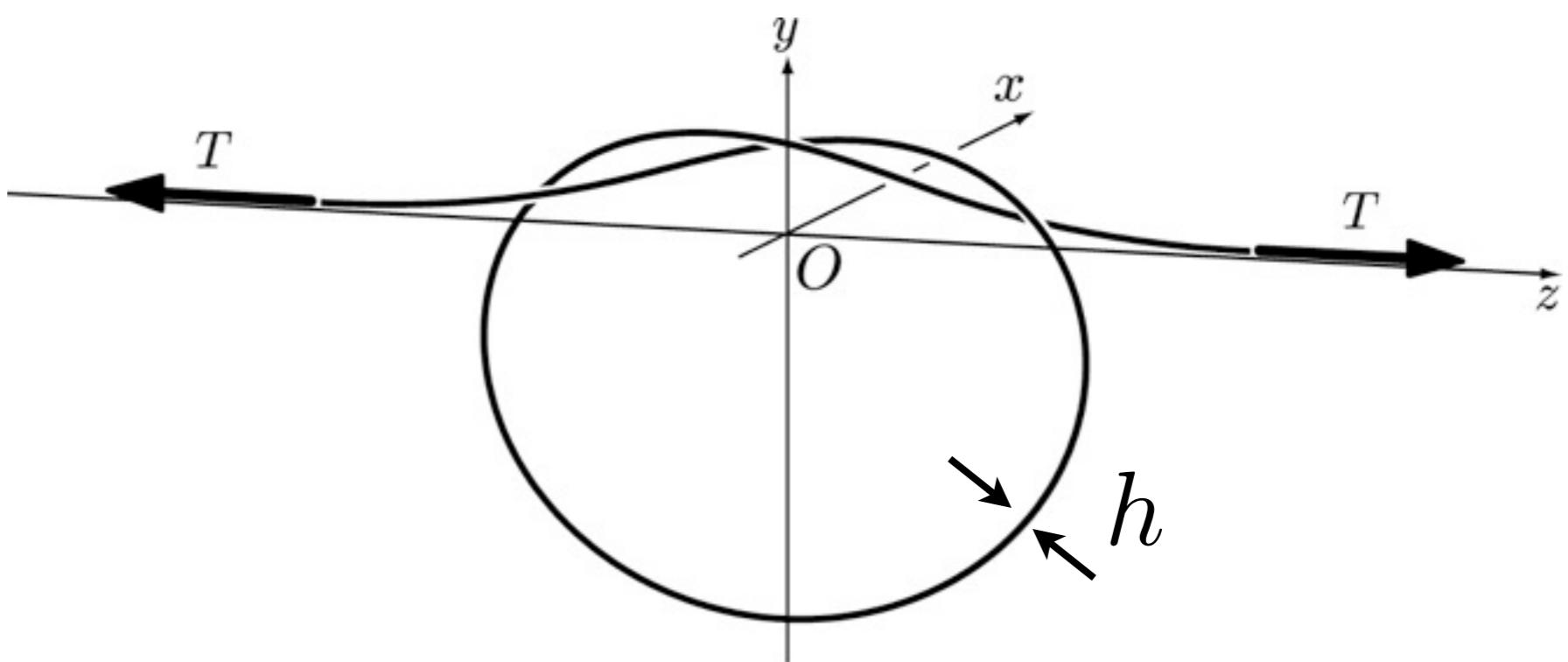
Perturbative problem

$$\epsilon = 0 \\ (h = 0)$$

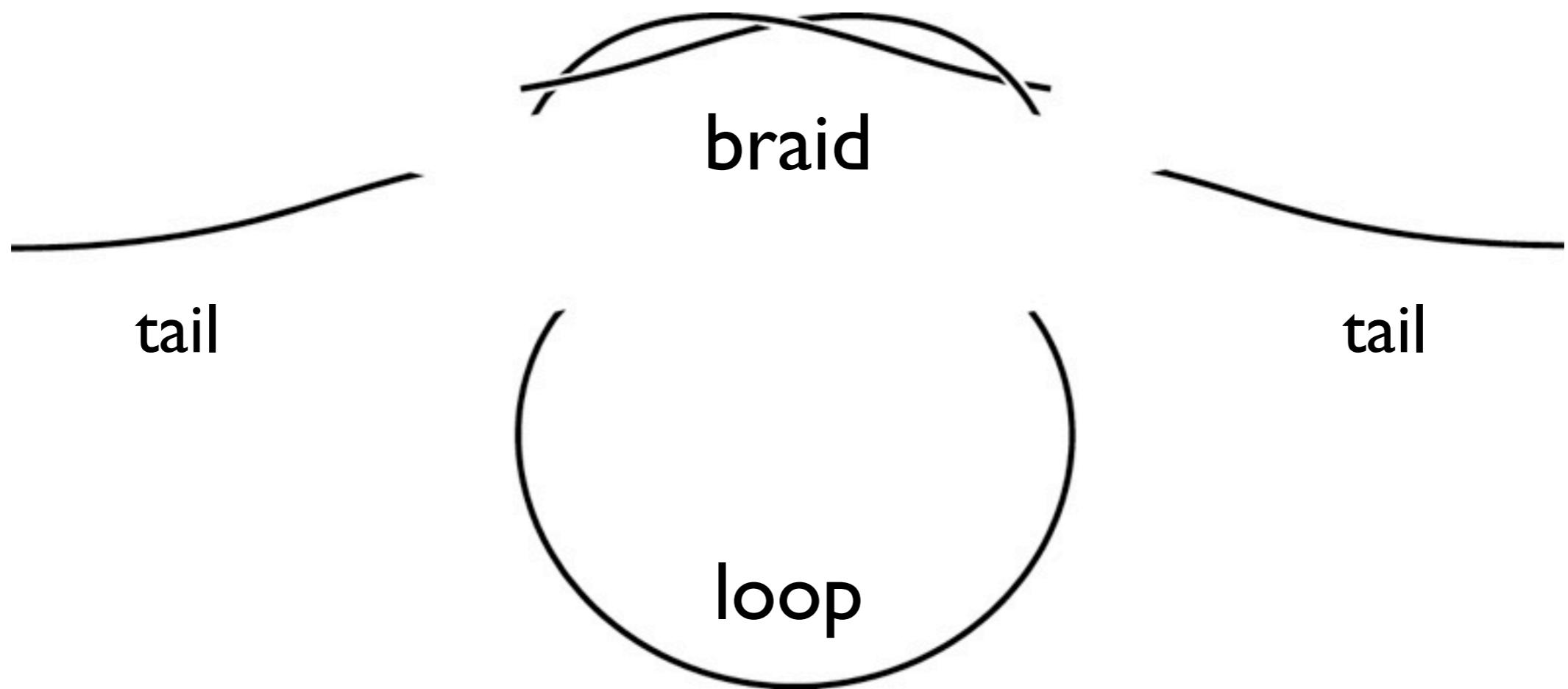


small parameter

$$\epsilon = \left(\frac{2h^2T}{EI} \right)^{1/4} \ll 1$$

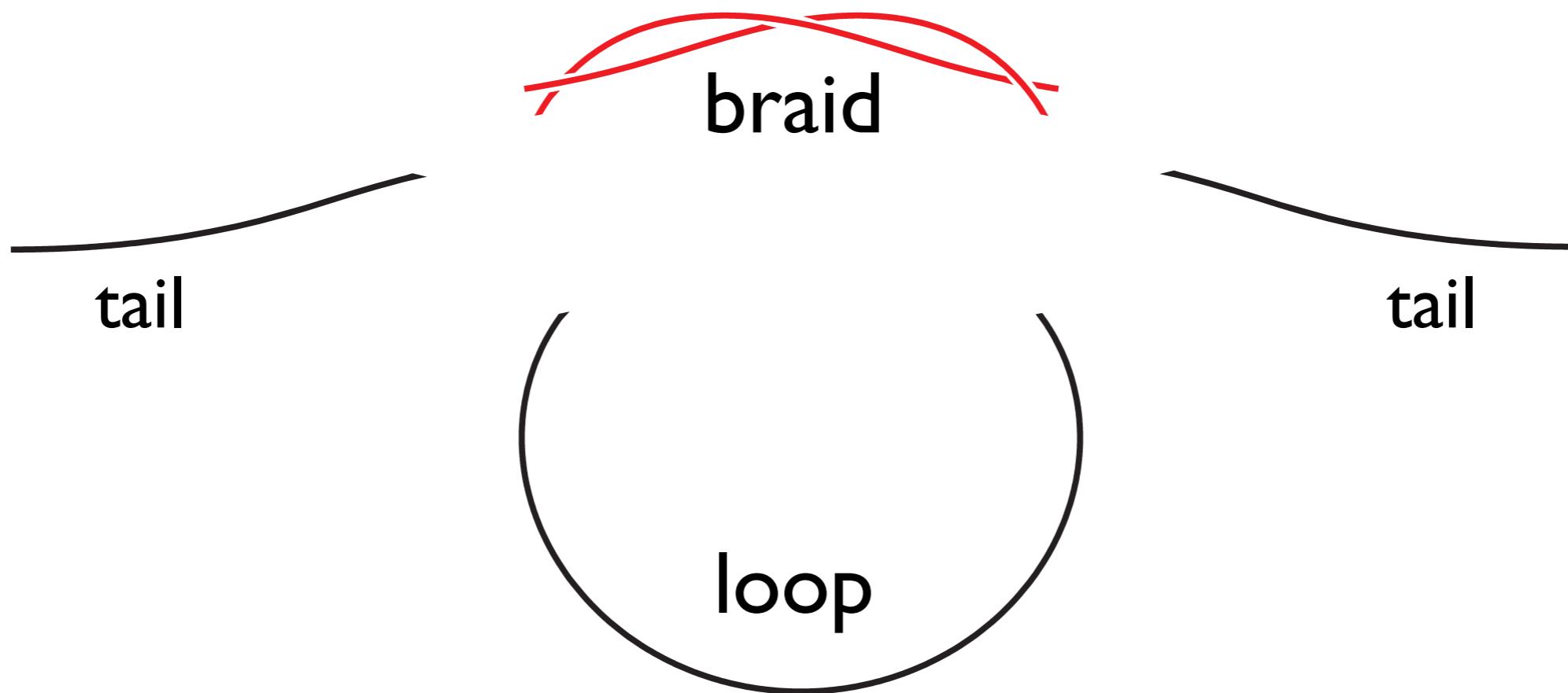


Matched asymptotic expansions

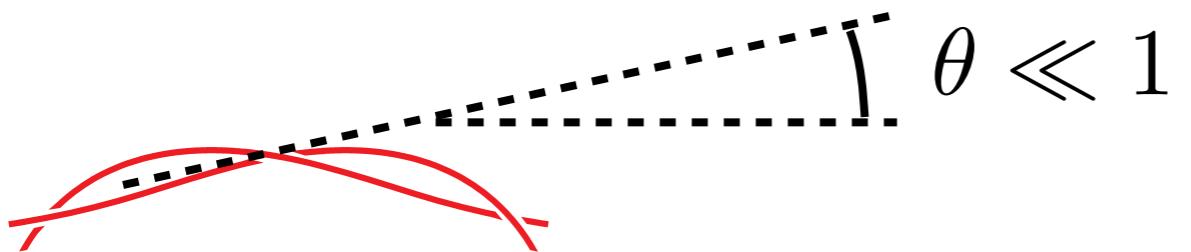


small parameter : $\epsilon = \left(\frac{2h^2T}{EI} \right)^{1/4} \ll 1$

Braid : self-contact zone



Braid : linear superposition

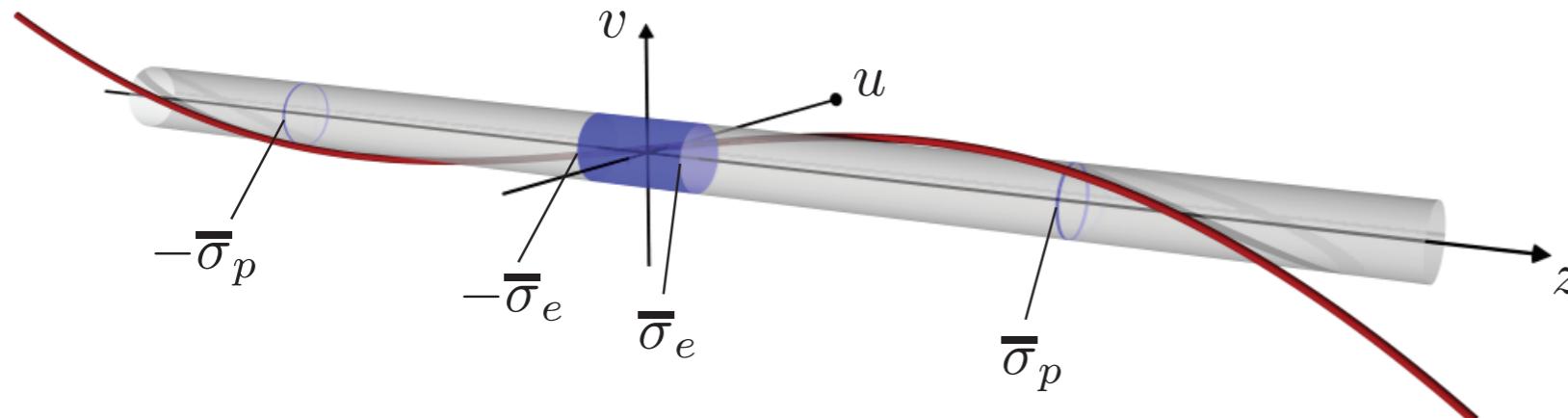


small deflections => linear problem

$$x^A, x^B = x^B + x^A + x^B - x^A$$

The diagram illustrates the decomposition of a braid into a sum of simpler components. On the left, two red curves representing rigid bodies x^A and x^B are shown. An equals sign follows. To the right of the equals sign is a red curve representing the sum $x^B + x^A$. Below this sum is a plus sign ($+$) enclosed in a circle. To the right of the plus sign is a red curve representing the difference $x^B - x^A$, which forms a complex, knotted loop. A downward arrow points from the difference curve to the text "self-contact => contact with obstacle". To the left of the equals sign, another downward arrow points from the sum curve to the text "twice more rigid curvature: $\frac{1}{2} \frac{1}{R}$ ".

Braid : variational formulation



Kirchhoff equations => minimizing an energy

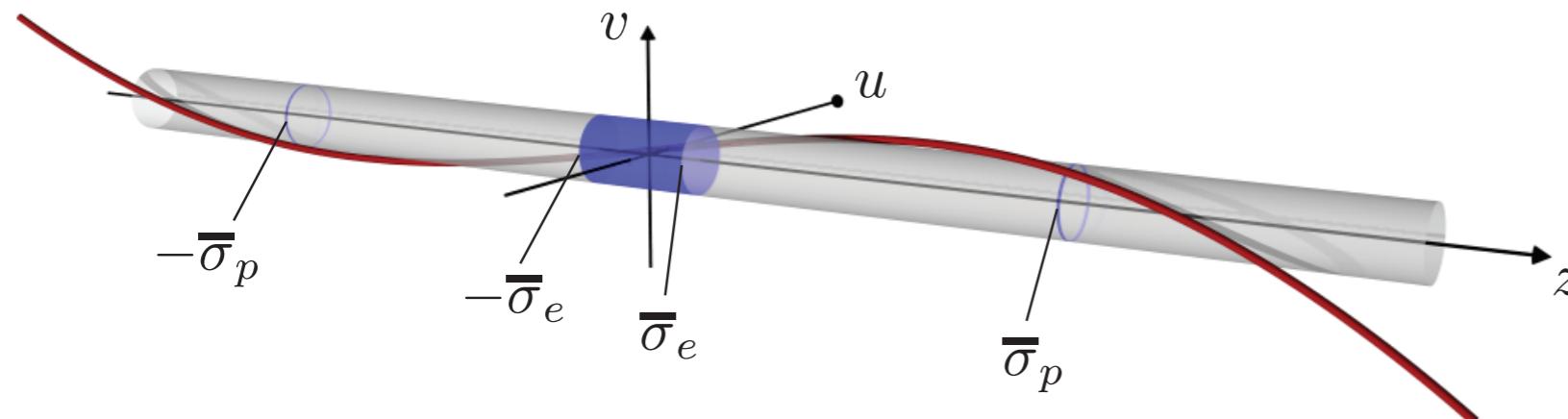
$$V = \frac{1}{2} \int_{-\infty}^{+\infty} \left(u''^2 + v''^2 \right) d\sigma + \underbrace{v'(+\infty) + v'(-\infty)}_{\text{work of external applied moments}}$$

with constraint:

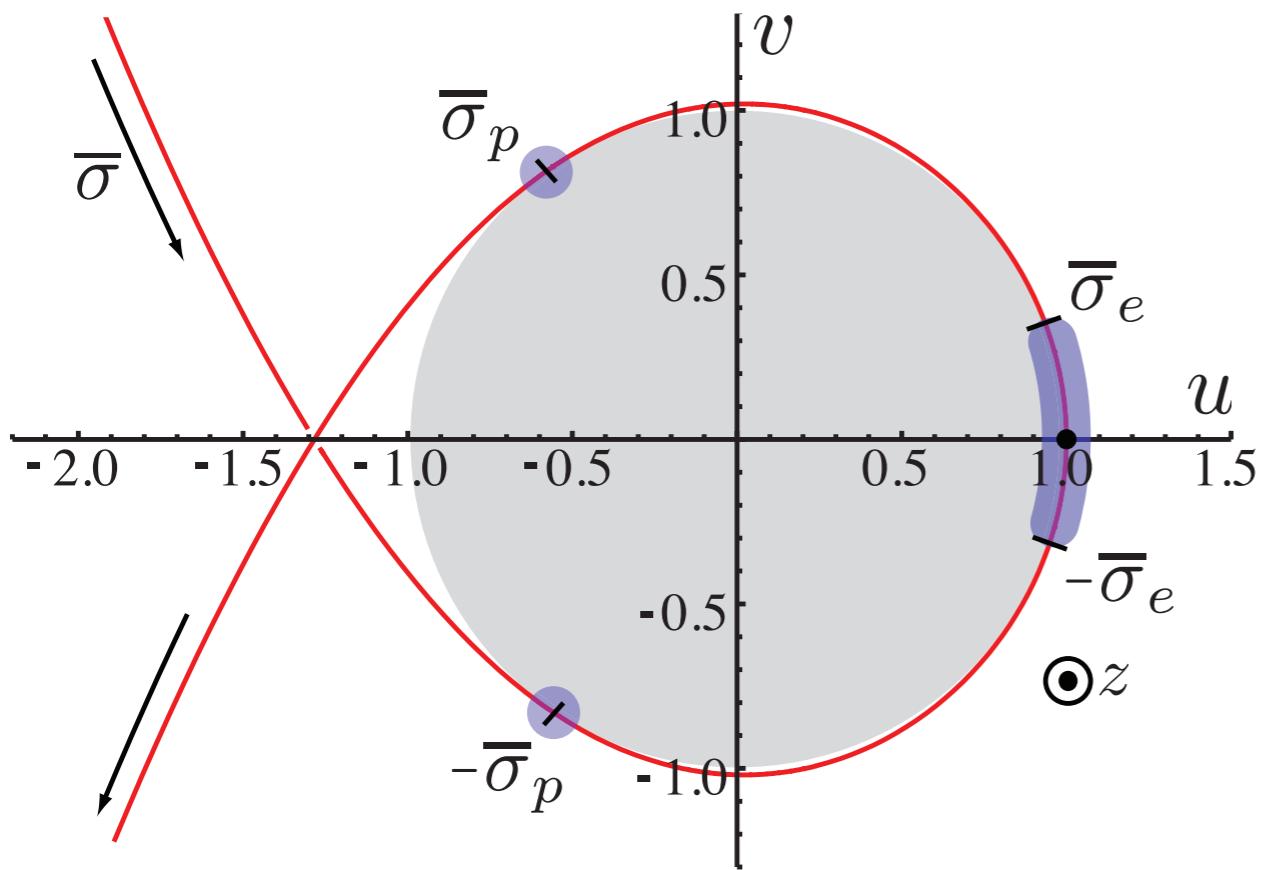
$$u^2(\sigma) + v^2(\sigma) \geq 1, \forall \sigma$$

work of external
applied moments

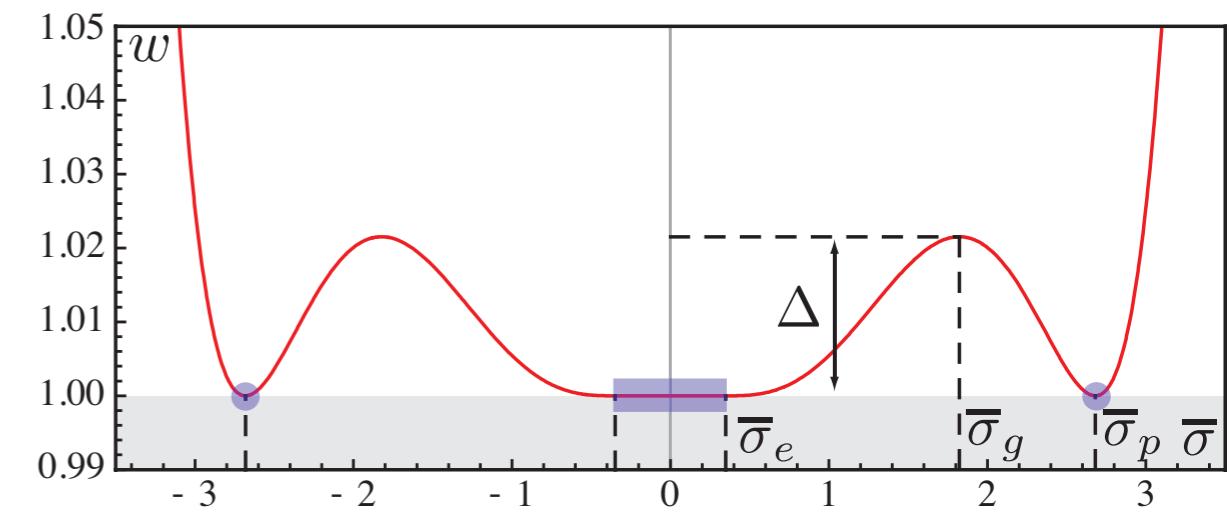
Braid : contact topology



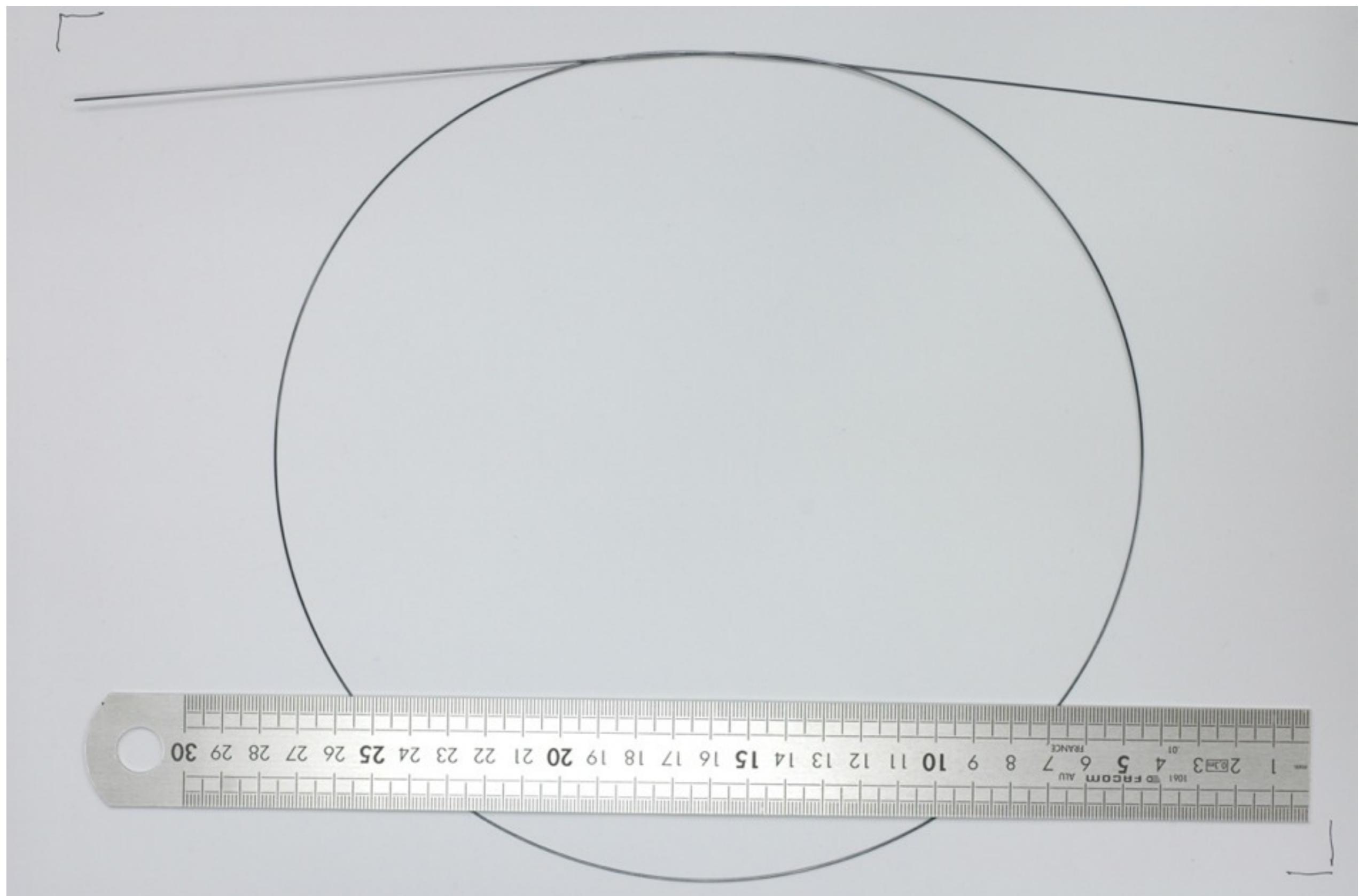
side view



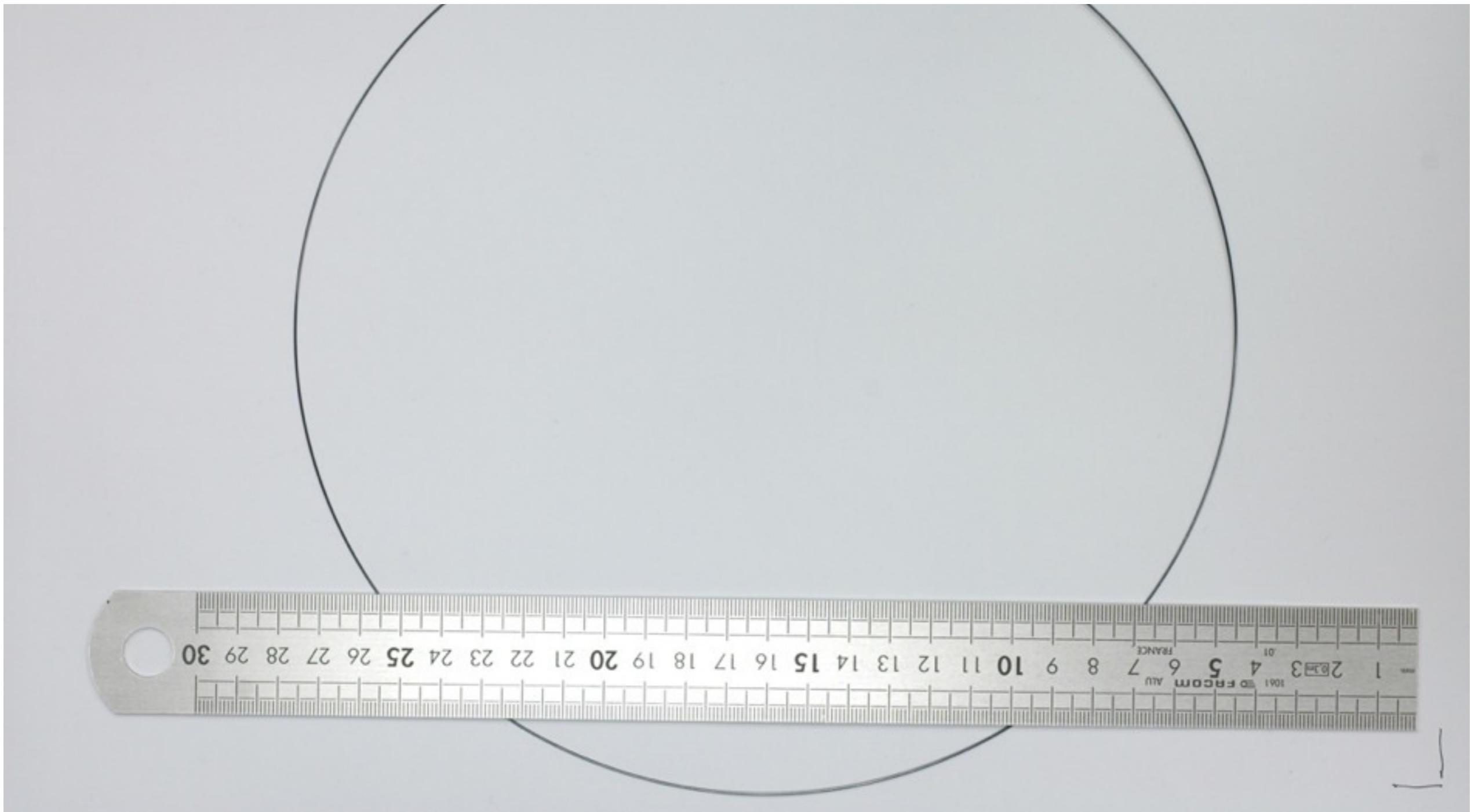
inter-strand distance



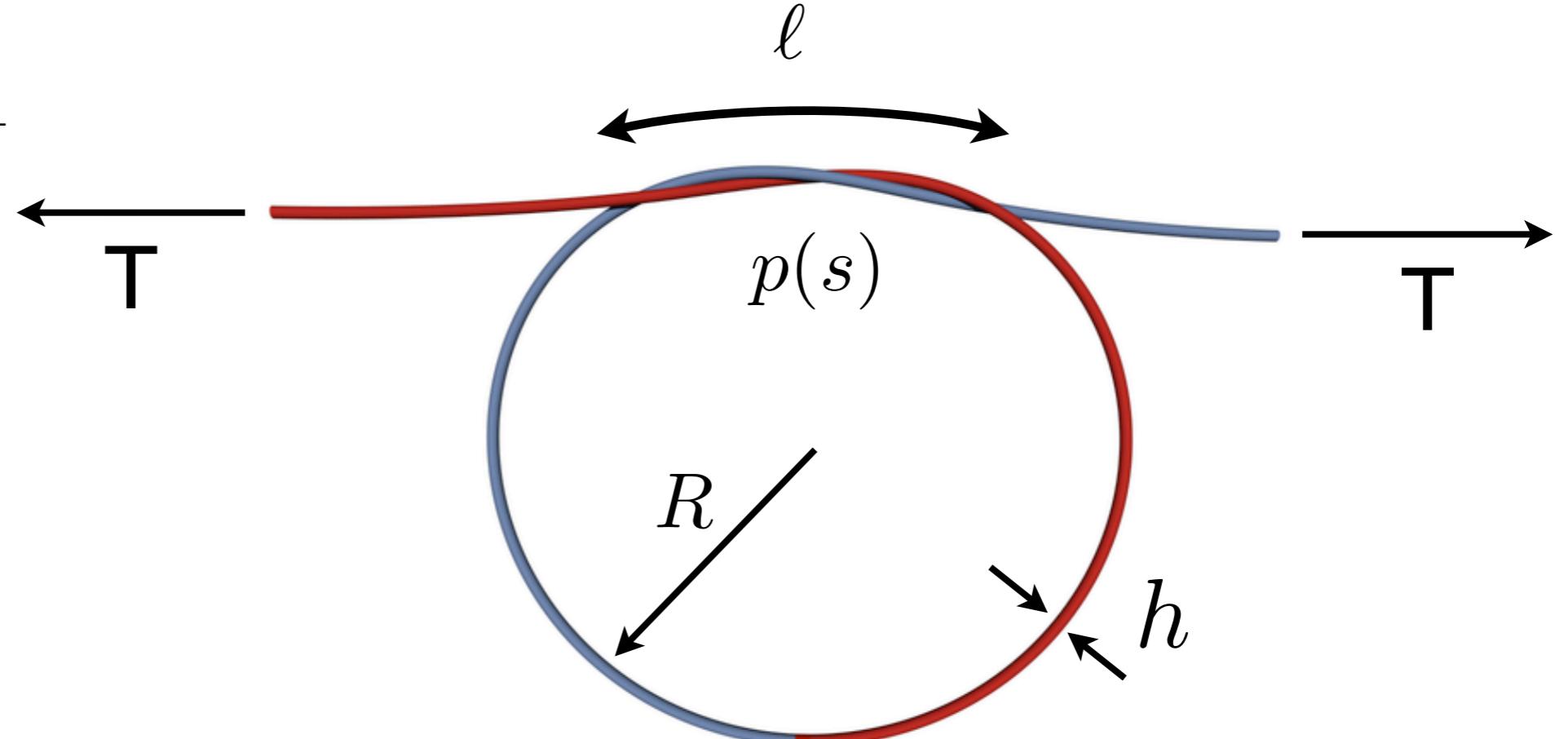
Braid : contact topology



Braid : contact topology



Results



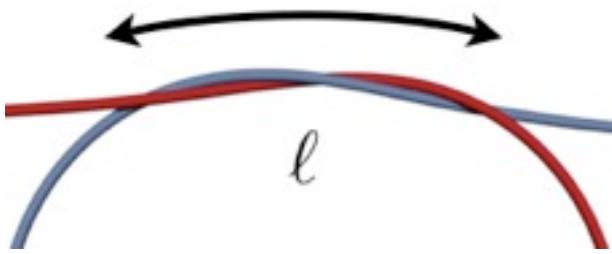
$$R = \sqrt{\frac{EI}{2T}}$$

$$\ell = 9.91 h^{1/2} (EI)^{1/4} T^{-1/4}$$

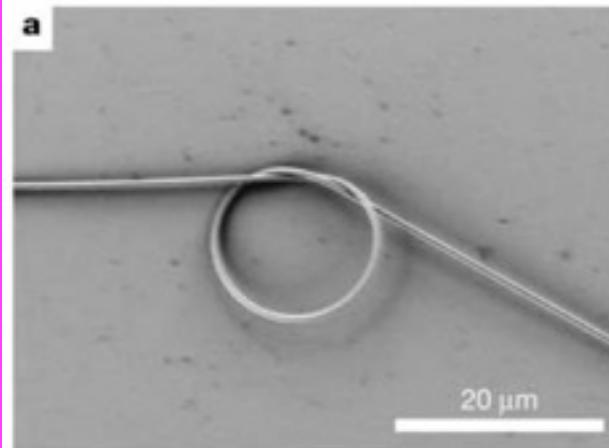
Contact pressure $p(s)$

Total contact force $P = \int_0^\ell p(s)ds = 0.82 h^{-1/2} (EI)^{1/4} T^{3/4}$

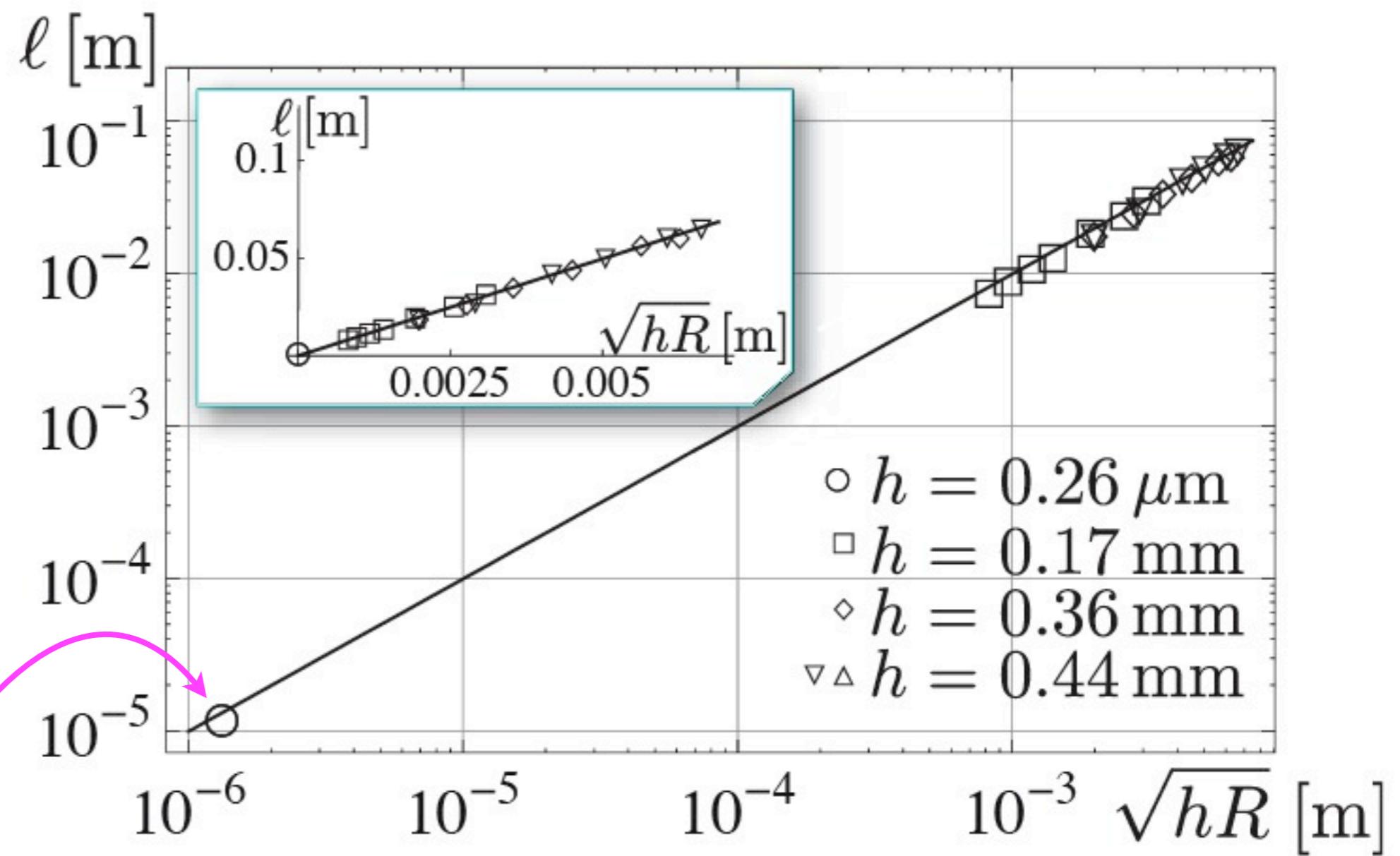
Experiments



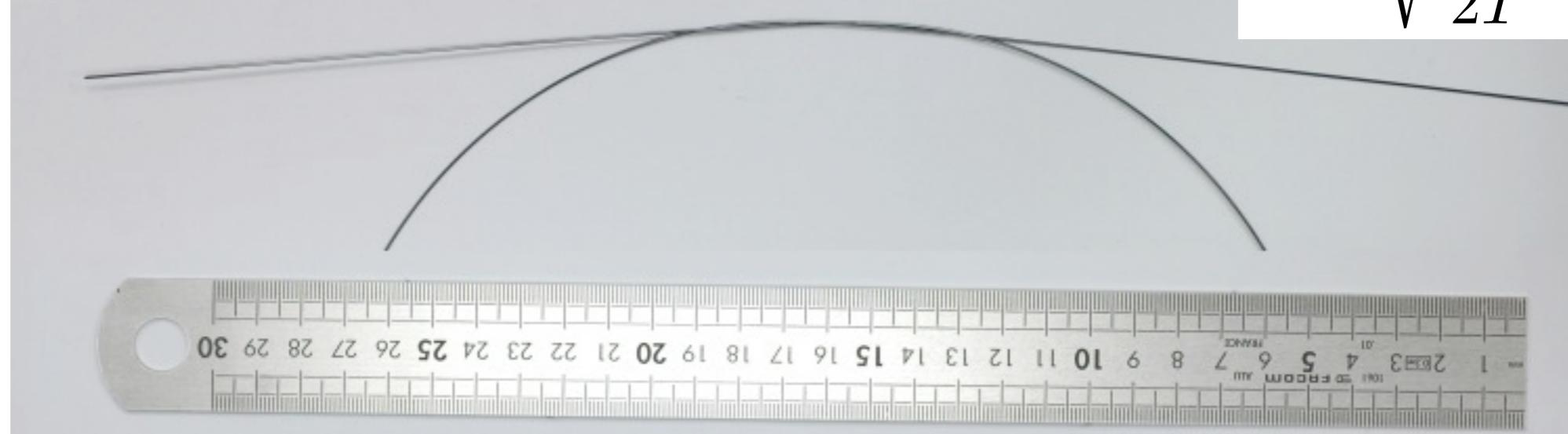
Tong et al., Nature 2003



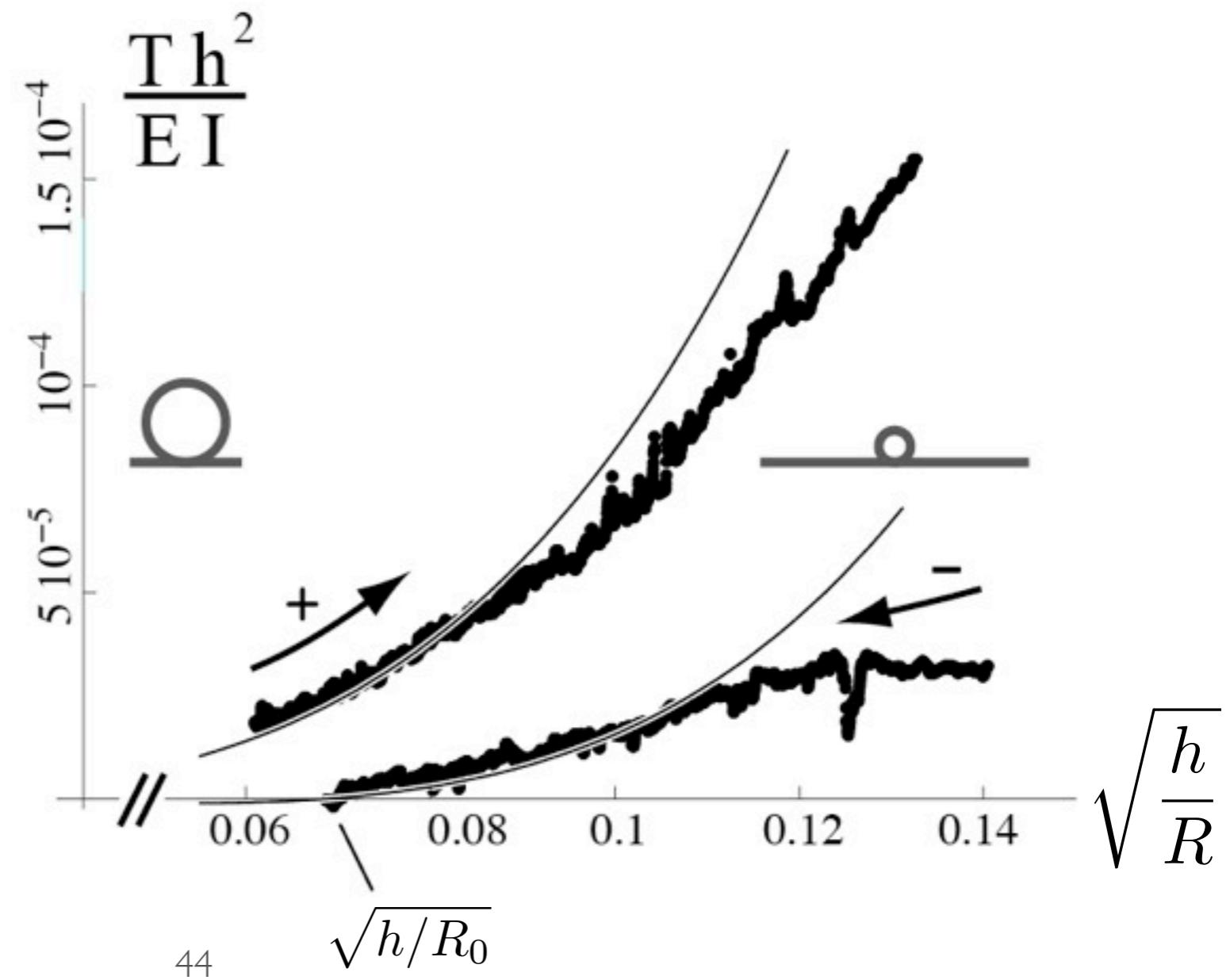
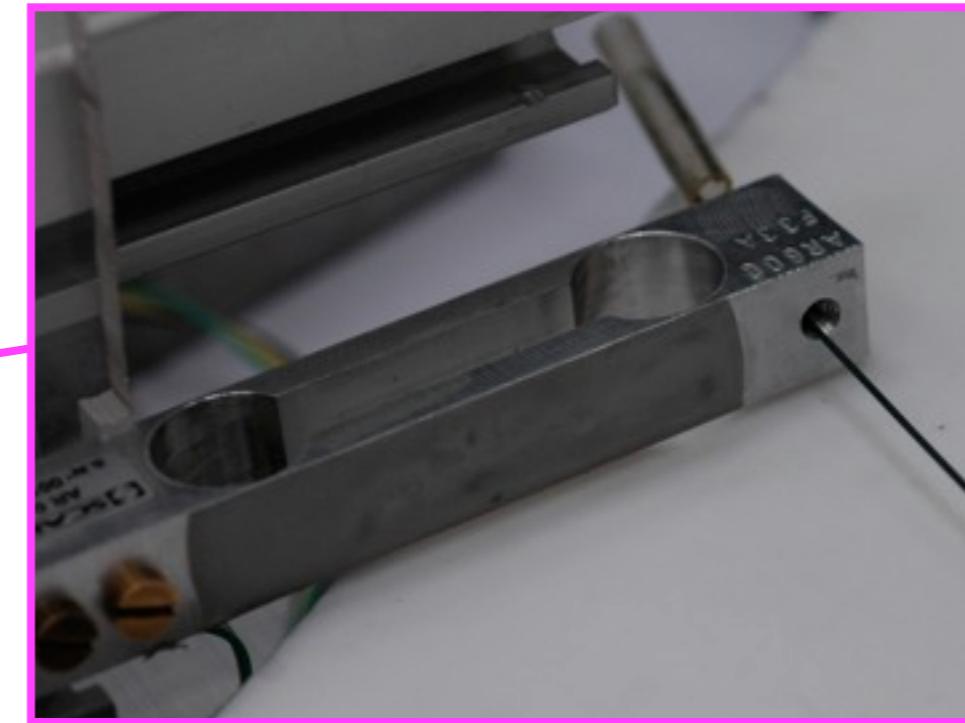
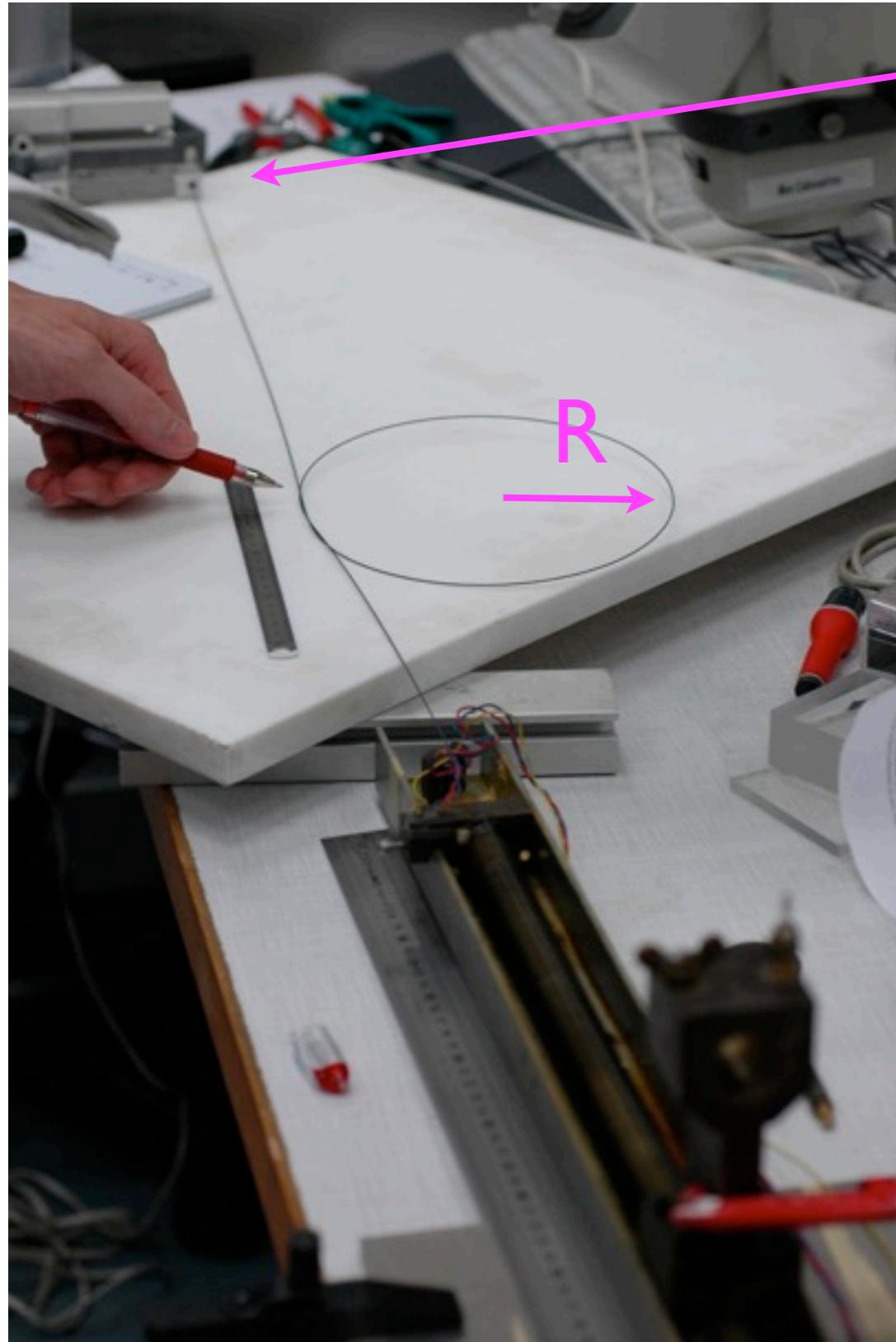
silica wire
 $h = 1/2 \text{ micron}$



$$R = \sqrt{\frac{EI}{2T}}$$



Experiments

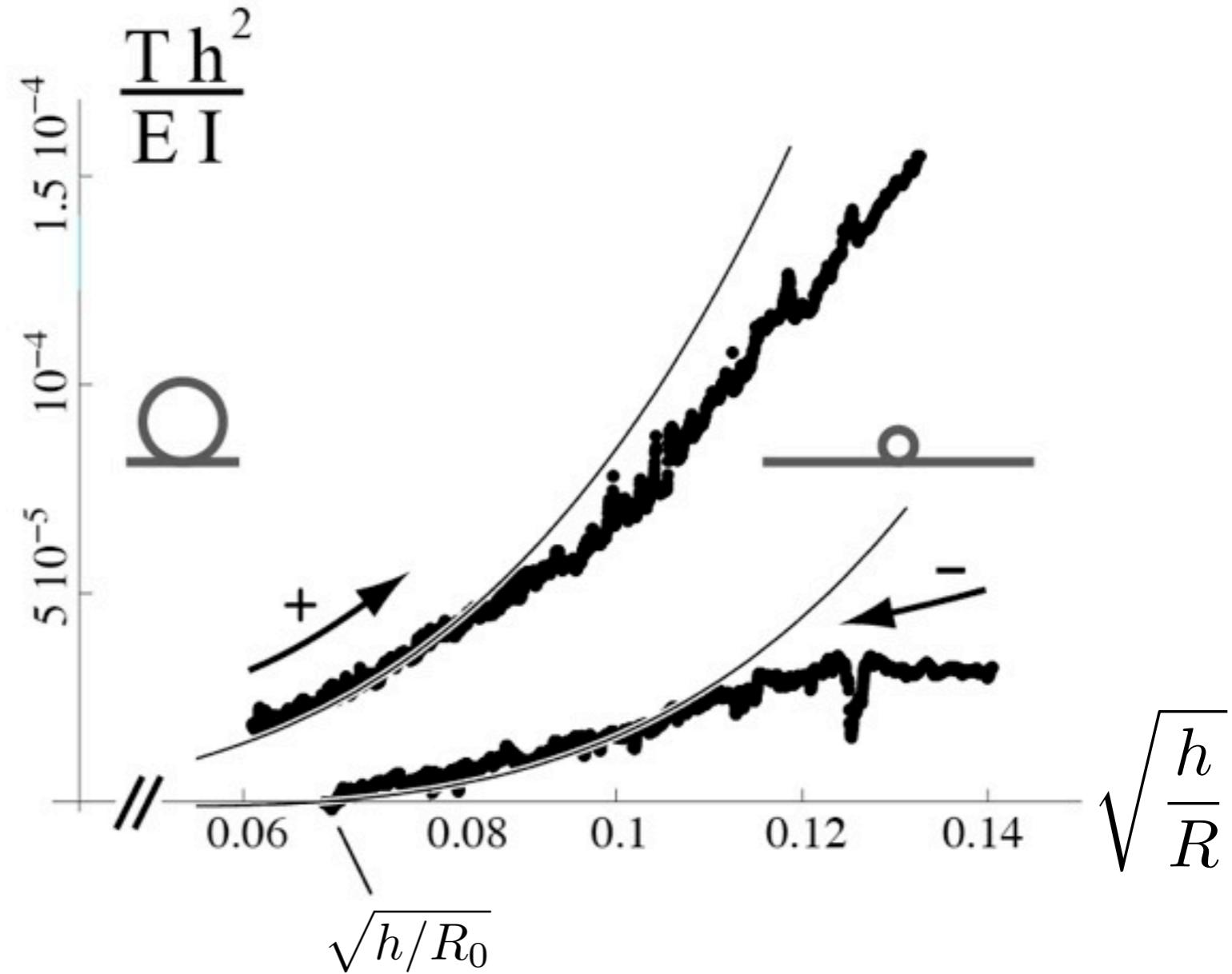


Expériences

sans frottement

$$T = \frac{1}{2} \frac{EI}{R^2}$$

$$\Rightarrow \frac{Th^2}{EI} = \frac{1}{2} \frac{h^2}{R^2}$$



avec frottement

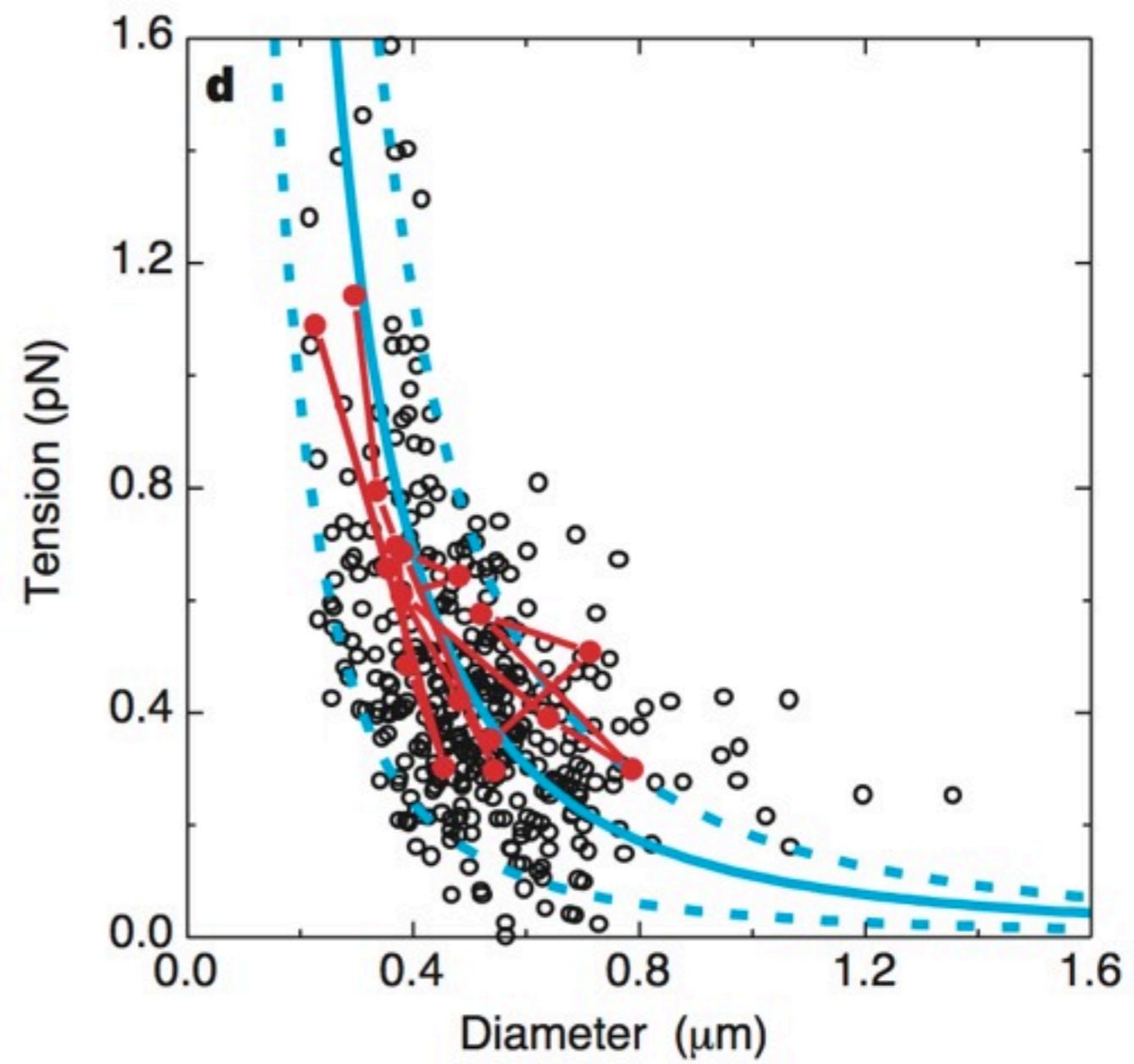
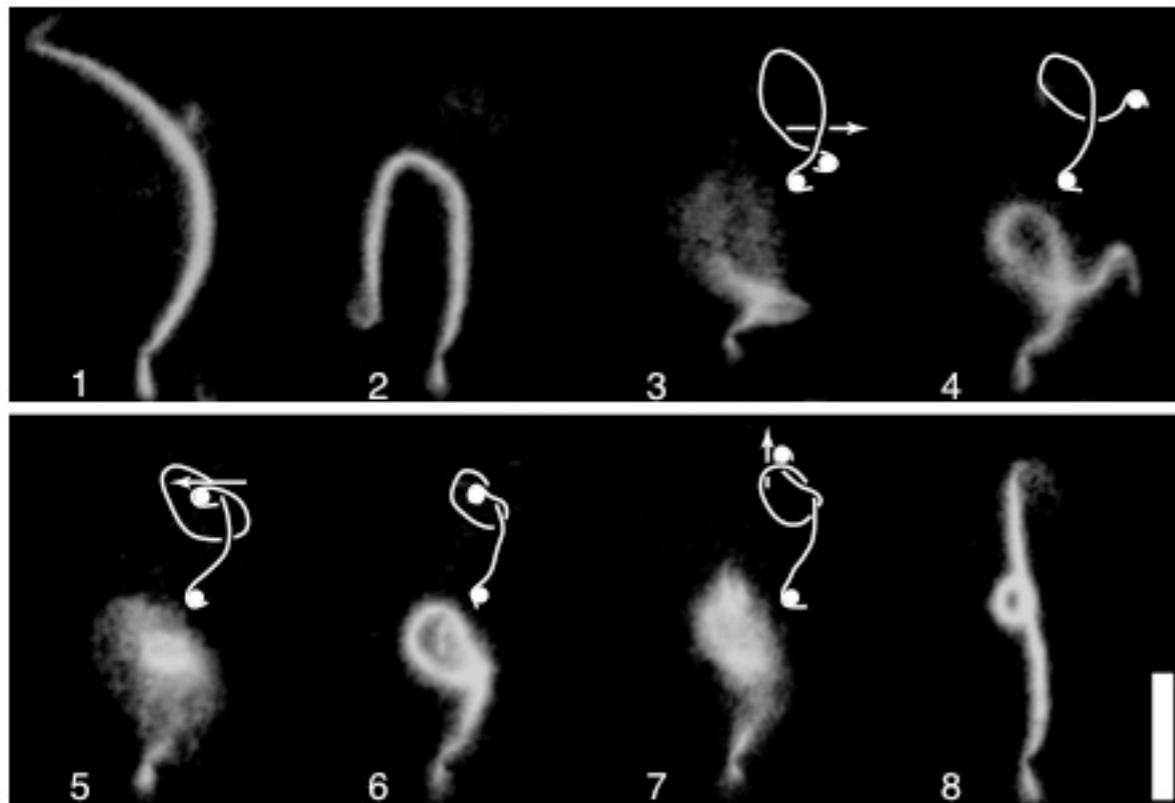
$$\left| \frac{Th^2}{EI} - \frac{1}{2} \frac{h^2}{R^2} \right| \leq \mu P = 0.49 \mu \left(\frac{h}{R} \right)^{3/2}$$

si $T = 0$: glissement jusqu'à $R = R_0$ tel que : $\mu = 1.02 \sqrt{\frac{h}{R_0}}$

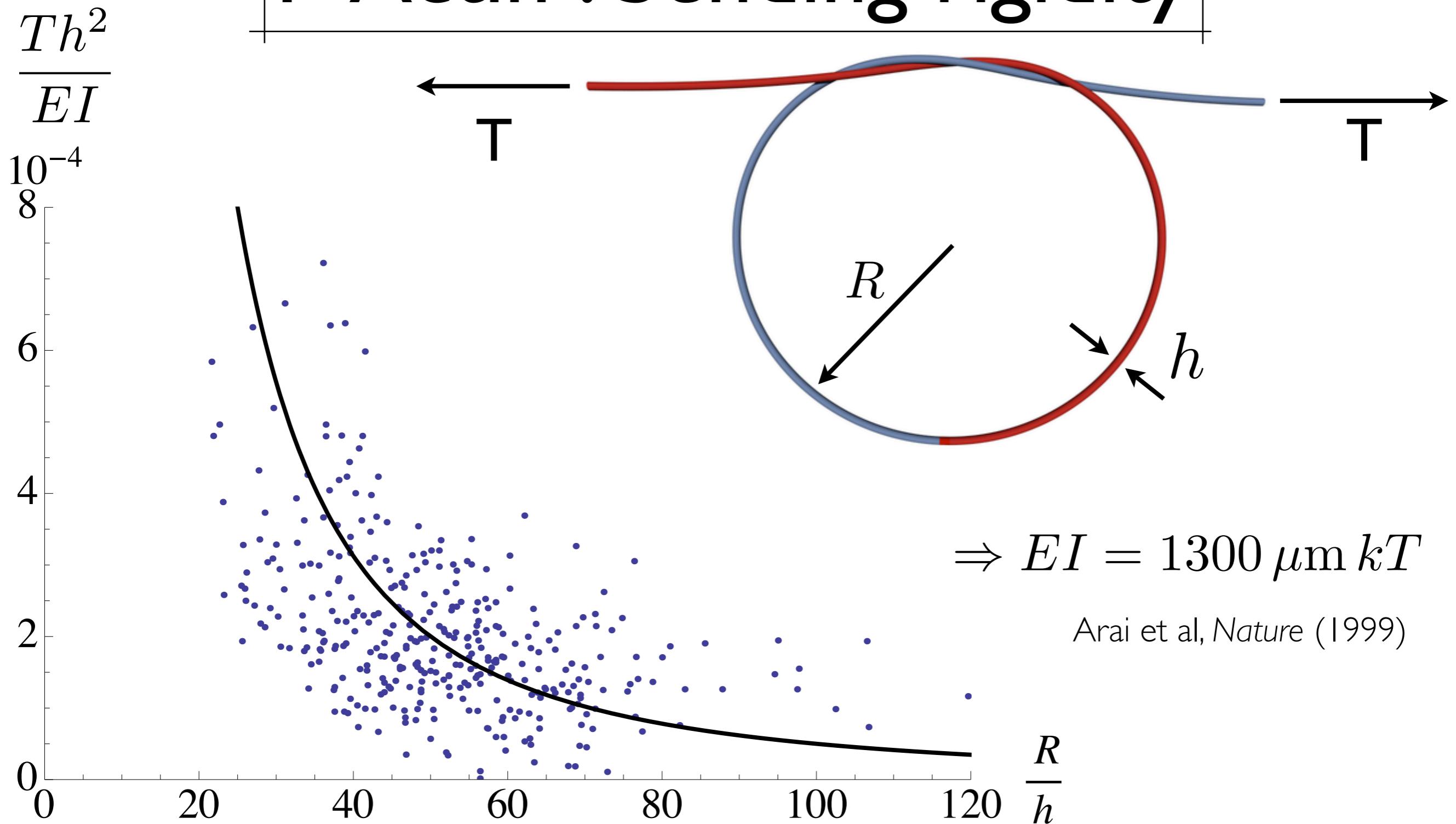
F-Actin

Cytoskelton filament Active network Locomotion

Single molecule experiment
with knotted F-Actin filaments
Arai et al, *Nature* (1999)



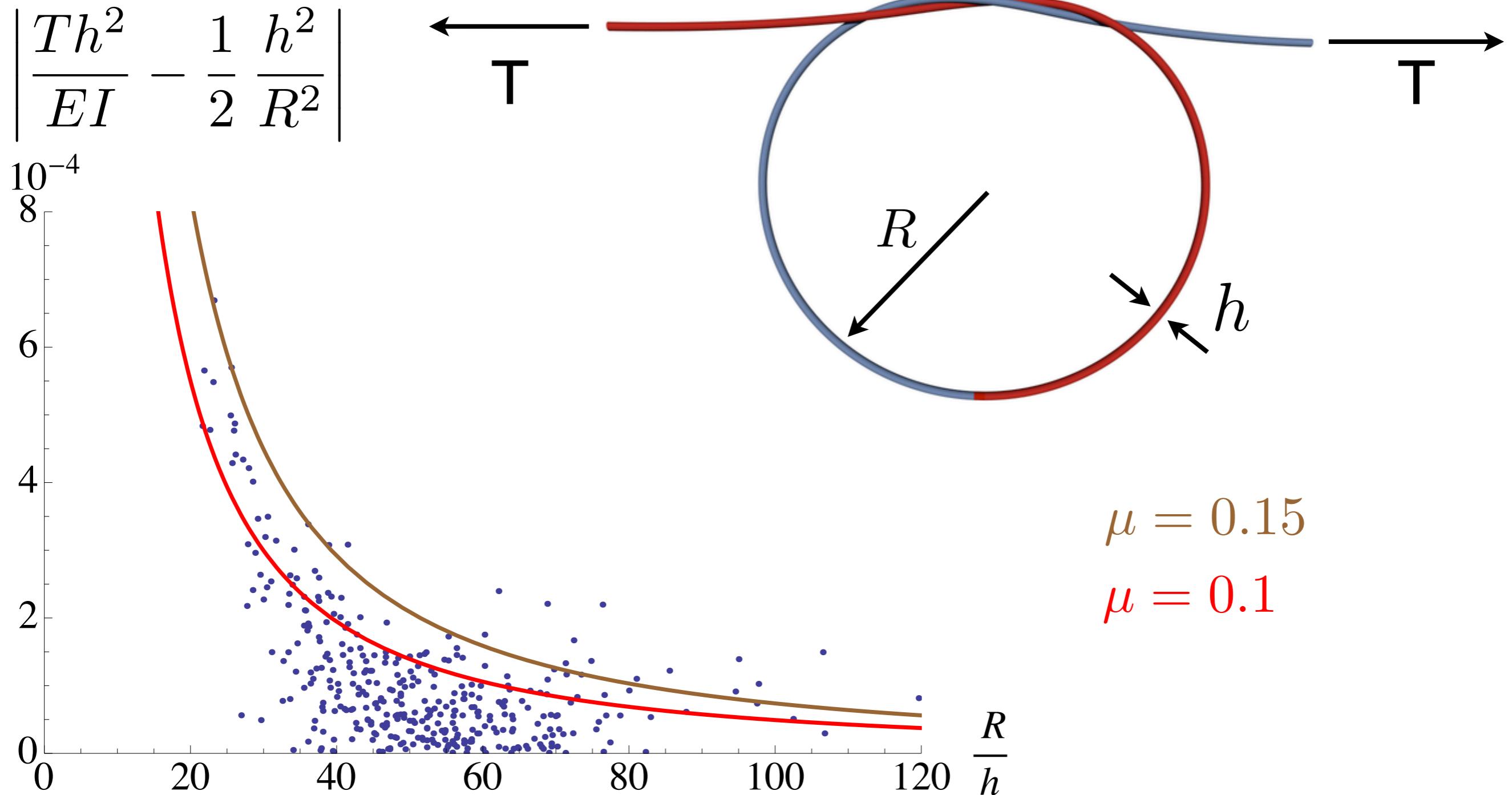
F-Actin : bending rigidity



sans frottement

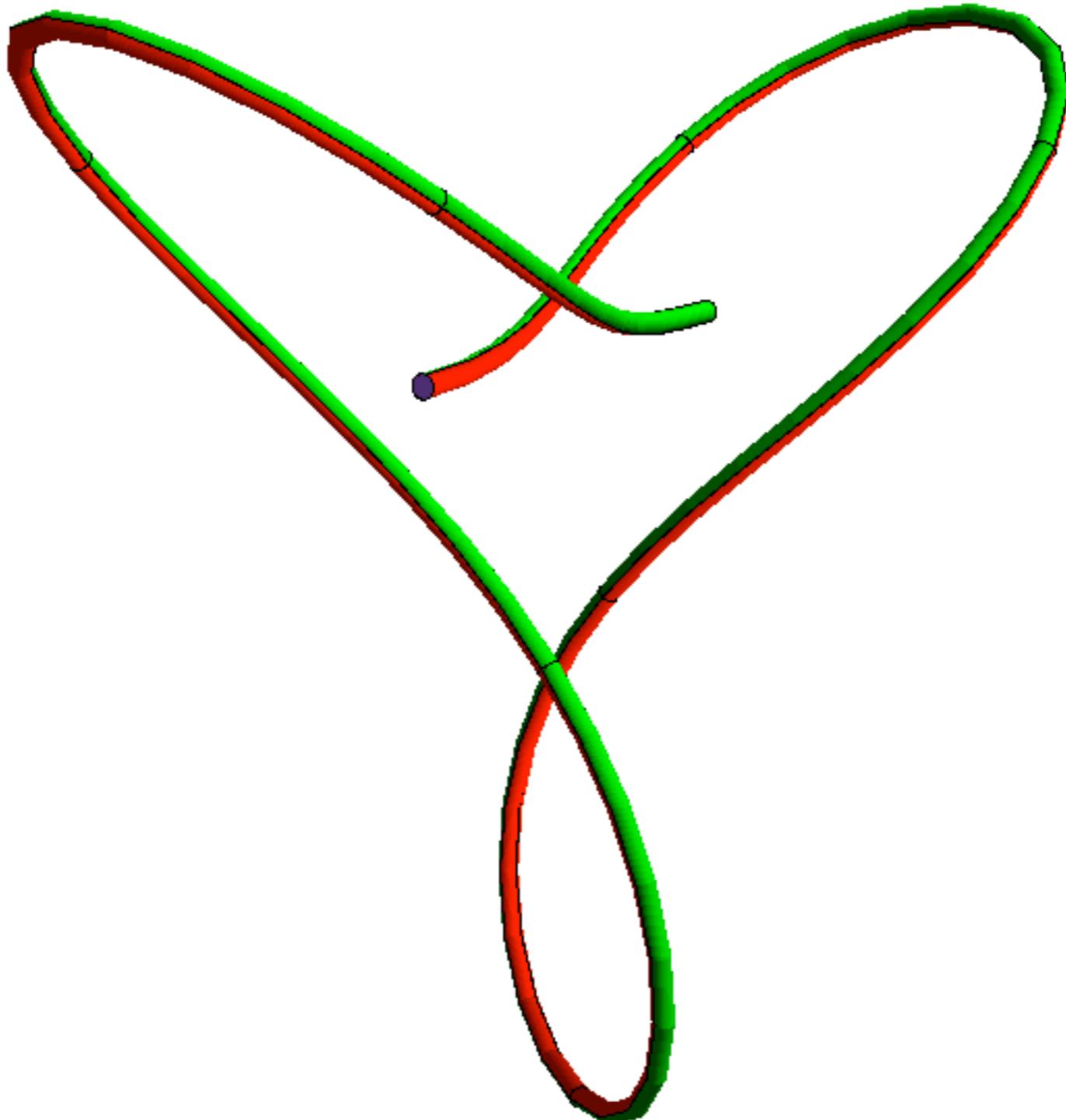
$$\frac{Th^2}{EI} = \frac{1}{2} \frac{h^2}{R^2}$$

F-Actin :self- friction coefficient



avec frottement

$$\left| \frac{Th^2}{EI} - \frac{1}{2} \frac{h^2}{R^2} \right| \leq \mu P = 0.49 \mu \left(\frac{h}{R} \right)^{3/2}$$



do stable open trefoil knotted configurations exist ?

Langer + & Singer (J. London Math. Soc) 1984 conjecture that no.
(for closed configurations though)

Fin

www.ida.upmc.fr/~neukirch

B. Audoly, N. Clauvelin, and S. Neukirch. *Physical Review Letters*, 99 (2007) 164301.

N. Clauvelin, B. Audoly, and S. Neukirch. *Journal of the Mechanics and Physics of Solids*, 57 (2009) 1623–1656.

H. O. Kirchner and S. Neukirch. *Journal of the Mechanical Behavior of Biomedical Materials*, 3 (2010) 121–123.

Variational formulation

$$E = \int_{-\infty}^{+\infty} \left(\frac{B}{2} \kappa^2 + \frac{C}{2} \tau^2 \right) ds + TD_\infty - UR_\infty,$$

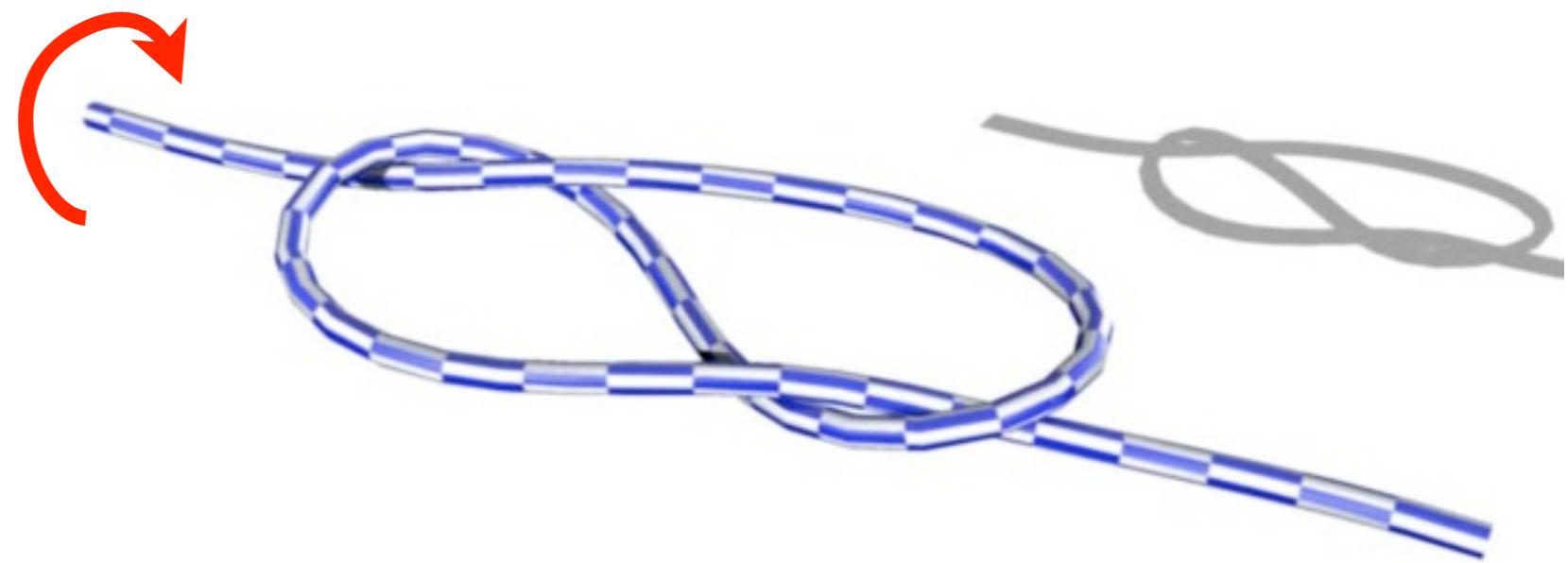
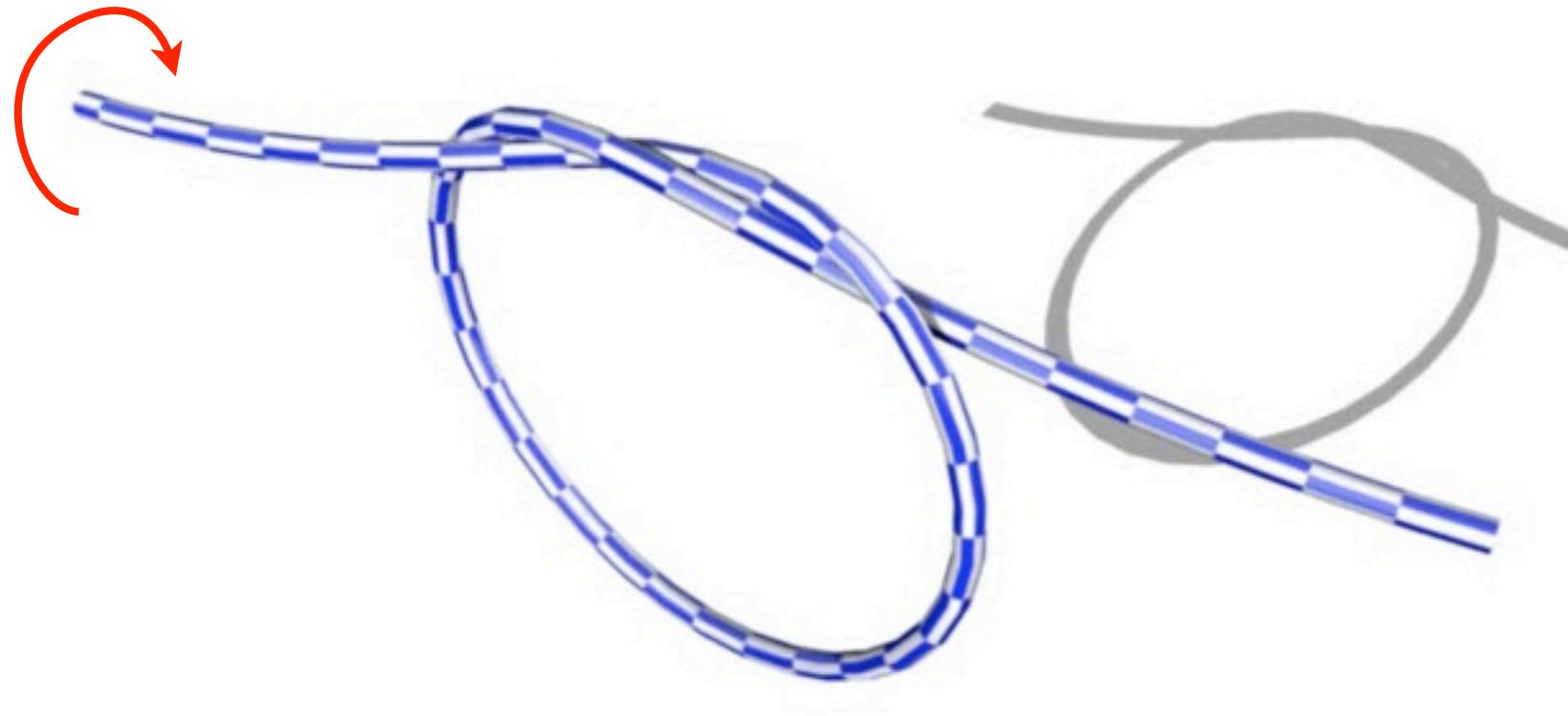
where κ and τ stand for the curvature
and the twist of the rod.

$$\kappa = |\mathbf{t}'(s)|.$$

$$|\mathbf{r}(s_1) - \mathbf{r}(s_2)| \geq 2h,$$

for any s_1 and s_2 such that $|s_1 - s_2| > 4h$.

Twist Instability

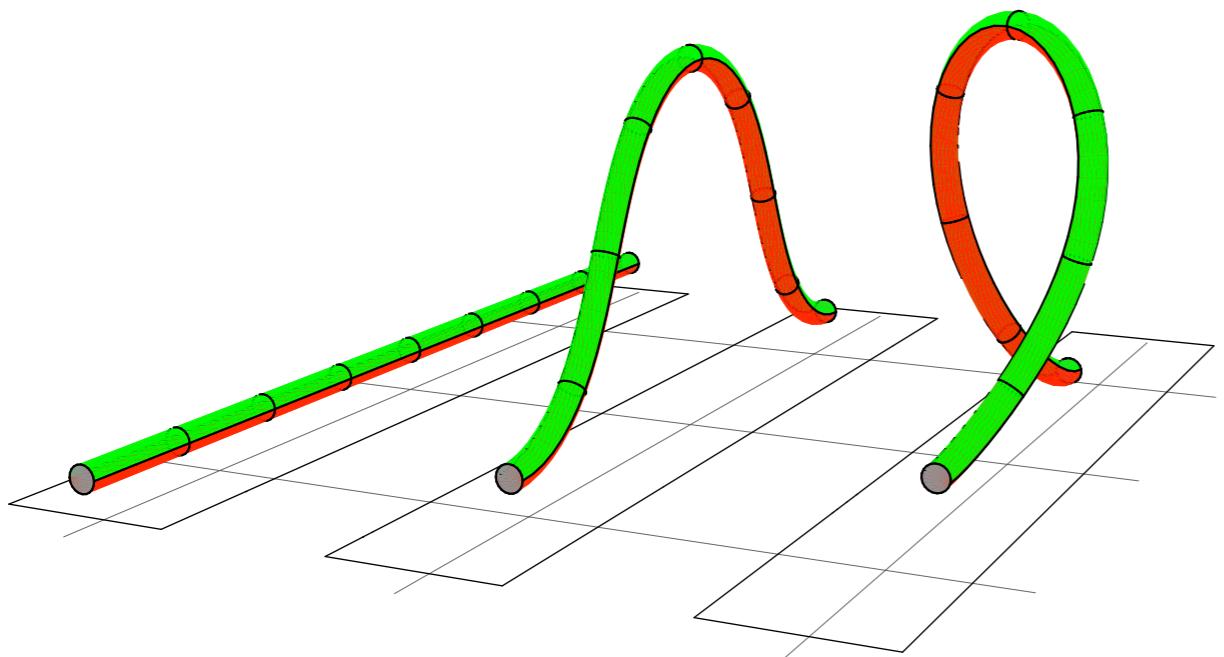


numerical simulations : M. Bergou, M. Wardetzky, S. Robinson, B. Audoly, and E. Grinspun.

ACM Transactions on Graphics (SIGGRAPH), 2008

Twisted rods : the ideal case

if rod is uniform, isotropic, naturally straight

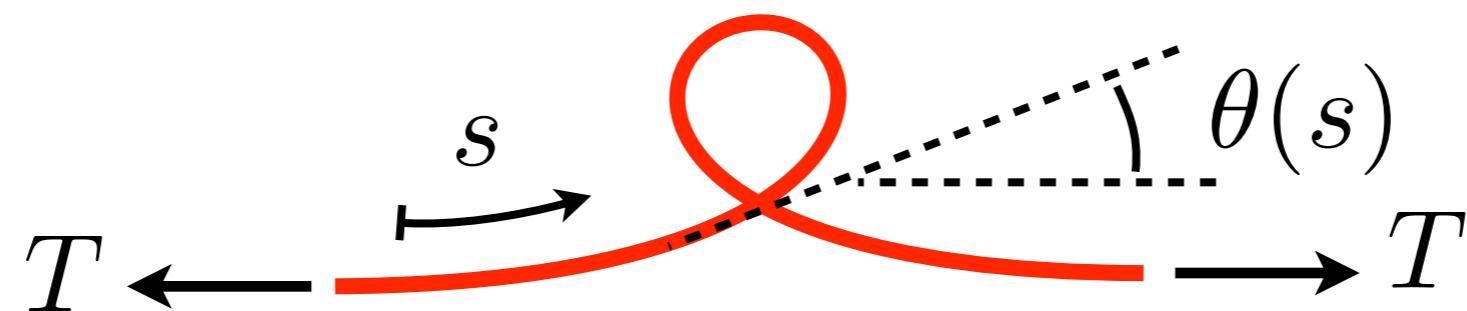


system reduction
 $2 \text{D} \Rightarrow 6\text{D}$

$$r' = d_3$$

$$d'_3 = (F \times r + M_0) \times d_3$$

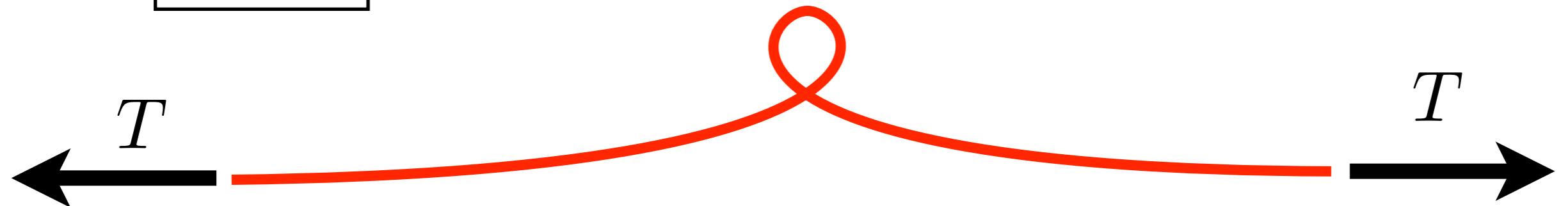
Planar Elastica



$$EI\theta'' = T \sin \theta$$

Planar Elastica

large T

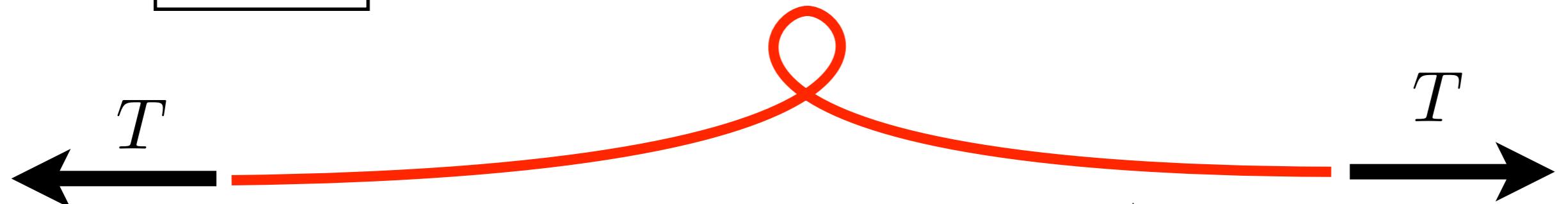


$$\frac{EI}{T} \theta'' = \sin \theta$$

singular
perturbation

Planar Elastica

large T

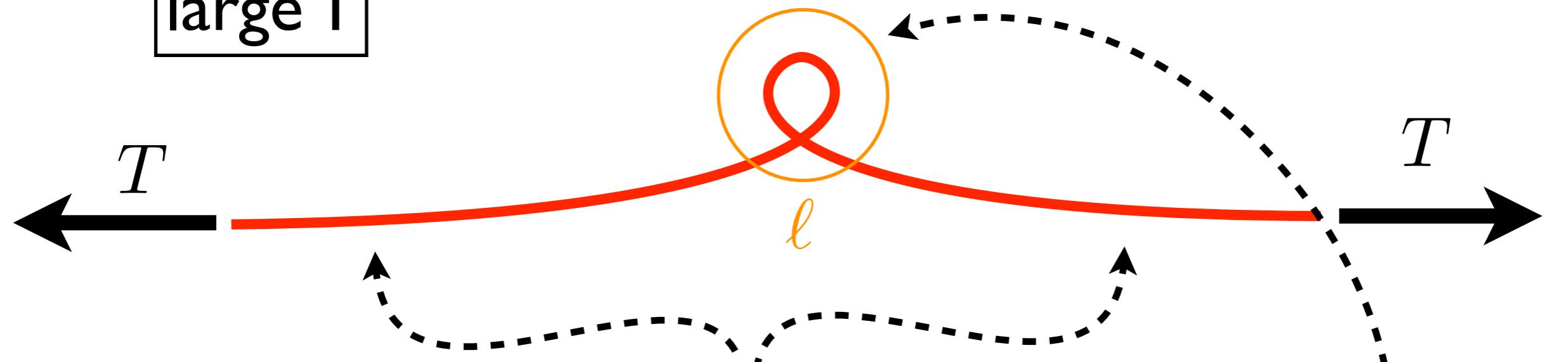


$$\frac{EI}{T} \theta'' = \sin \theta \quad \left\{ \begin{array}{l} \sin \theta \approx 0 \Rightarrow \theta(s) \approx 0 \end{array} \right.$$

singular
perturbation

Planar Elastica

large T



$$\frac{EI}{T} \theta'' = \sin \theta \quad \left\{ \begin{array}{l} \sin \theta \approx 0 \Rightarrow \theta(s) \approx 0 \\ \theta(s) \text{ rapidly varying} \end{array} \right.$$

region size : $\ell \sim \sqrt{\frac{EI}{T}}$

singular
perturbation

inner layer

Kirchhoff Equations

$$\left\{ \begin{array}{l} \vec{F}' = -\vec{p} \\ \vec{M}' = \vec{F} \times \vec{t} \\ \vec{t}' = \frac{1}{EI} \vec{M} \times \vec{t} \\ \vec{R}' = \vec{t} \end{array} \right. \quad \begin{array}{l} \text{forces equil.} \\ \text{moments equil.} \\ \text{kinematics} \\ \text{tangent def.} \end{array}$$

$$' \equiv \frac{d}{ds}$$

constitutive relations:

$$\begin{array}{lll} M_\kappa & = & EI \kappa \quad \text{curvature } \kappa \\ M_\tau & = & GJ \tau \quad \text{twist } \tau \end{array}$$

$\vec{p}(s)$ ext. pressure

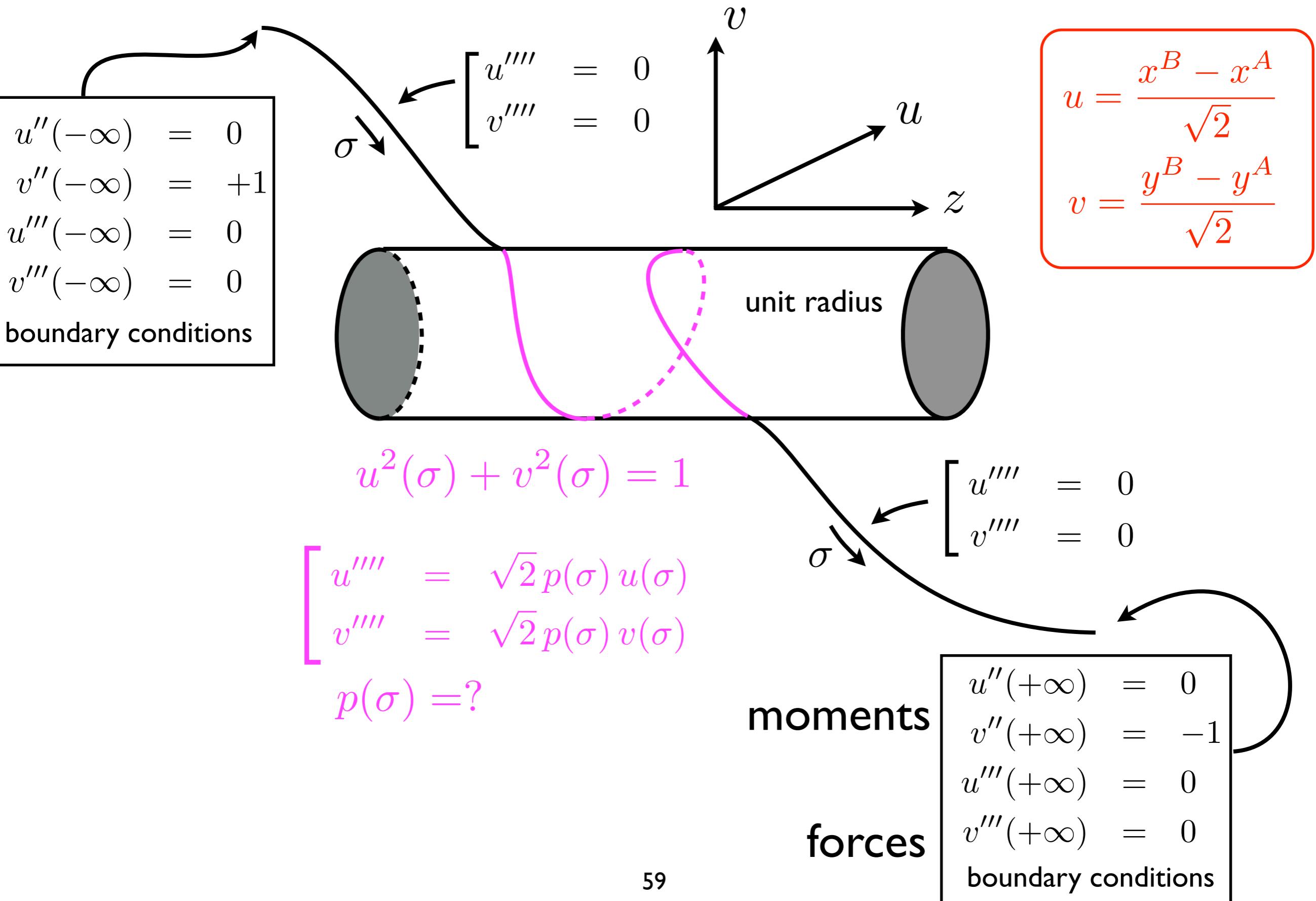
$\vec{M}(s)$ internal moment

$\vec{F}(s)$ internal force

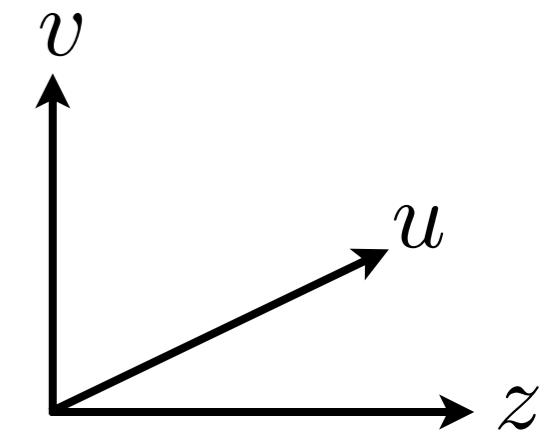
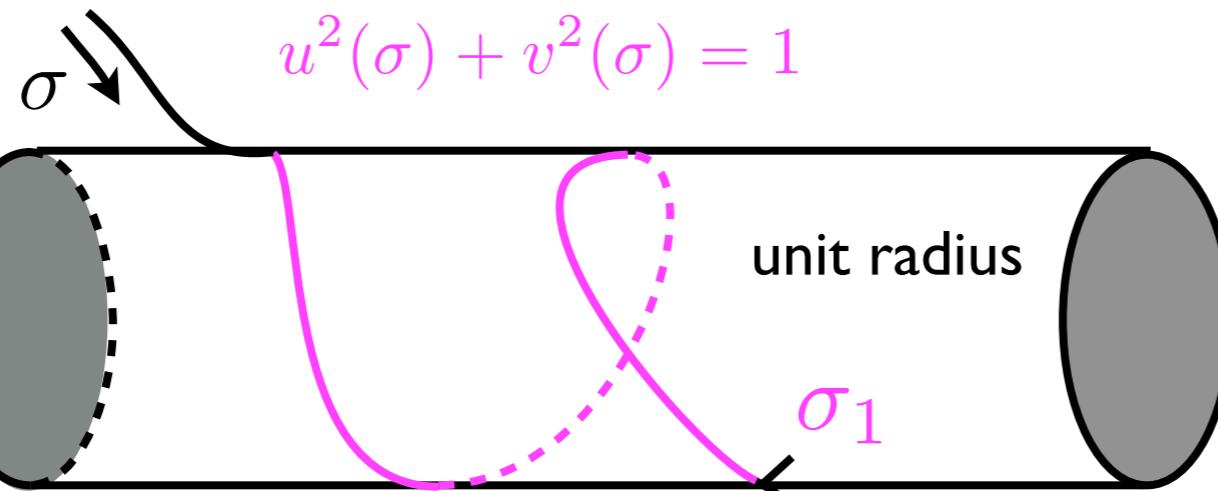
$\vec{R}(s)$ position

$\vec{t}(s)$ tangent

Braid : boundary value problem (BVP)



Braid : first kind of solutions



boundary conditions

$$u''(\sigma_1) = 0$$

$$v''(\sigma_1) = -1$$

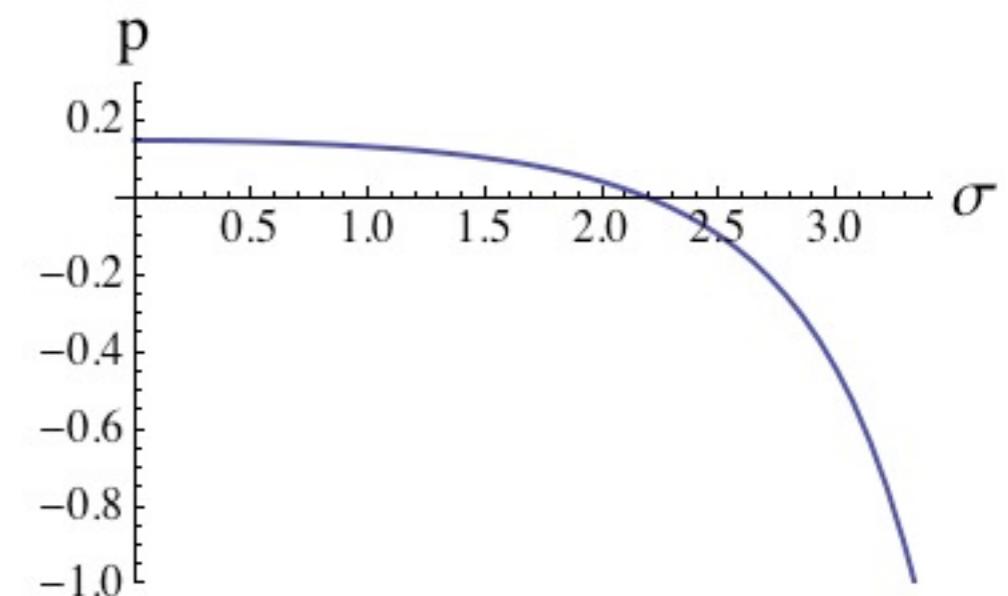
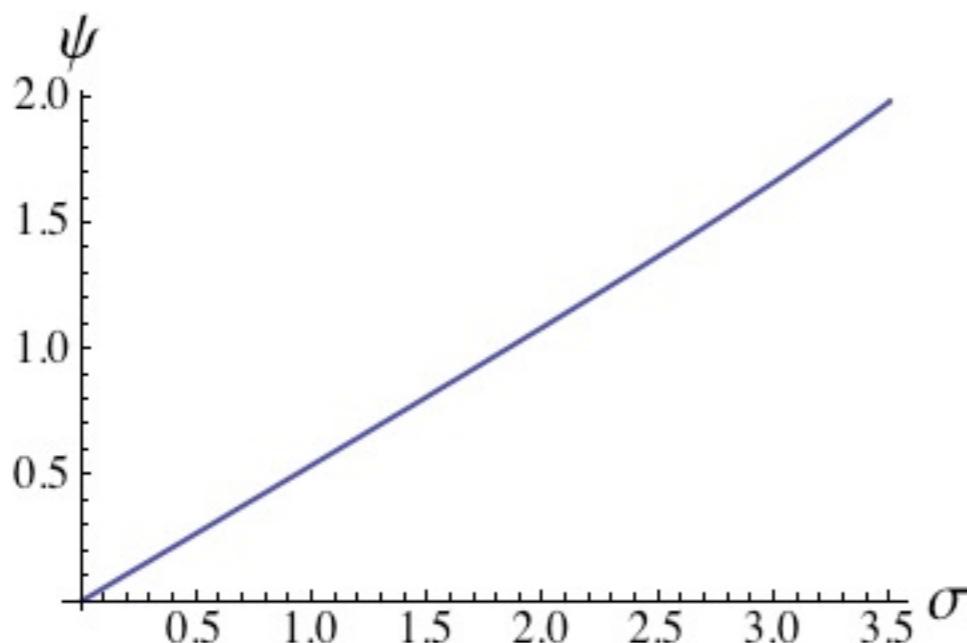
$$u'''(\sigma_1) = 0$$

$$v'''(\sigma_1) = 0$$

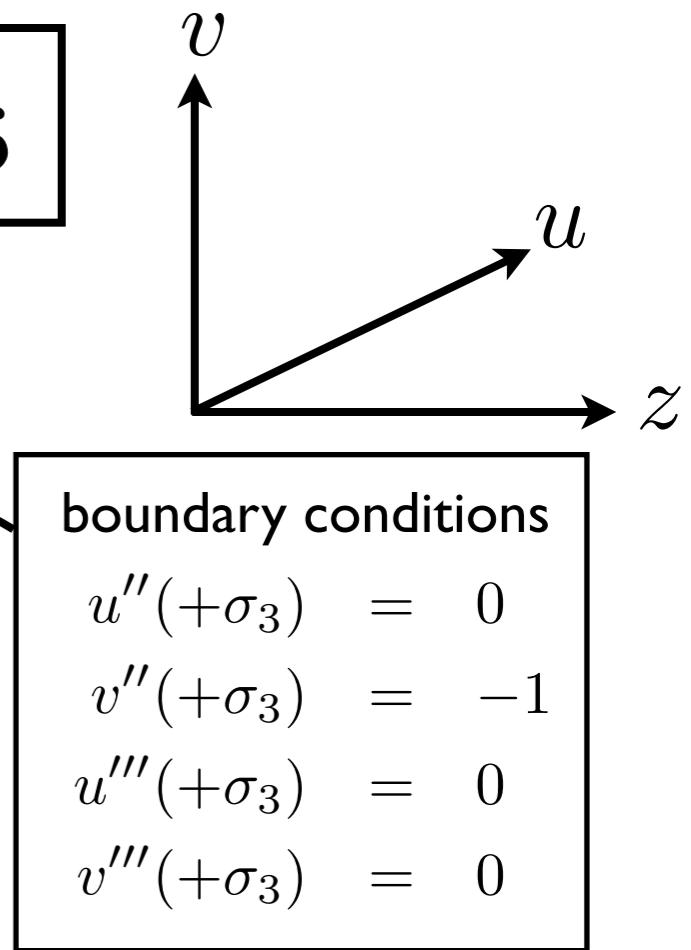
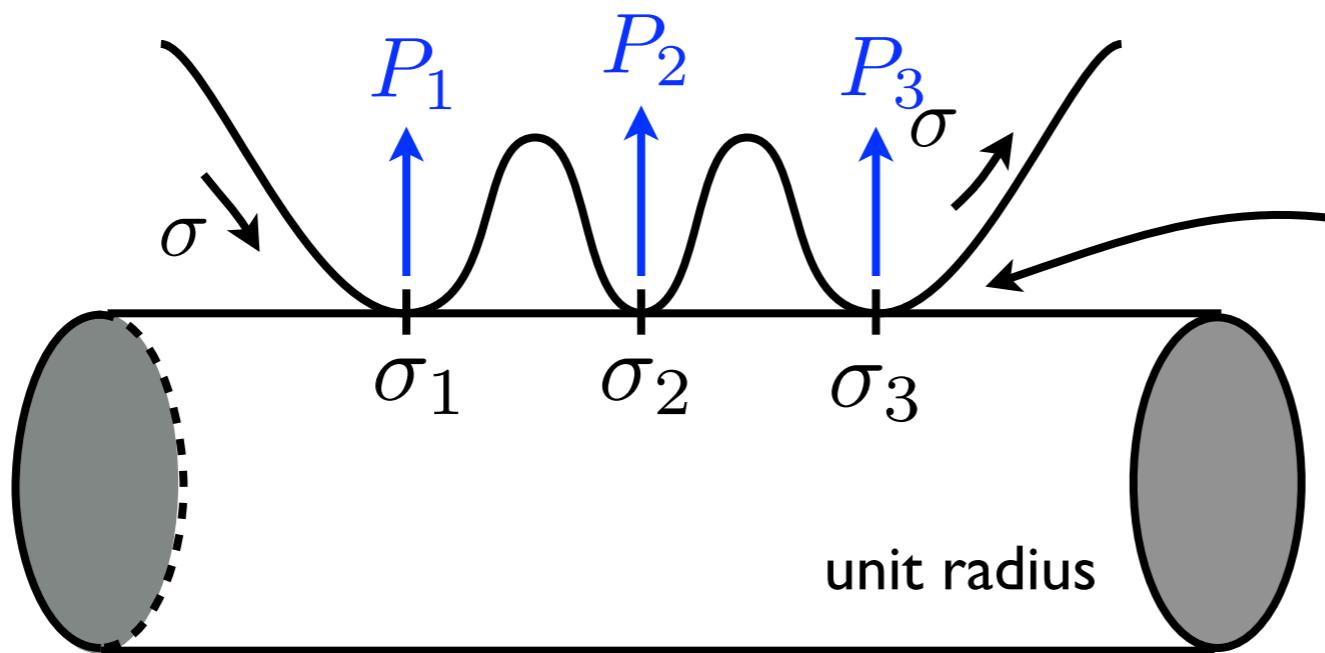
$$\begin{cases} u'''' = \sqrt{2} p(\sigma) u(\sigma) \\ v'''' = \sqrt{2} p(\sigma) v(\sigma) \end{cases} \quad \begin{cases} u = \cos(\psi(\sigma)) \\ v = \sin(\psi(\sigma)) \end{cases}$$

$$\begin{cases} \psi'''' = 6 \psi'' \psi'^2 \\ p(\sigma) = (\psi'^4 - 3\psi''^2 - 4\psi' \psi''') / \sqrt{2} \end{cases}$$

$\psi(0)$	=	0
$\psi'(0)$	=	0.54
$\psi''(0)$	=	0
$\psi'''(0)$	=	0.004
σ_1	=	3.50



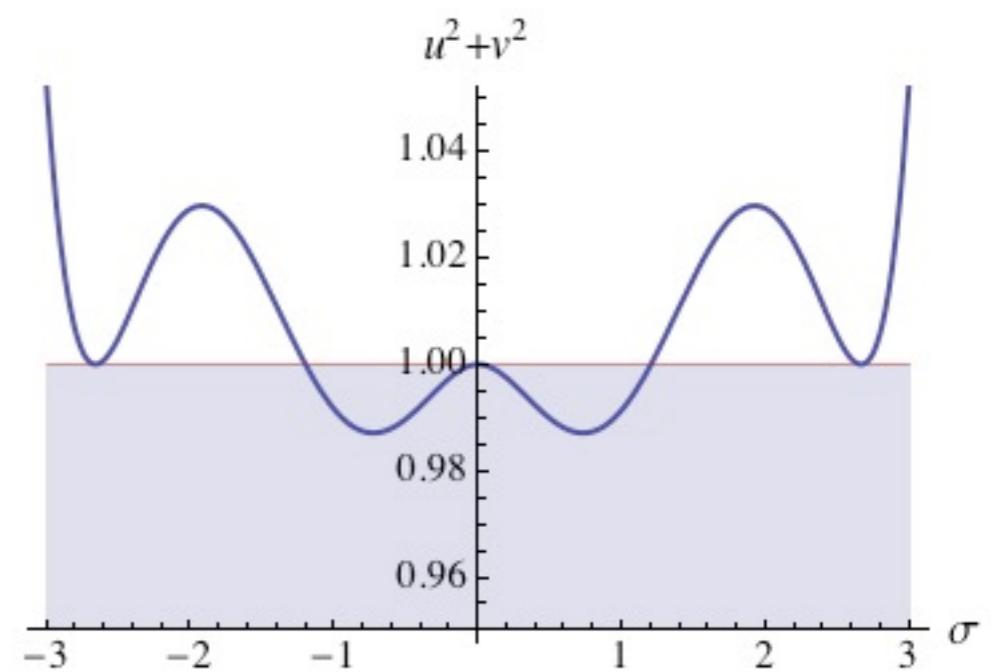
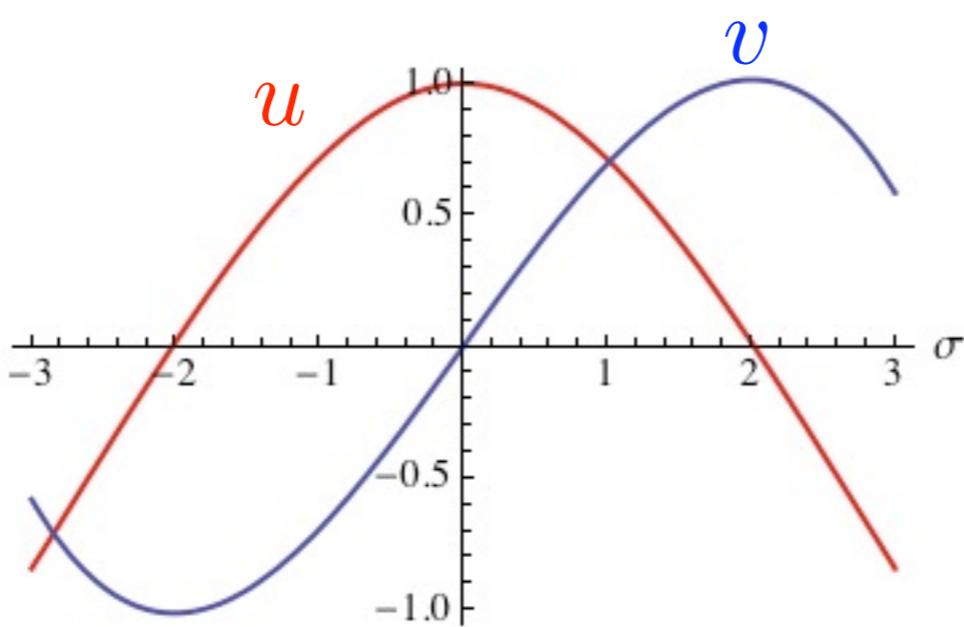
Braid : second kind of solutions



$$\begin{cases} u'''' = \sqrt{2} p(\sigma) u(\sigma) \\ v'''' = \sqrt{2} p(\sigma) v(\sigma) \end{cases} \quad \text{avec} \quad p(\sigma) = P_1 \delta(\sigma - \sigma_1) + P_2 \delta(\sigma - \sigma_2) + P_3 \delta(\sigma - \sigma_3)$$

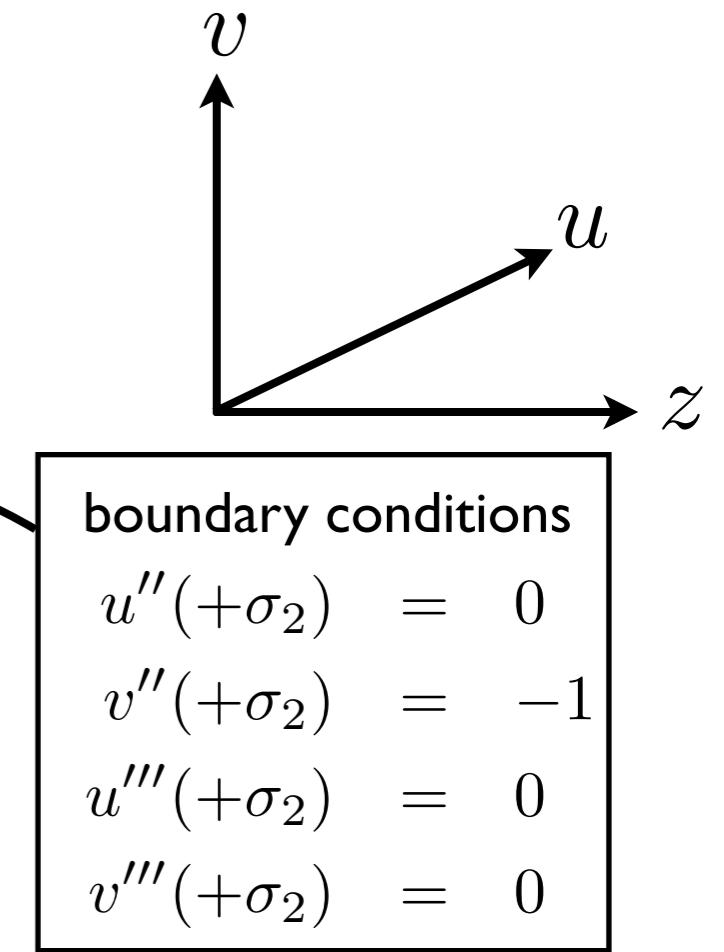
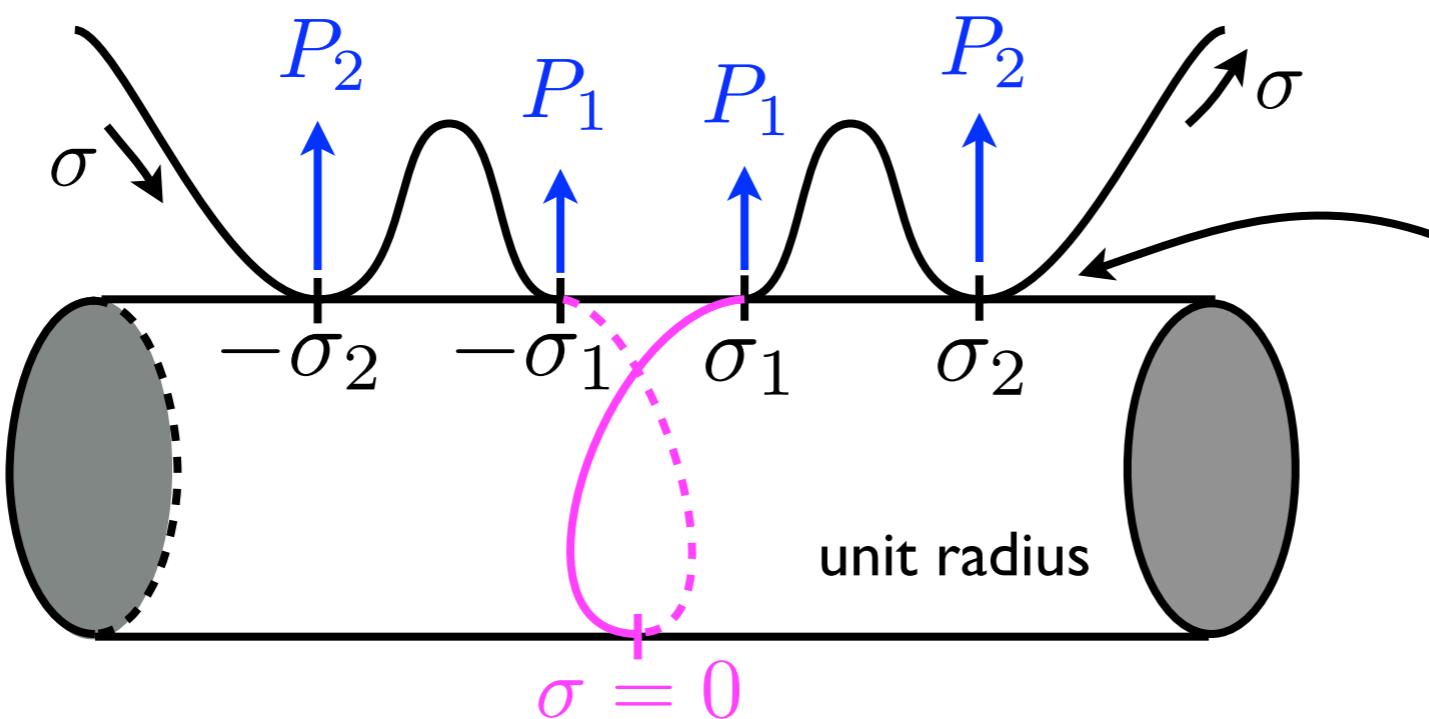
$$\begin{aligned} u(0) &= 1 \\ u'(0) &= 0 \\ u''(0) &= -0.66 \\ u'''(0) &= 0.25 \end{aligned}$$

$$\begin{aligned} v(0) &= 0 \\ v'(0) &= 0.76 \\ v''(0) &= 0 \\ v'''(0) &= -0.38 \end{aligned}$$



$$\sigma_3 = -\sigma_1 = 2.66 ; \sigma_2 = 0 ; P_1 = P_3 = 0.32 ; P_2 = 0.35$$

Braid : third kind of solutions



$$\begin{aligned}\sigma_1 &= 0.35 \\ \sigma_2 &= 2.68 \\ P_1 &= 0.12 \\ P_2 &= 0.31\end{aligned}$$

