

Slender beams vibrations: Frequency jumps at buckling

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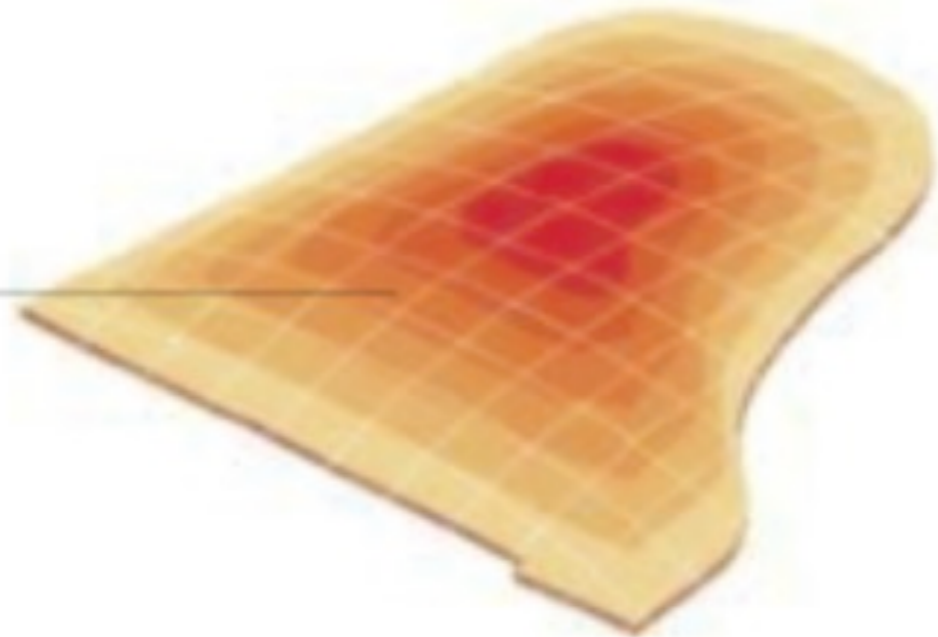
Centre for Applied Math. (OCCAM), Oxford, U.K.

Piano soundboard



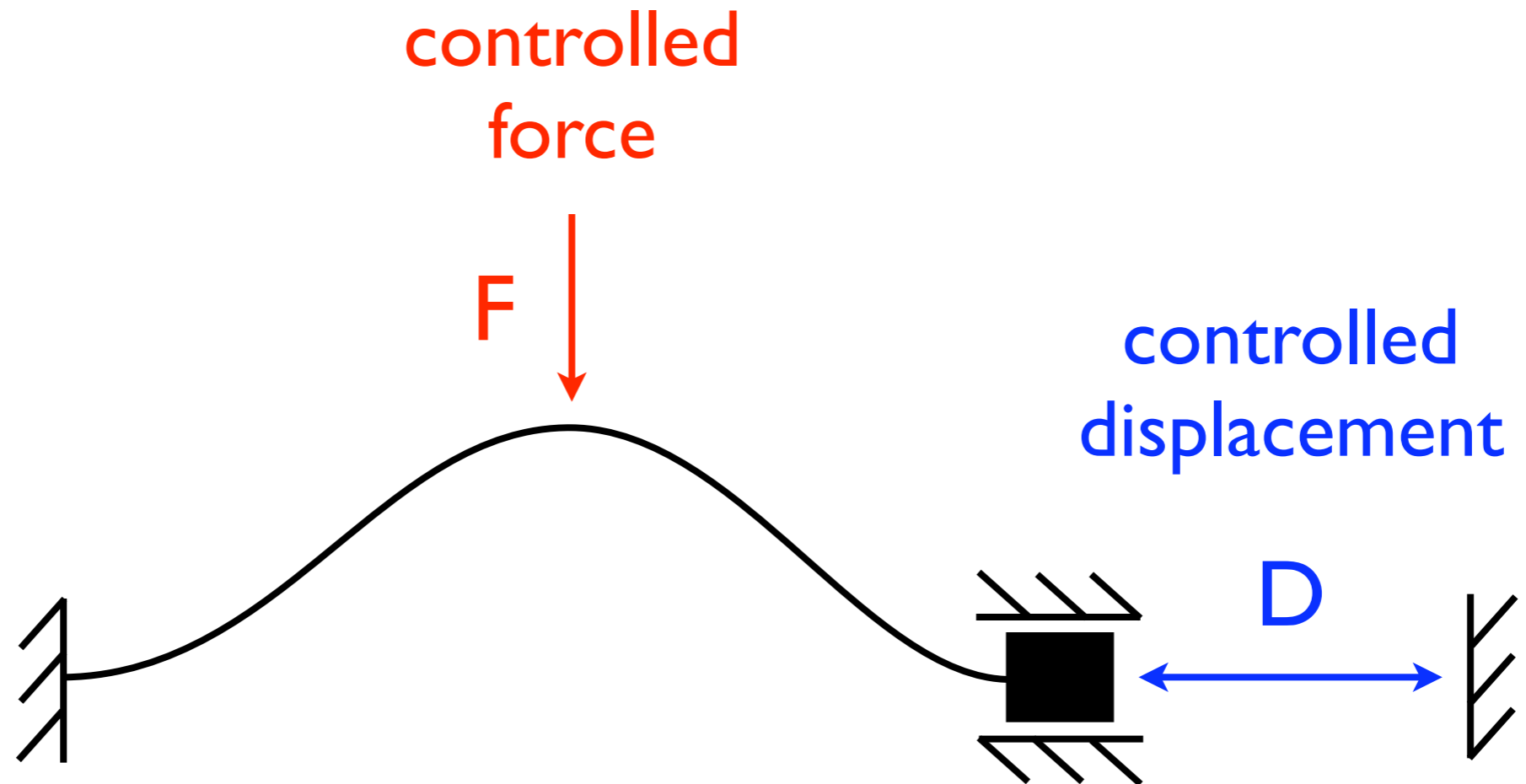
Piano soundboard

acoustic radiation from the soundboard (not the strings)



Model: pre-stressed beam

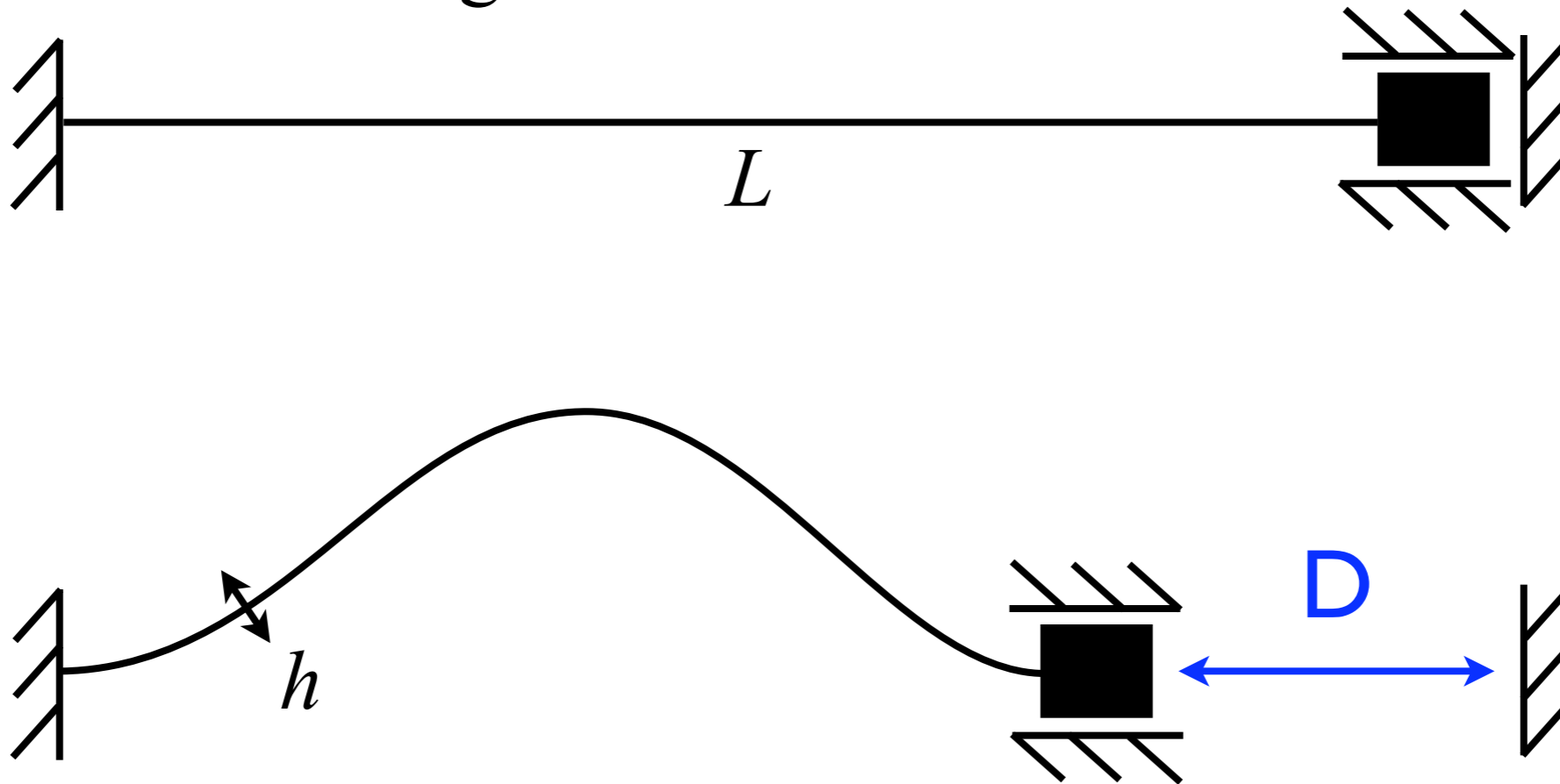
vibrations slender elastic beam in the plane



Influence of F , D
on the frequencies ?

Elastic beam in the plane

L : length in unstressed state



h : section thickness

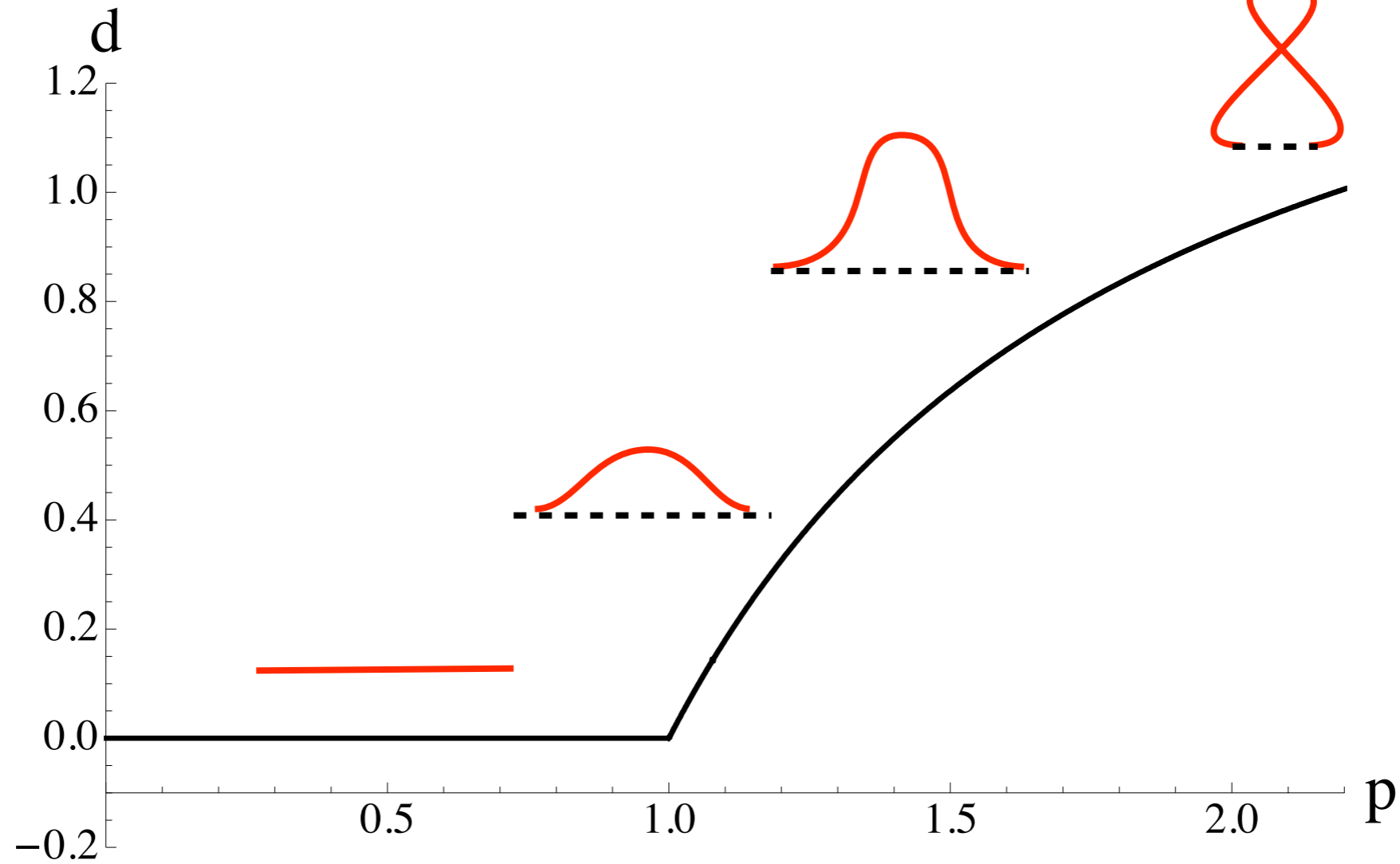
e : section width

$$I = \frac{1}{12} h^3 e$$

$$A = h e$$

Equilibrium (numerical study)

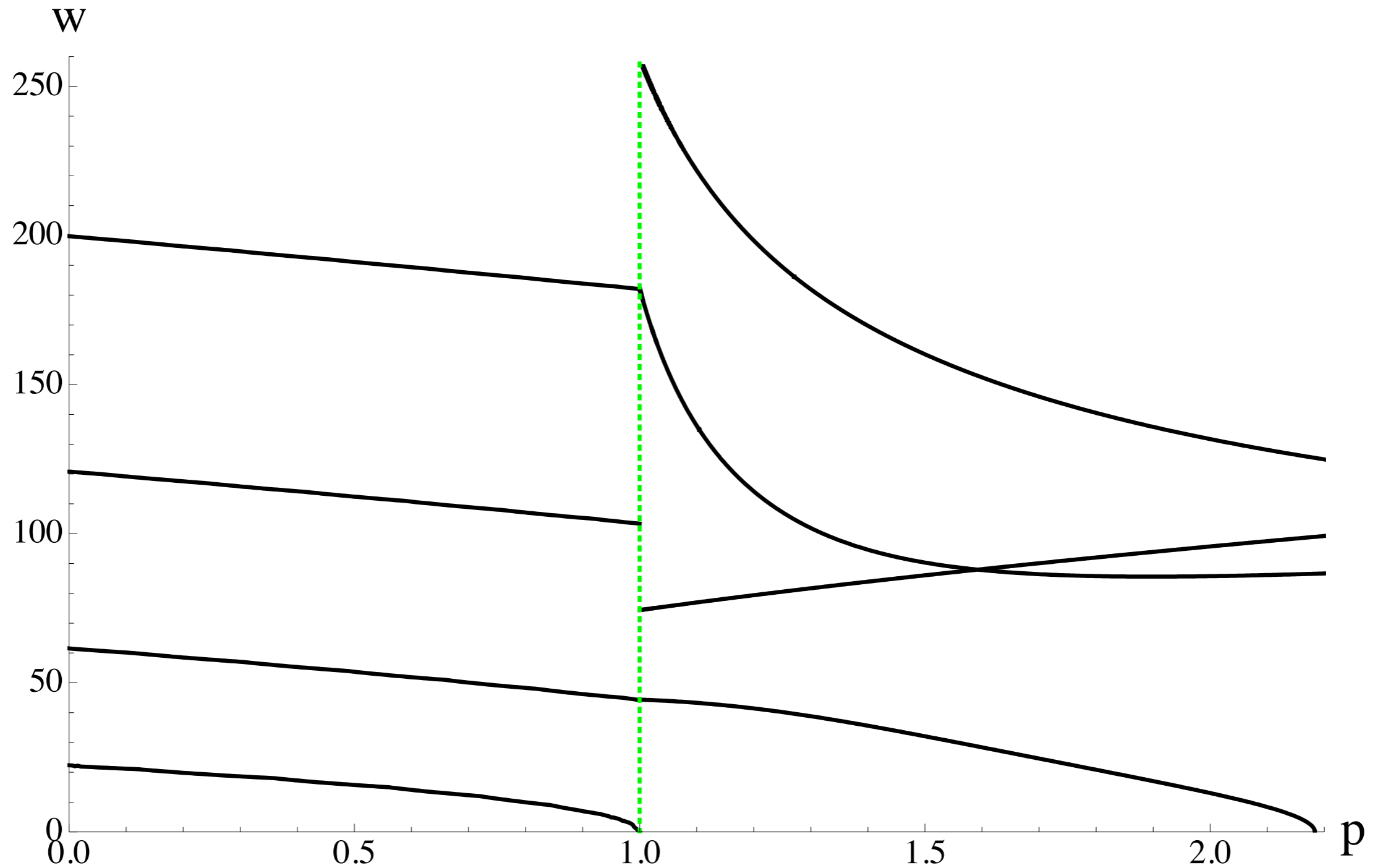
$$d = \frac{D}{L}$$



$$p = \frac{PL^2}{4\pi^2 EI}$$

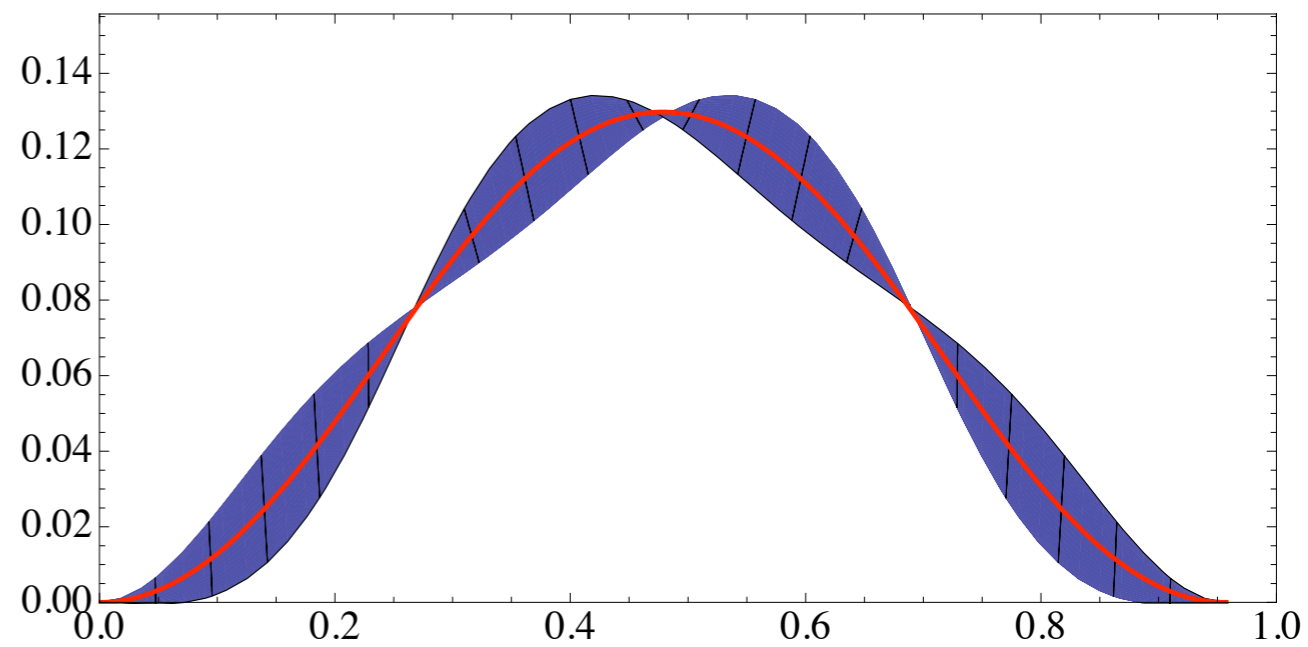
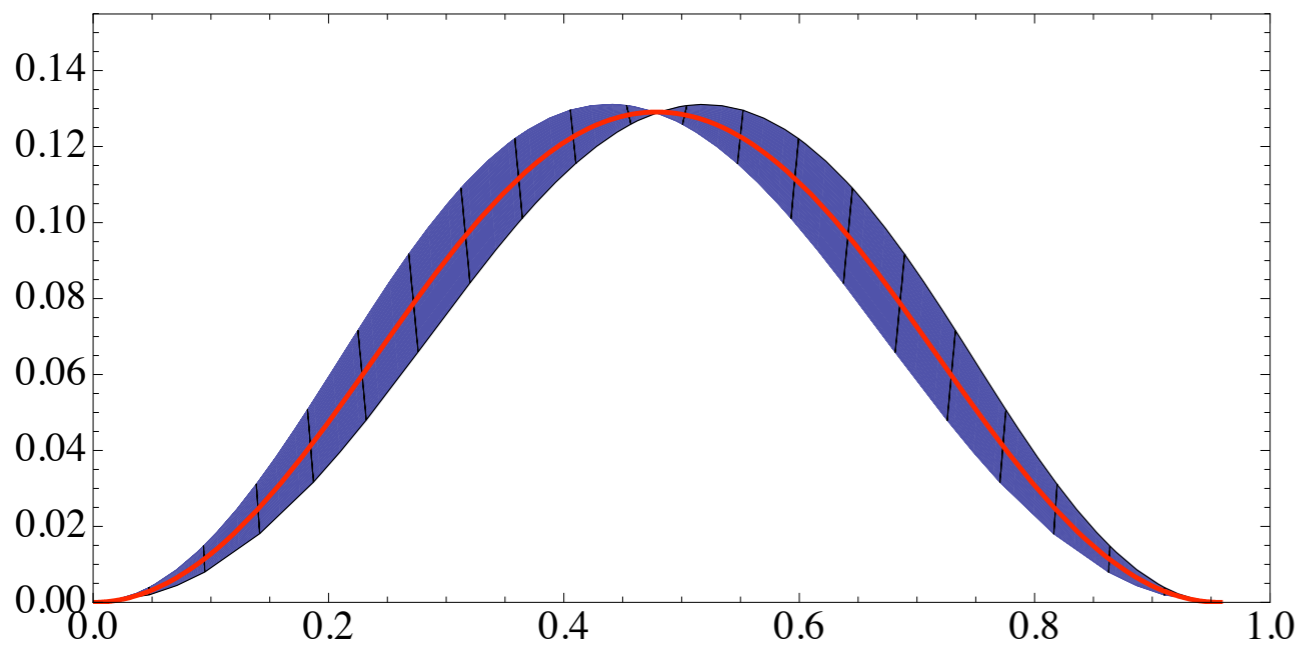
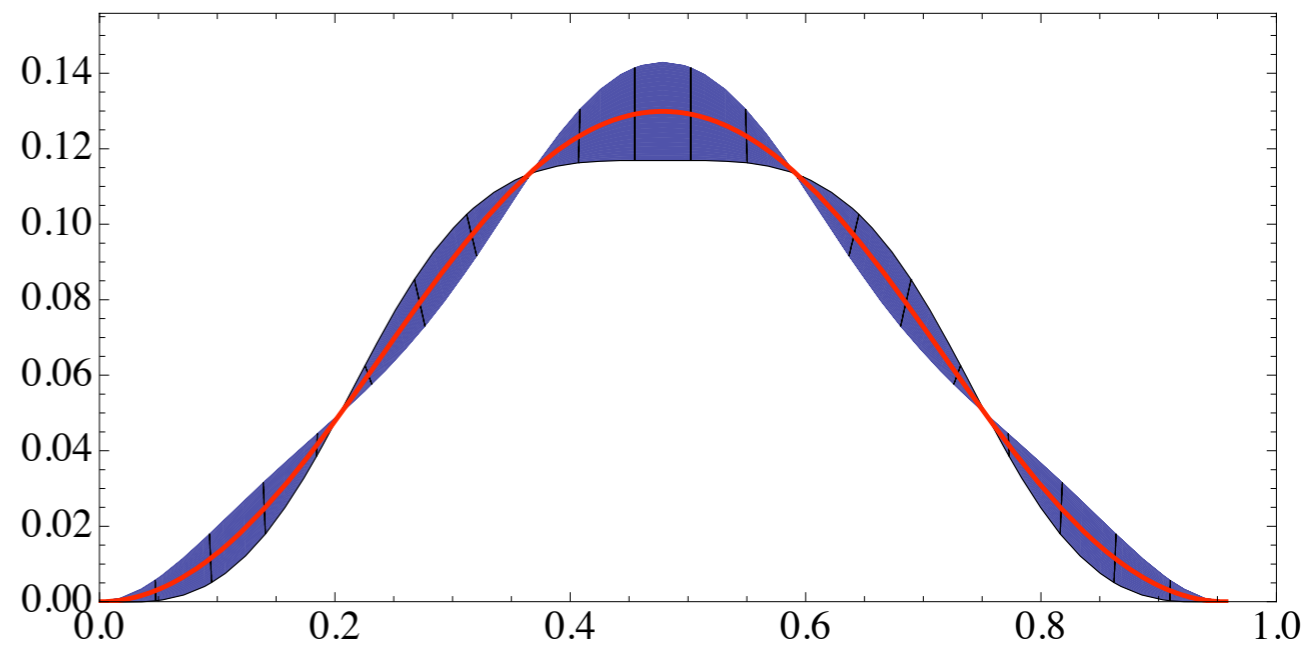
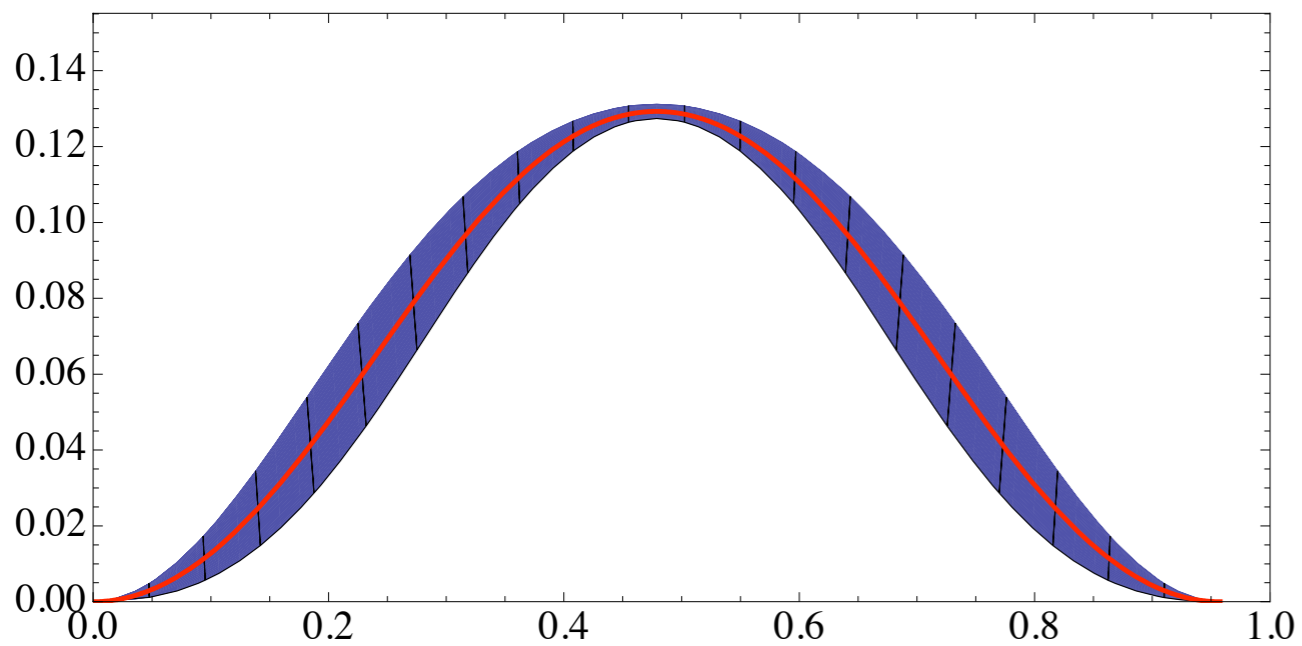
$$\omega = \Omega L^2 \left(\frac{\rho A}{EI} \right)^{1/2}$$

Vibrations (numerical study)



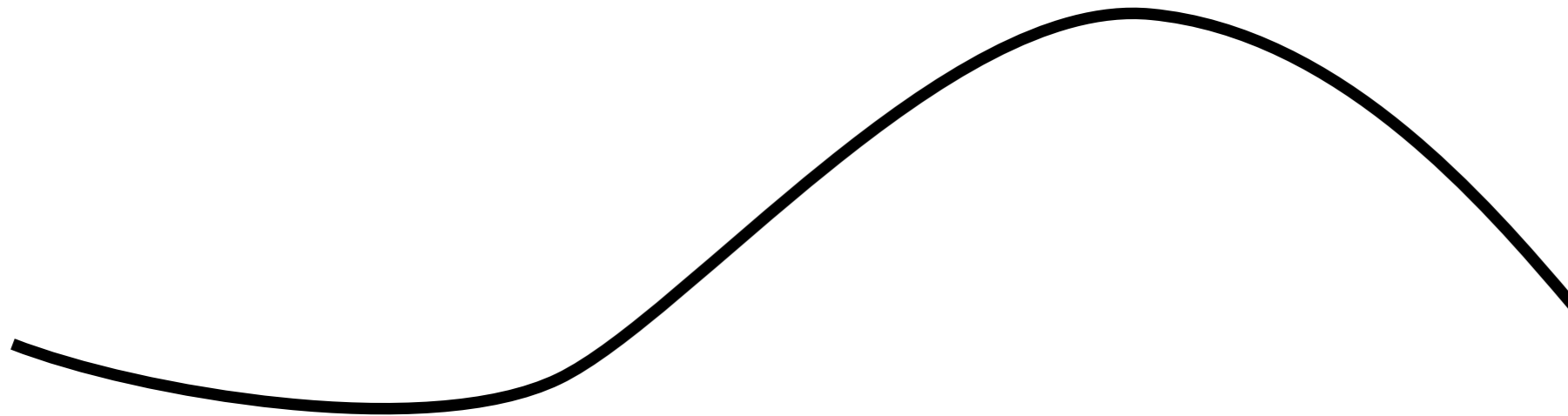
$$p = \frac{PL^2}{4\pi^2 EI}$$

Vibrations: first 4 modes



Kirchhoff equations

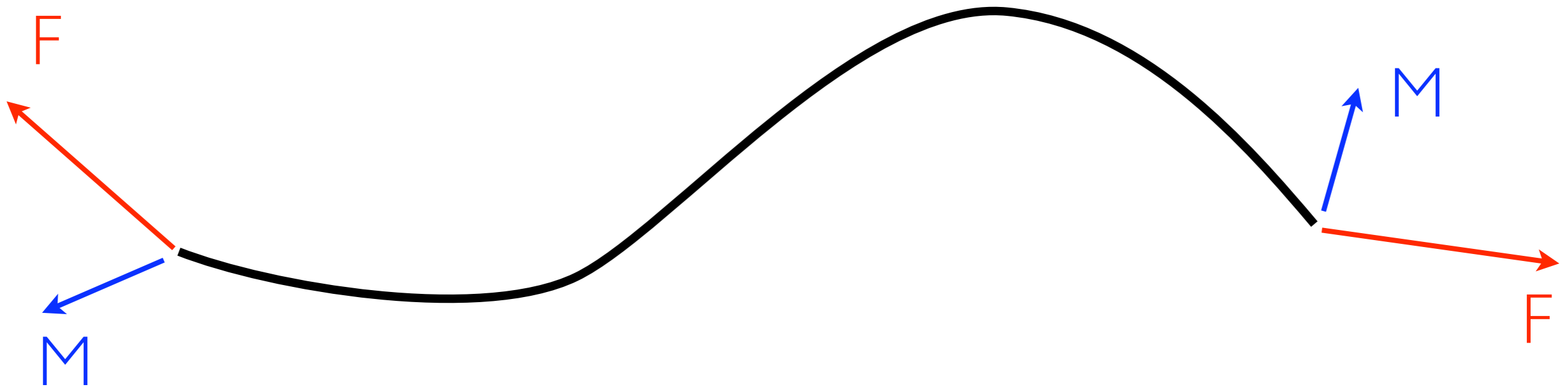
short
review



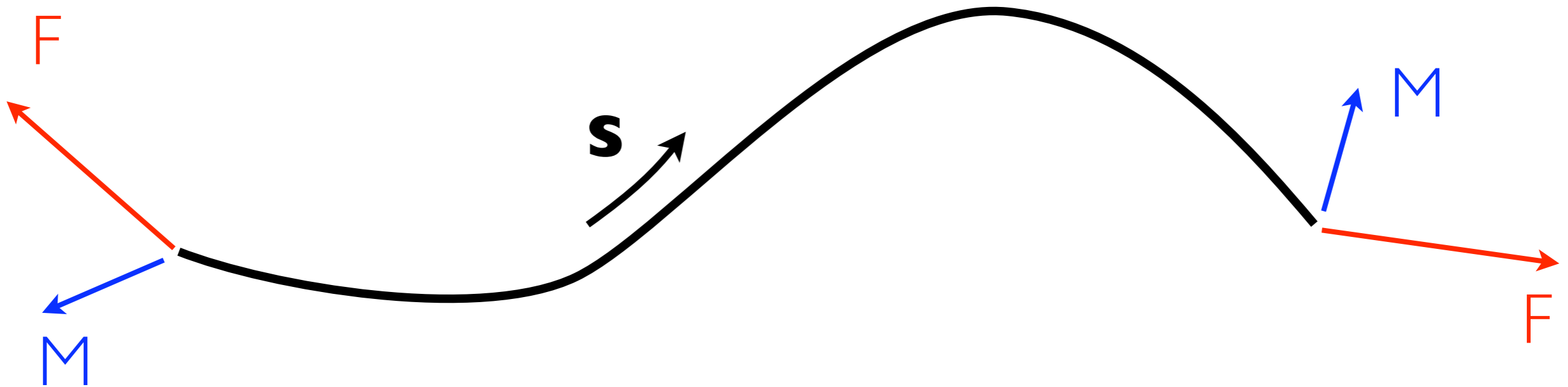
apply to :

- slender bodies
- not too bent

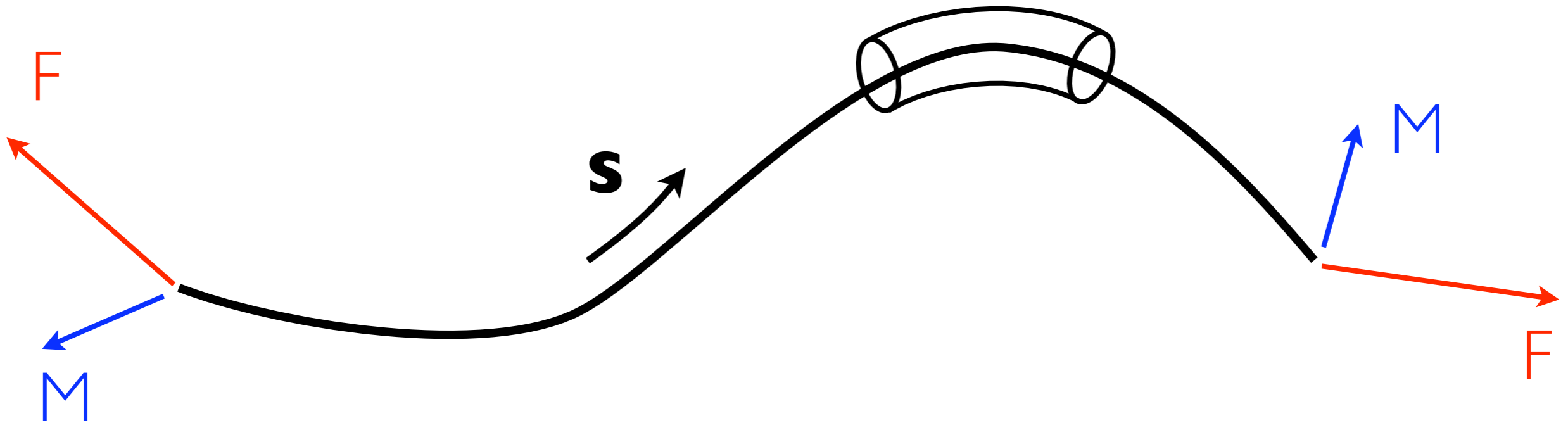
Kirchhoff equations



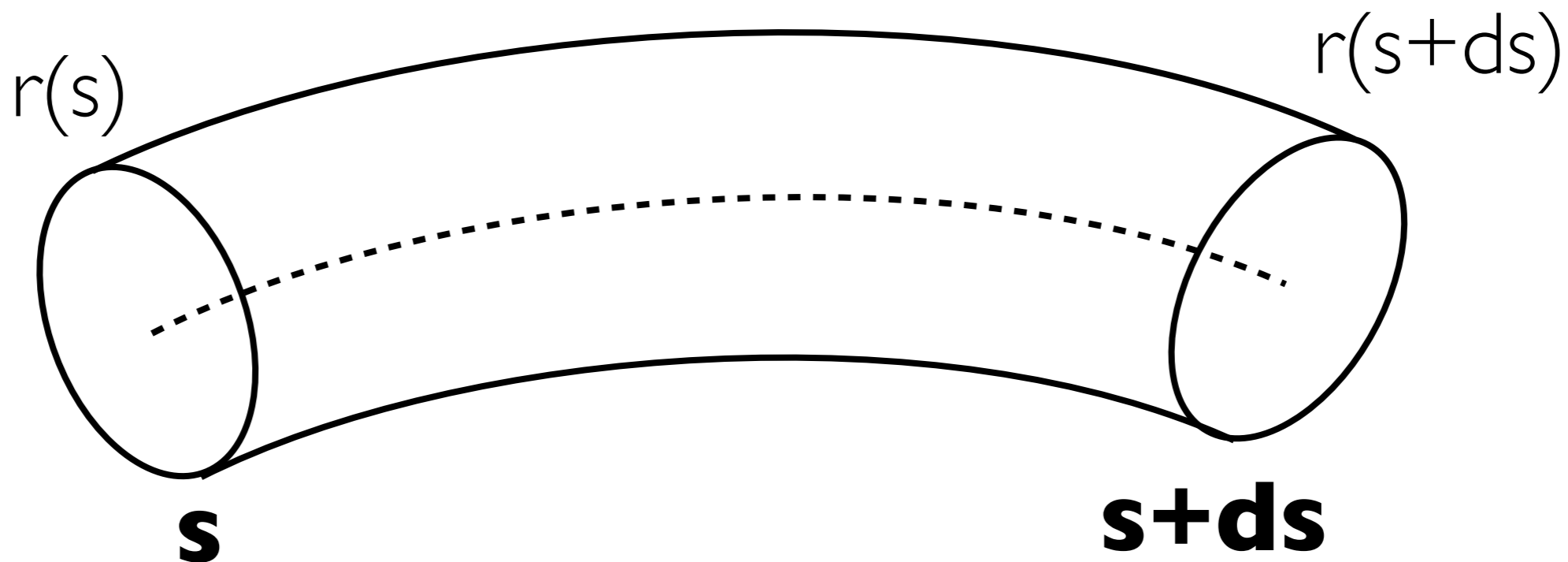
Kirchhoff equations



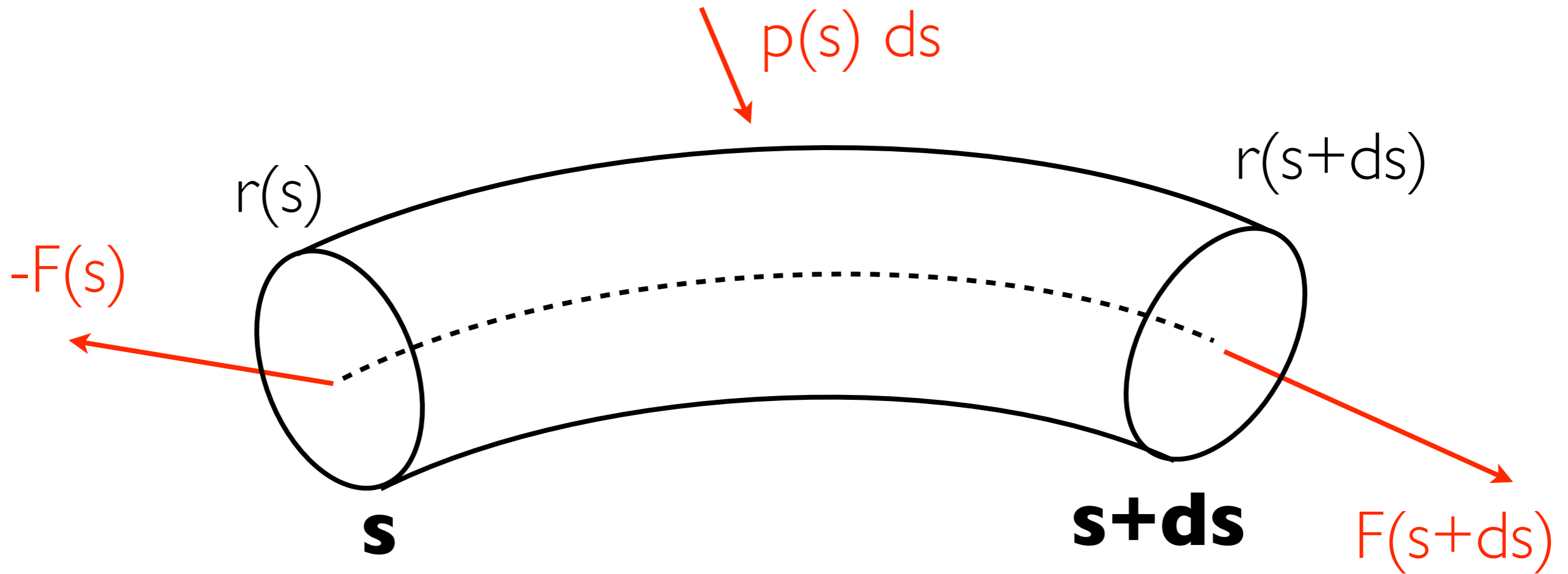
Kirchhoff equations



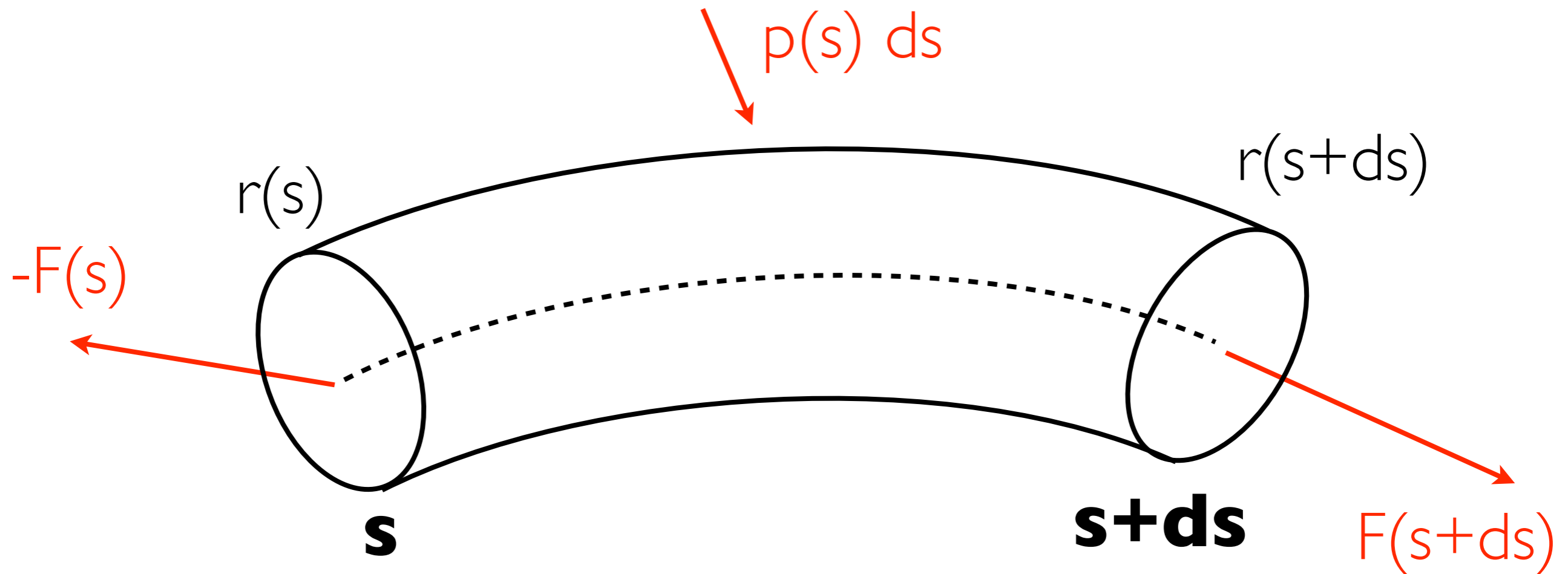
Kirchhoff equations



Kirchhoff equations



Kirchhoff equations



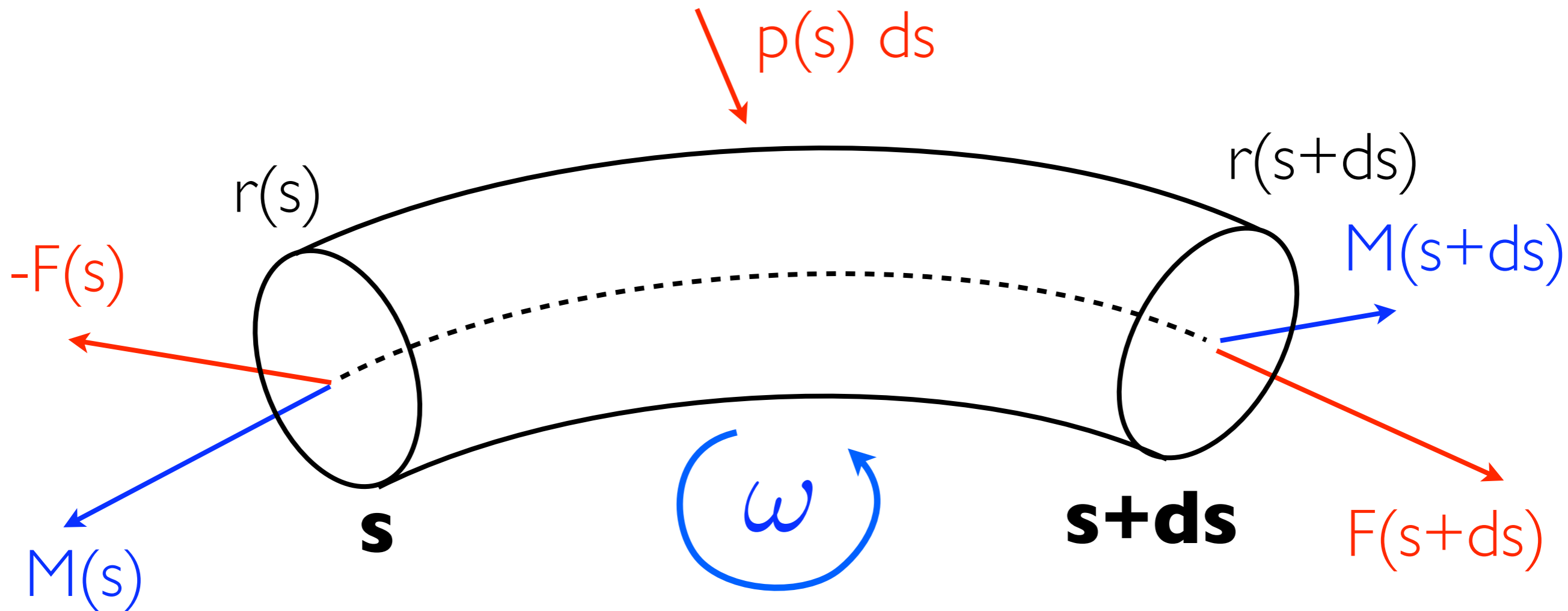
Dynamics (linear momentum):

$$F(s+ds, t) - F(s, t) + p(s, t) ds = \rho A ds \ddot{r}(s, t)$$

$$F'(s, t) + p(s, t) = \rho A \ddot{r}(s, t)$$

$$\dot{\quad} \equiv \frac{d}{dt} \quad , \quad ' \equiv \frac{d}{ds}$$

Kirchhoff equations

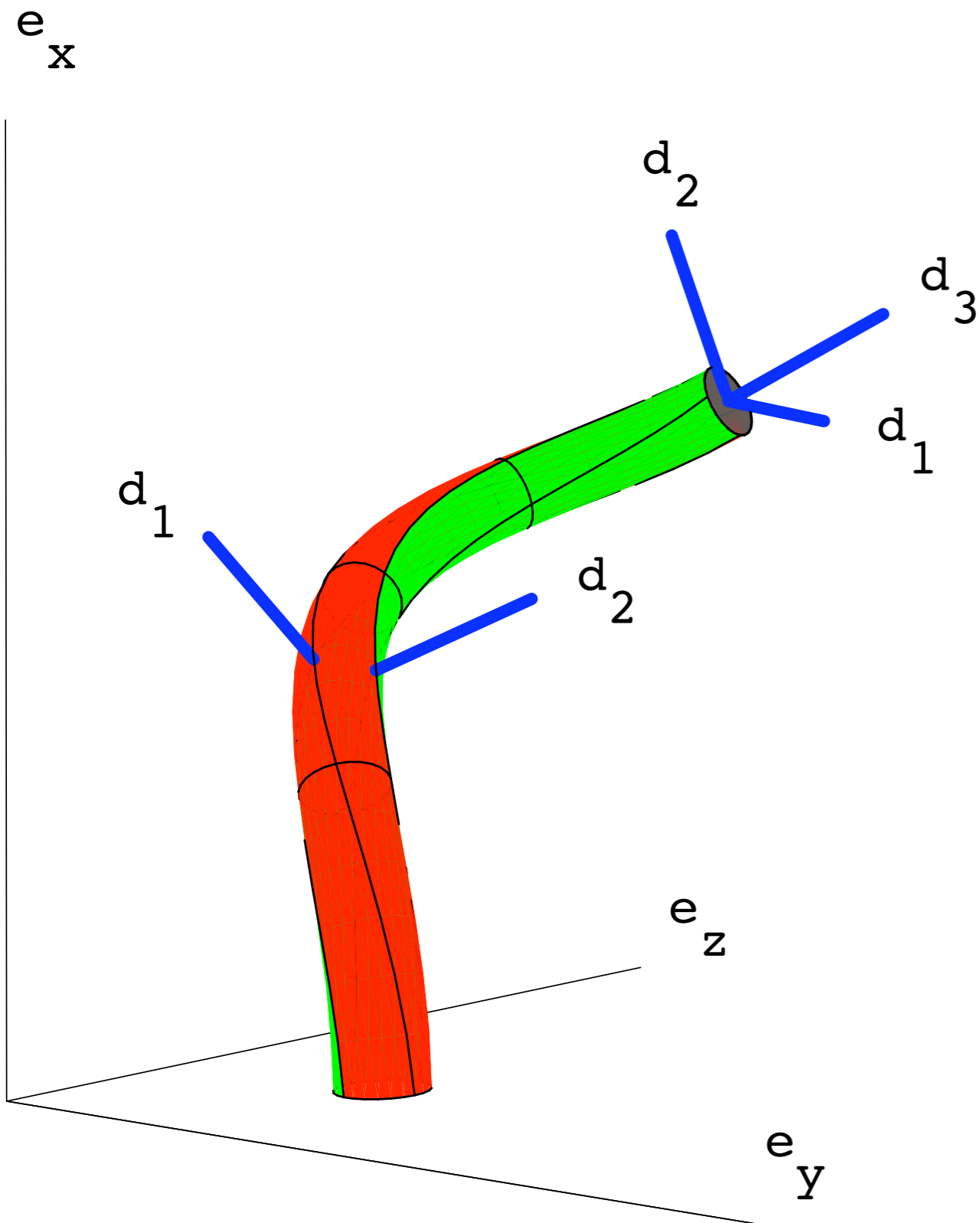


Dynamics (angular momentum):

$$M'(s, t) + r'(s, t) \times F(s, t) = \rho I \dot{\omega}(s, t)$$

$$\dot{\quad} \equiv \frac{d}{dt} \quad , \quad ' \equiv \frac{d}{ds}$$

Kirchhoff equations: kinematics



Cosserat frame

$$d'_1 = u \times d_1$$

$$d'_2 = u \times d_2$$

$$d'_3 = u \times d_3$$

$$u = \{ \kappa_1, \kappa_2, \tau \} d_i$$

curvatures

twist

Kirchhoff equations: constitutive relations

curvature

$M = EI\kappa$

with $\kappa(s) = \theta'(s)$

- G shear modulus
- E Young's modulus
- I second moment of area
- A area of section

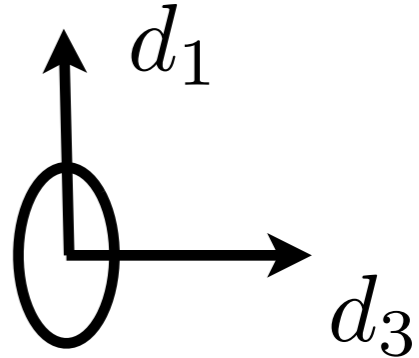
extension

$F_3 = EA(v_3 - 1)$

(v_1, v_3)
shear strains

shear

$F_1 = GA v_1$



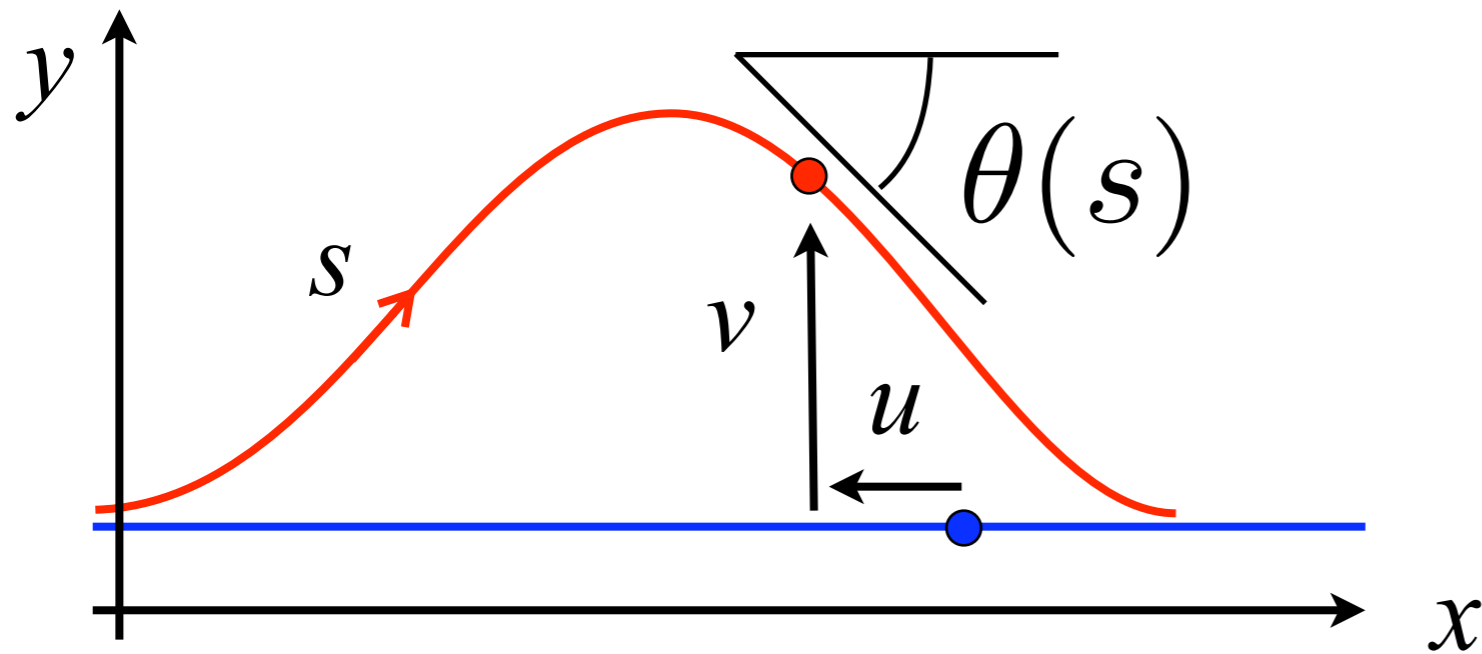
Special Cosserat theory of rods

S.Antman, *Nonlinear problems of elasticity*, (2004).

$$\begin{array}{rcl} F' & = & \rho A \ddot{R} \\ M' & = & F \times R' + \rho I \ddot{\theta} \\ R' & = & V = v_1 d_1 + v_3 d_3 \\ M & = & EI \theta' \\ F \cdot d_1 & = & GA v_1 \\ F \cdot d_3 & = & EA (v_3 - 1) \end{array} \left. \begin{array}{l} \\ \\ \\ \\ \\ \end{array} \right\} \begin{array}{l} \text{dynamics} \\ \\ \text{kinematics} \\ \\ \text{constitutive} \\ \text{relations} \end{array}$$

in the (x,y) plane

Strength of materials notations

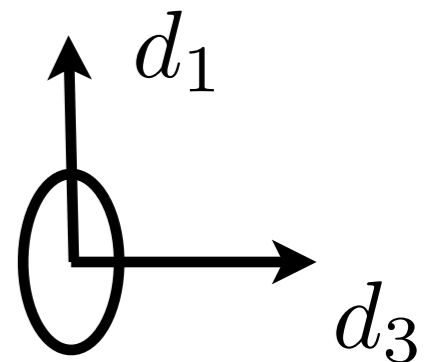


$$\begin{cases} x(s) & = & s - u(s) \\ y(s) & = & v(s) \\ x'(s) & = & 1 - u'(s) \\ y'(s) & = & v'(s) \end{cases}$$

$$R'(s) = \begin{pmatrix} x'(s) \\ y'(s) \end{pmatrix} = V(s) = v_1 d_1 + v_3 d_3$$

$$d_1(s) = \begin{pmatrix} -\sin \theta \\ \cos \theta \end{pmatrix} \quad d_3(s) = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}$$

with $\begin{cases} v_1(s) & = & (F \cdot d_1)/GA \\ v_3(s) & = & 1 + (F \cdot d_3)/EA \end{cases}$



Equilibrium equations (adim)

$$f = \frac{FL^2}{EI} \quad m = \frac{ML}{EI} \quad x = \frac{X}{L} \quad s = \frac{S}{L} \quad \nu : \text{Poisson}$$

$$\left\{ \begin{array}{l} x' = \cos \theta + \epsilon (f_3 \cos \theta - 2(1 + \nu) f_1 \sin \theta) \\ y' = \sin \theta + \epsilon (f_3 \sin \theta + 2(1 + \nu) f_1 \cos \theta) \\ \theta' = m \\ m' = -f_1 + \epsilon f_1 f_3 (1 - 2\nu) \\ f'_x = 0 \\ f'_y = 0 \end{array} \right.$$

boundary conditions

$$\begin{array}{l} x(0) = 0 \\ y(0) = 0 = y(1) \\ \theta(0) = 0 = \theta(1) \end{array}$$

$$\epsilon = \frac{I}{AL^2} = \frac{1}{12} \left(\frac{h}{L} \right)^2 \ll 1$$

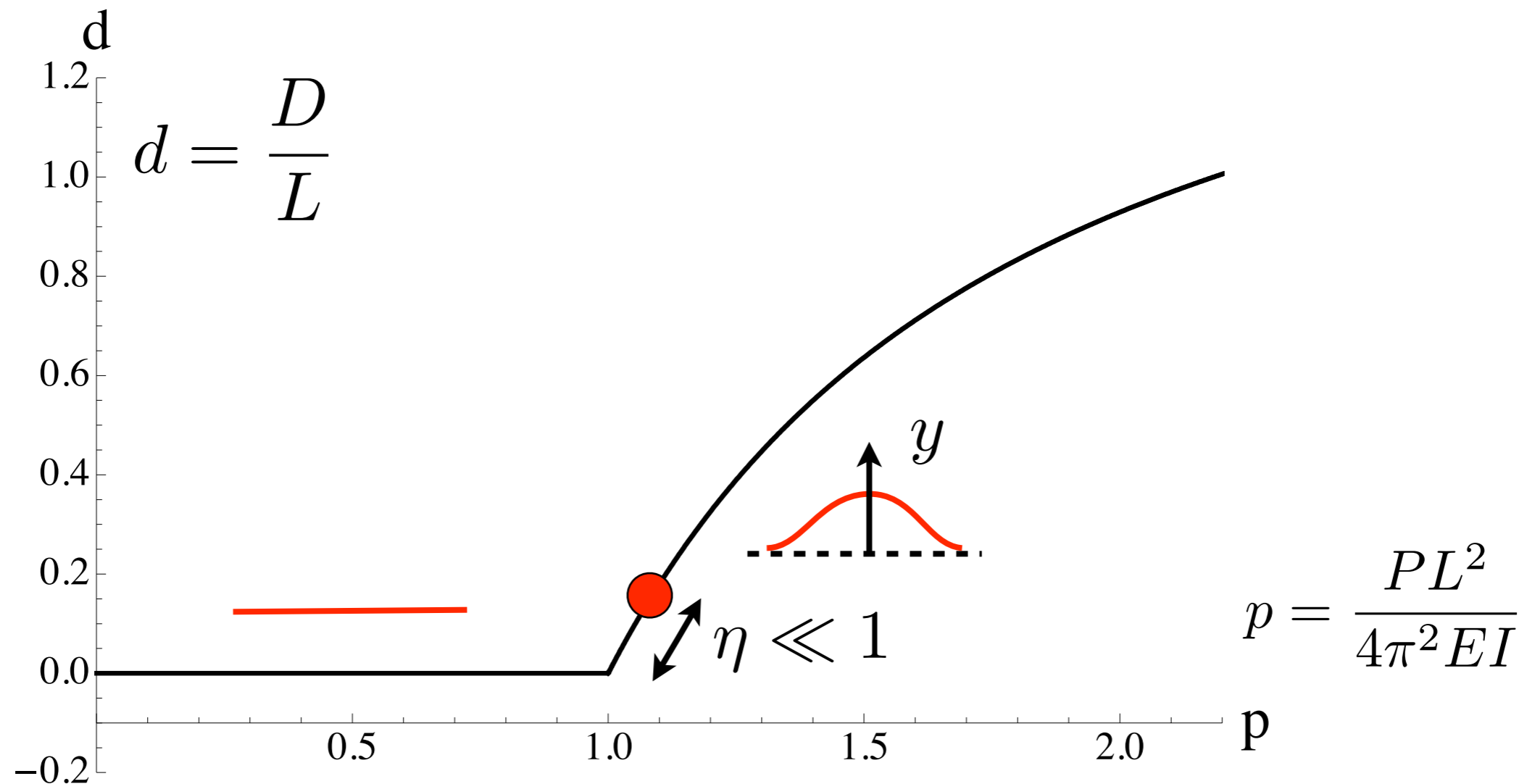
with $\left\{ \begin{array}{l} f_1 = -f_x \sin \theta + f_y \cos \theta \\ f_3 = f_x \cos \theta + f_y \sin \theta \end{array} \right.$

$\epsilon = 0$ Euler-Bernoulli beam

$\epsilon > 0$ Timoshenko beam

Equilibrium: analytical study

$\epsilon = 0$ Euler-Bernoulli



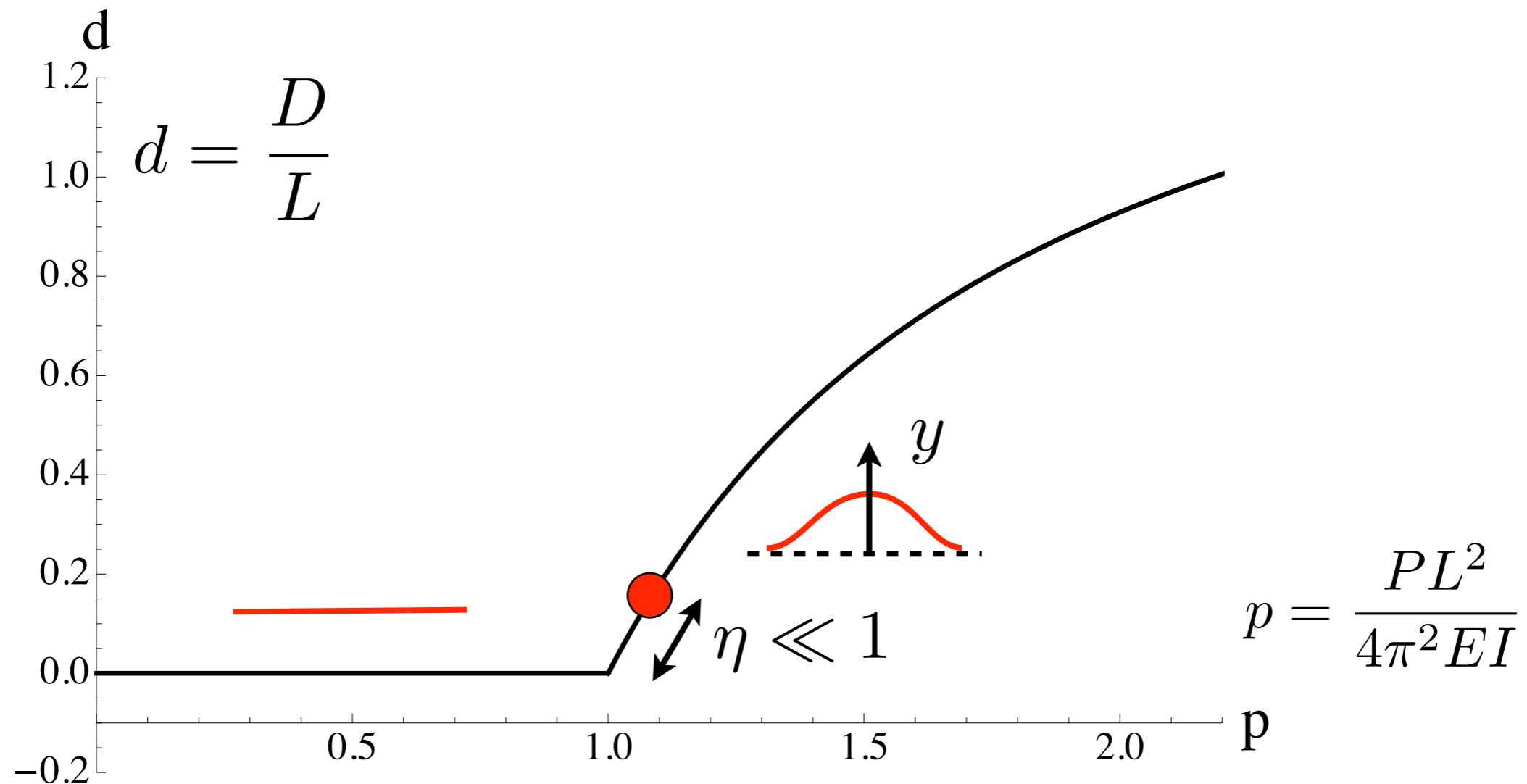
$$\theta(s) = \eta \theta_1(s) + \eta^2 \theta_2(s) + \eta^3 \theta_3(s) + O(\eta^4)$$

$$y(s) = \eta y_1(s) + \eta^2 y_2(s) + \eta^3 y_3(s) + O(\eta^4)$$

$$p = p_0 + \eta p_1 + \eta^2 p_2 + \eta^3 p_3 + O(\eta^4)$$

Equilibrium: analytical study

$\epsilon = 0$ Euler-Bernoulli



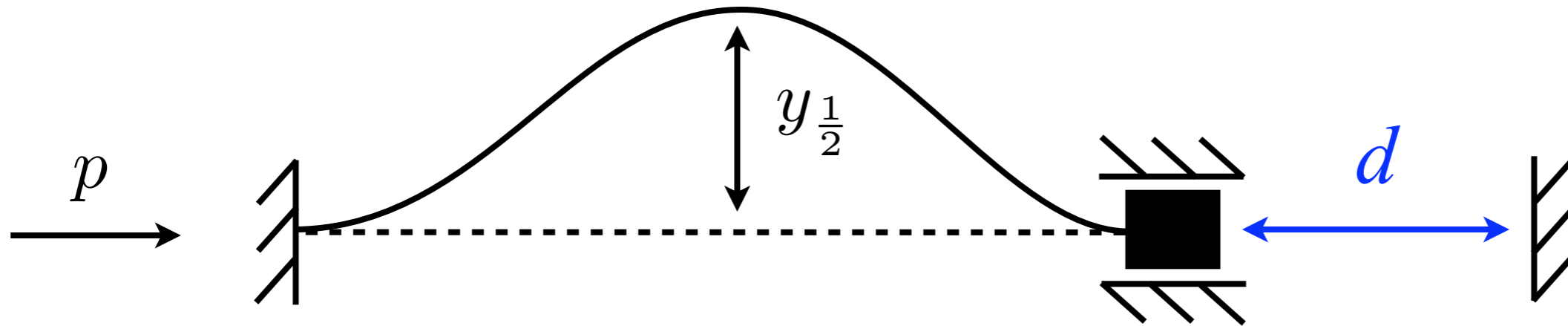
$$\theta(s) = \eta \sin 2\pi s + \frac{\eta^3}{48} \cos^2(2\pi s) \sin(2\pi s) + O(\eta^4)$$

$$y(s) = \frac{\eta}{2\pi} (1 - \cos 2\pi s) + \frac{\eta^3}{384\pi} (-20 + 23 \cos(2\pi s) - 3 \cos(6\pi s)) + O(\eta^4)$$

$$p = 4\pi^2 + \eta^2 \pi^2 / 2 + O(\eta^4)$$

Equilibrium: analytical study

$\epsilon = 0$ Euler-Bernoulli



$$p - p_c = \eta^2 \frac{\pi^2}{2} \quad \text{with } (p_c = 4\pi^2)$$

$$y_{\frac{1}{2}} = \frac{\eta}{\pi}$$

$$d = \frac{\eta^2}{4}$$

Vibrations: analytical study

$\epsilon = 0$ Euler-Bernoulli

small amplitude vibrations around pre or post-buckled equilibrium

$$y(s, t) = y_E(s) + \delta y(s) e^{i\omega t} \quad \text{with } |\delta y(s)| \ll 1$$

$$\left\{ \begin{array}{l} \delta m' = \delta \theta f_{3E} - \delta f_y \cos \theta_E + \delta f_x \sin \theta_E \\ \delta \theta' = \delta m \\ \delta y' = \cos \theta_E \delta \theta \\ \delta x' = -\sin \theta_E \delta \theta \\ \delta f'_x = -\omega^2 \delta x \\ \delta f'_y = -\omega^2 \delta y \end{array} \right.$$

boundary conditions

$$\delta x(0) = 0 = \delta x(1)$$

$$\delta y(0) = 0 = \delta y(1)$$

$$\delta \theta(0) = 0 = \delta \theta(1)$$

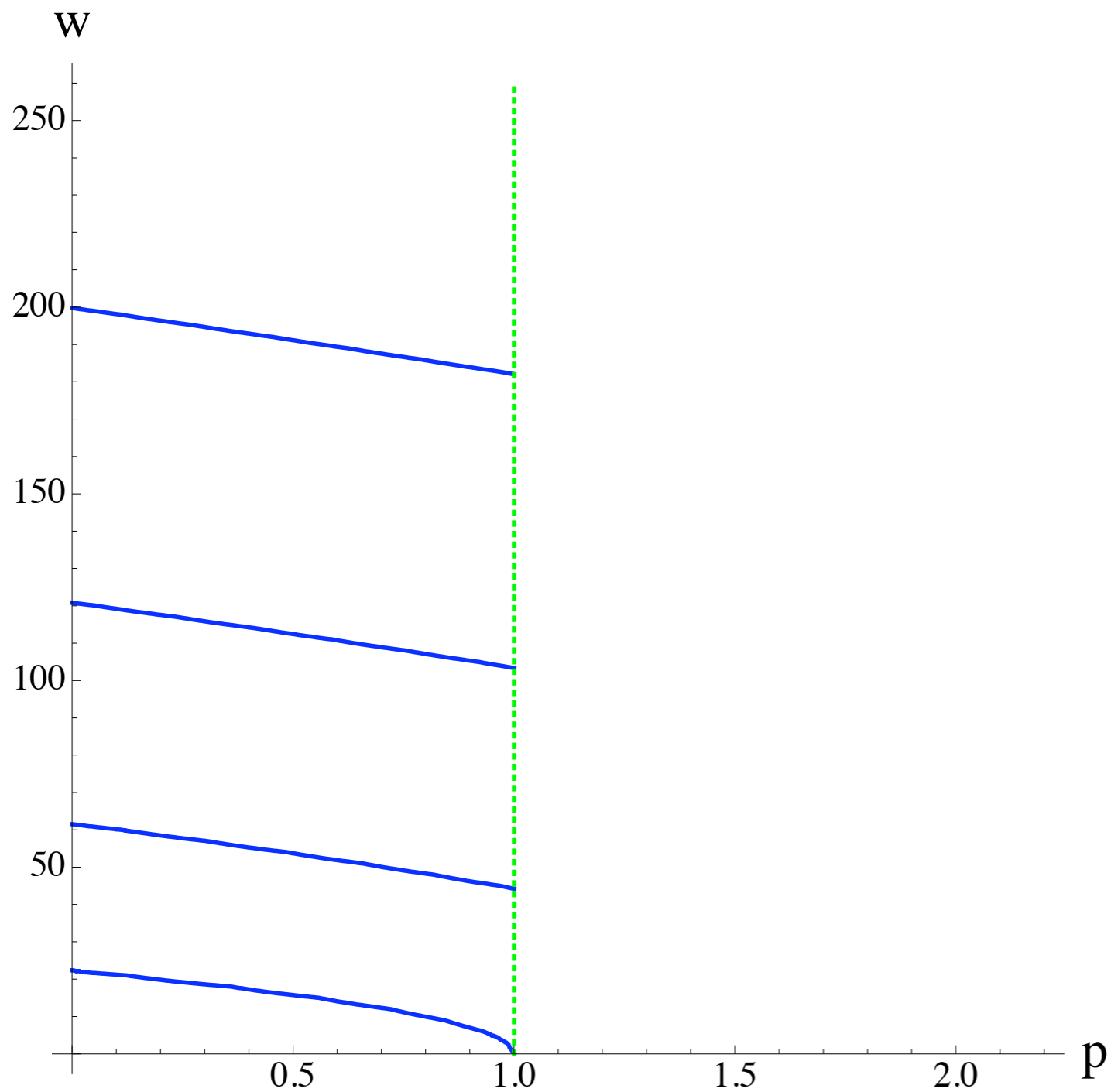
pre-buckled $\theta_E(s) \equiv 0$

post-buckled $\theta_E(s) = \eta \sin 2\pi s + \frac{\eta^3}{48} \cos^2(2\pi s) \sin(2\pi s) + O(\eta^4)$

Vibrations: analytical study

$\epsilon = 0$ Euler-Bernoulli

$$\text{poly0} = -2 n \sqrt{n^2 - p} + 2 n \sqrt{n^2 - p} \text{Cos}[n] \text{Cosh}[\sqrt{n^2 - p}] + p \text{Sin}[n] \text{Sinh}[\sqrt{n^2 - p}]$$

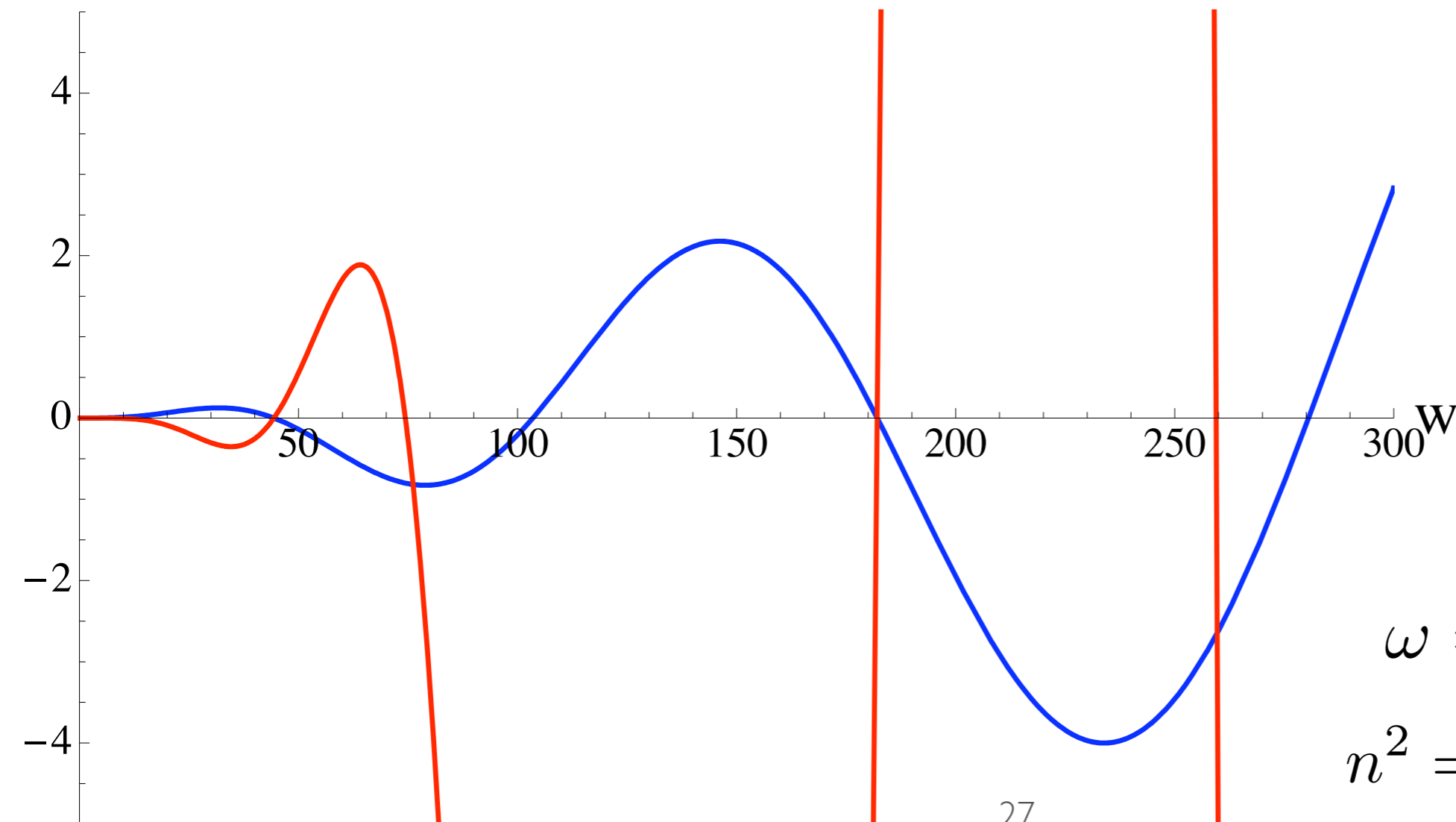


Vibrations: analytical study

$\epsilon = 0$ Euler-Bernoulli

$$\text{poly1} = -2 n \sqrt{n^2 - 4 \pi^2} + 2 n \sqrt{n^2 - 4 \pi^2} \cos[n] \cosh\left[\sqrt{n^2 - 4 \pi^2}\right] + 4 \pi^2 \sin[n] \sinh\left[\sqrt{n^2 - 4 \pi^2}\right];$$

$$\begin{aligned} \text{poly3} = & n \sqrt{n^2 - 4 \pi^2} \cosh\left[\sqrt{n^2 - 4 \pi^2}\right] \left(-n (n^2 - 4 \pi^2) \cos[n] + 4 (n^2 - 2 \pi^2) \sin[n]\right) + \\ & n \sqrt{n^2 - 4 \pi^2} \left(n^3 - 4 n \pi^2 + (-4 n^2 + 8 \pi^2) \sin[n]\right) + \\ & 2 (-n^2 + 4 \pi^2) \left(2 n^2 - 4 \pi^2 + (-2 n^2 + 4 \pi^2) \cos[n] + n \pi^2 \sin[n]\right) \sinh\left[\sqrt{n^2 - 4 \pi^2}\right]; \end{aligned}$$

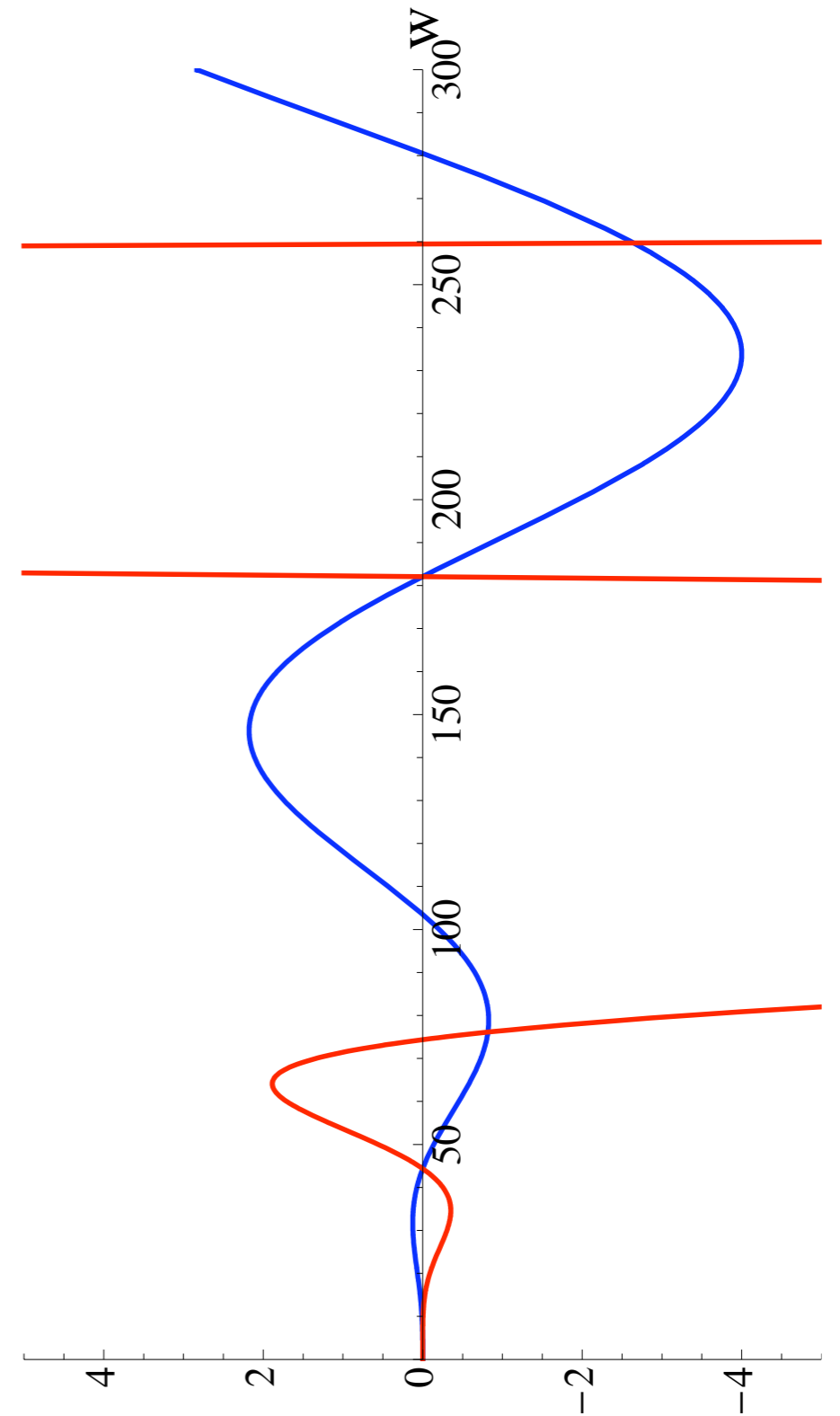


$$\omega = n \sqrt{n^2 - 4 \pi^2}$$

$$n^2 = 2 \pi^2 + \sqrt{\omega^2 + 4 \pi^4}$$

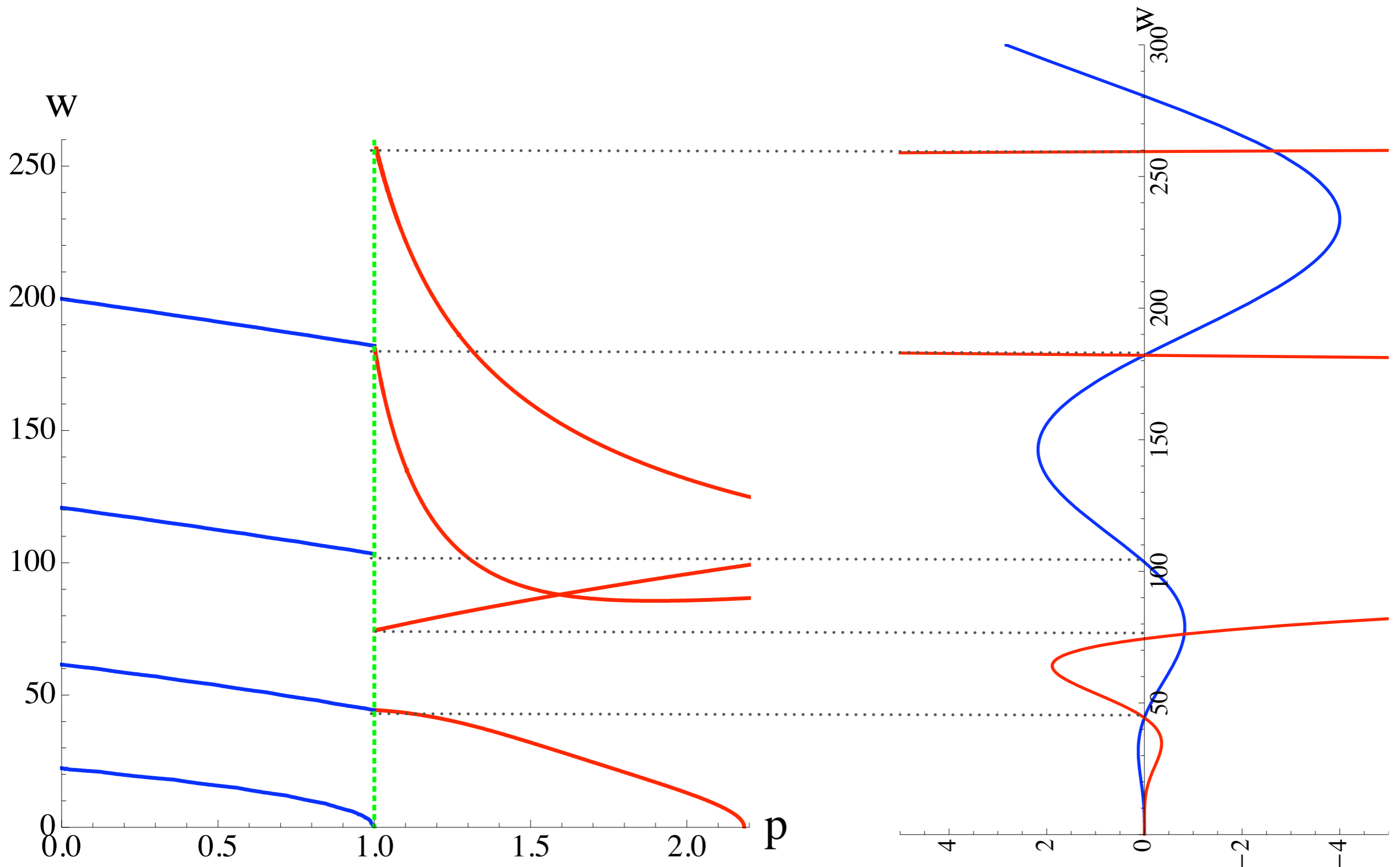
Vibrations: analytical study

$\epsilon = 0$ Euler-Bernoulli



Vibrations: analytical study

$\epsilon = 0$ Euler-Bernoulli



Dynamics

$\epsilon > 0$ Timoshenko

with shear, extension, and rotational inertia

$$0 < \epsilon = \frac{I}{AL^2} = \frac{1}{12} \left(\frac{h}{L} \right)^2 \ll 1$$

$$\left\{ \begin{array}{l} x' = \cos \theta + \epsilon (f_3 \cos \theta - 2(1 + \nu) f_1 \sin \theta) \\ y' = \sin \theta + \epsilon (f_3 \sin \theta + 2(1 + \nu) f_1 \cos \theta) \\ \theta' = m \\ m' = -f_1 + \epsilon f_1 f_3 (1 - 2\nu) + \epsilon \ddot{\theta} \\ f'_x = \ddot{x} \\ f'_y = \ddot{y} \end{array} \right.$$

ν : Poisson

$$\text{with } \begin{cases} f_1 = -f_x \sin \theta + f_y \cos \theta \\ f_3 = f_x \cos \theta + f_y \sin \theta \end{cases}$$

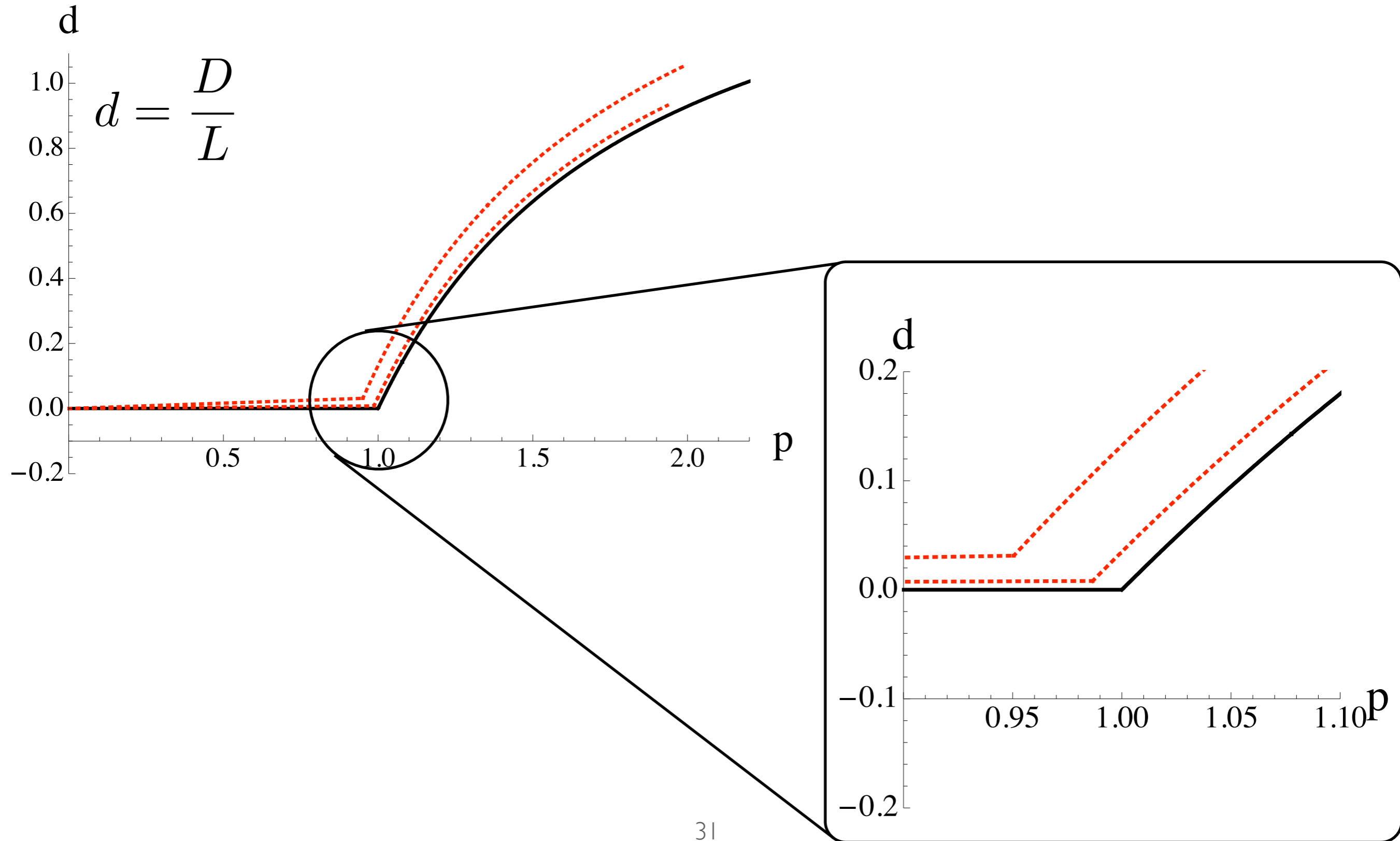
boundary conditions

$$\begin{array}{ll} x(0, t) = 0 & x(1, t) = 1 - d \\ y(0, t) = 0 & y(1, t) = 0 \\ \theta(0, t) = 0 & \theta(1, t) = 0 \end{array}$$

Equilibrium

$\epsilon > 0$ Timoshenko

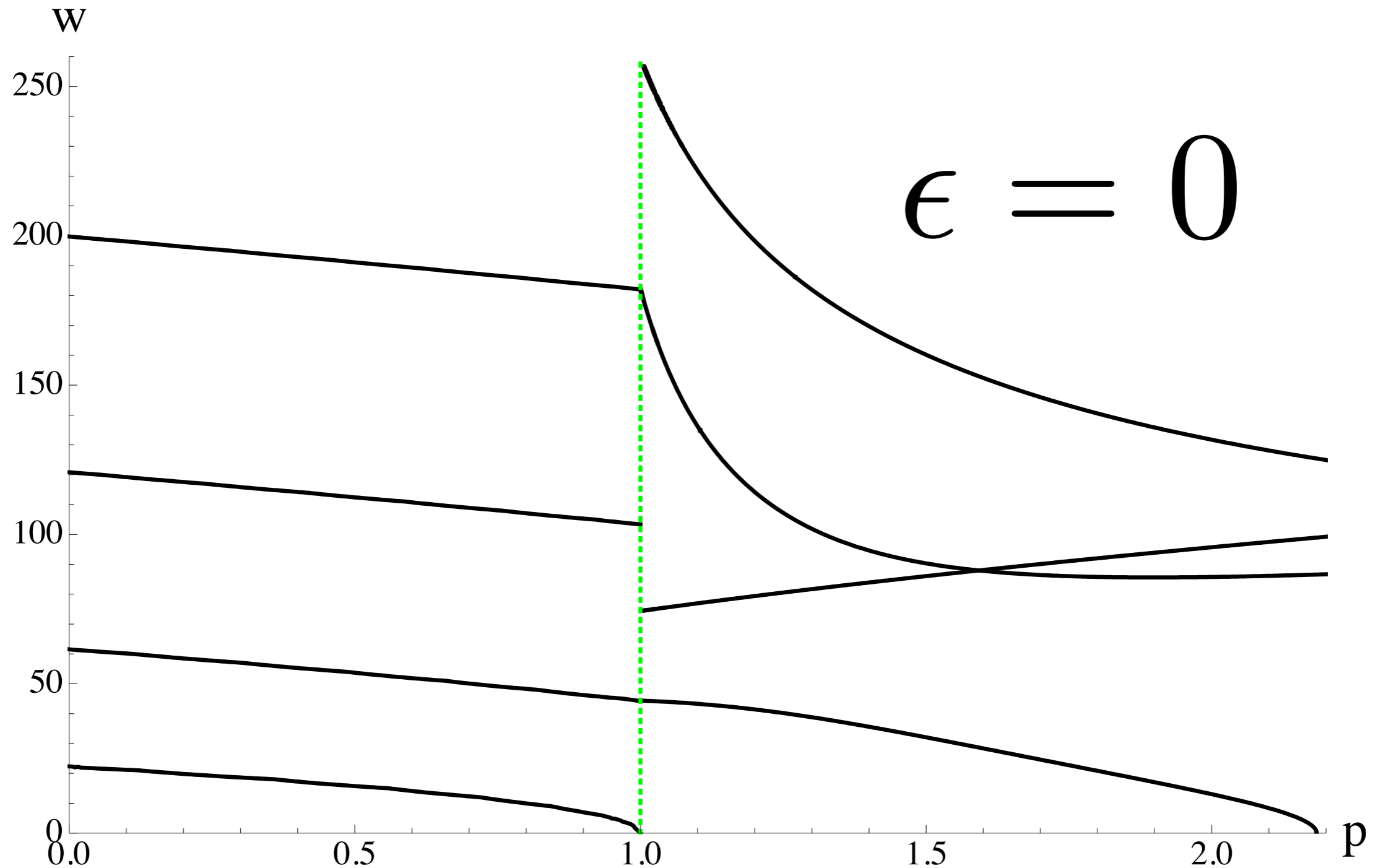
with shear, extension, and rotational inertia



Vibrations

$\epsilon > 0$ Timoshenko

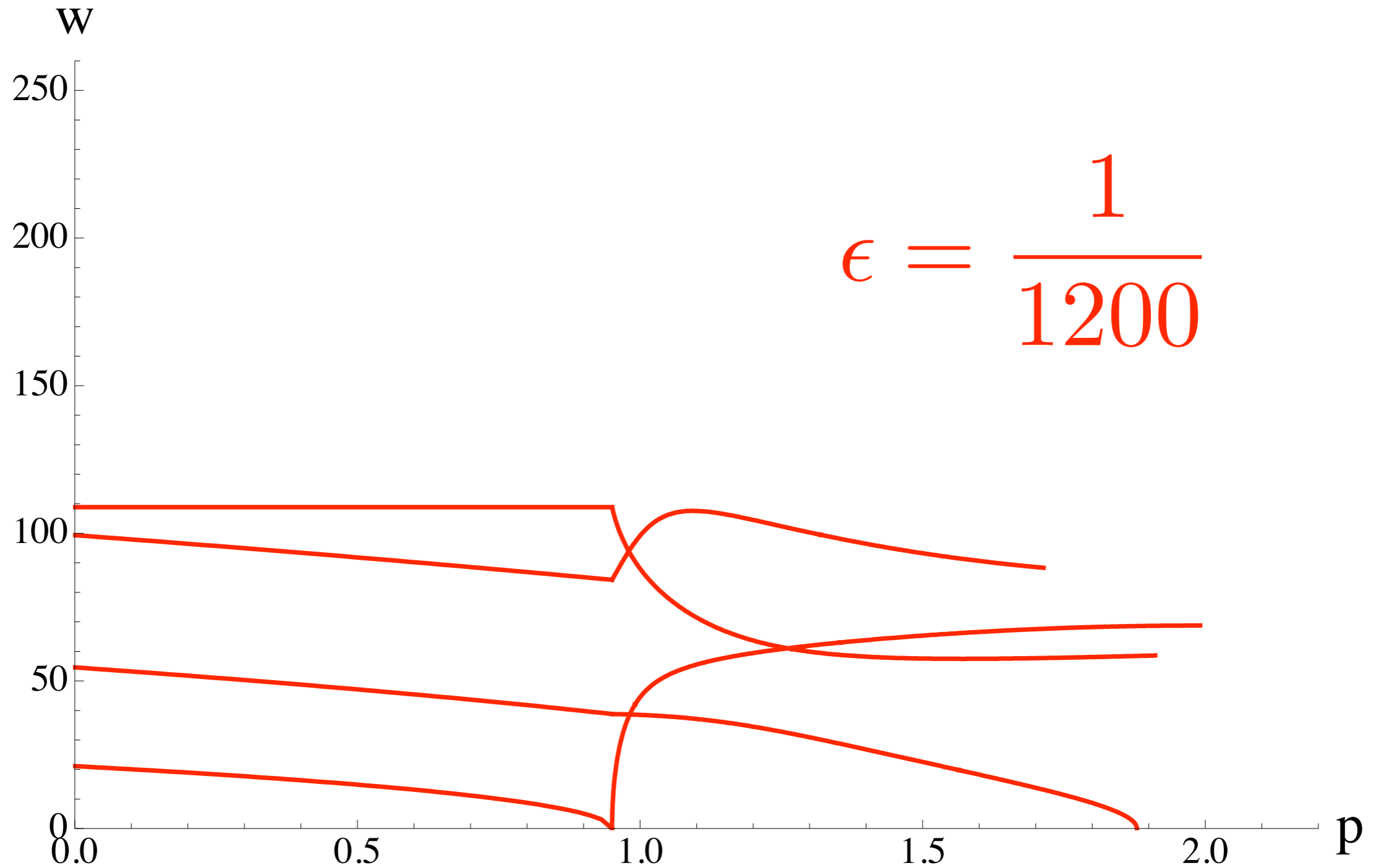
with shear, extension, and rotational inertia



Vibrations

$\epsilon > 0$ Timoshenko

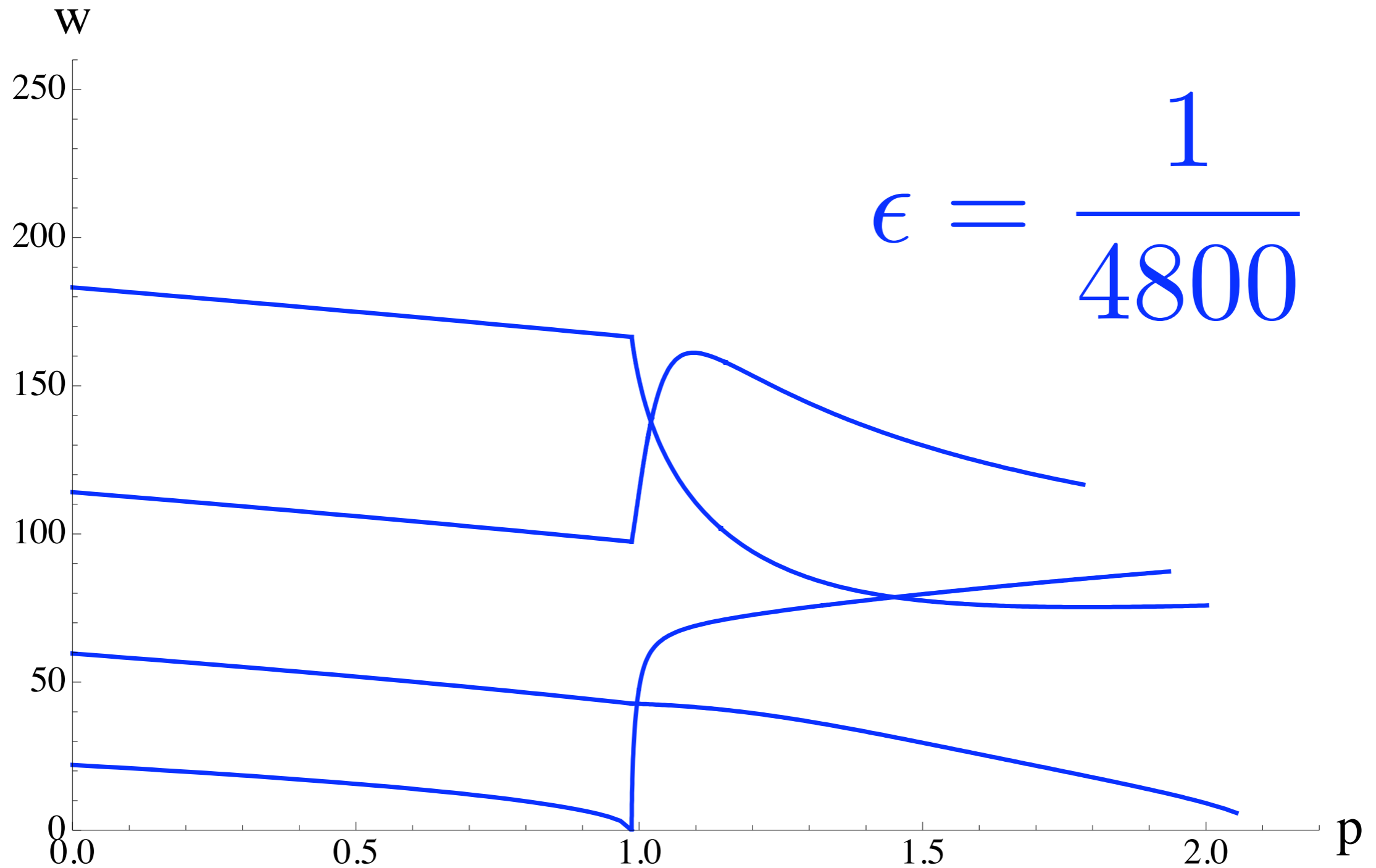
with shear, extension, and rotational inertia



Vibrations

$\epsilon > 0$ Timoshenko

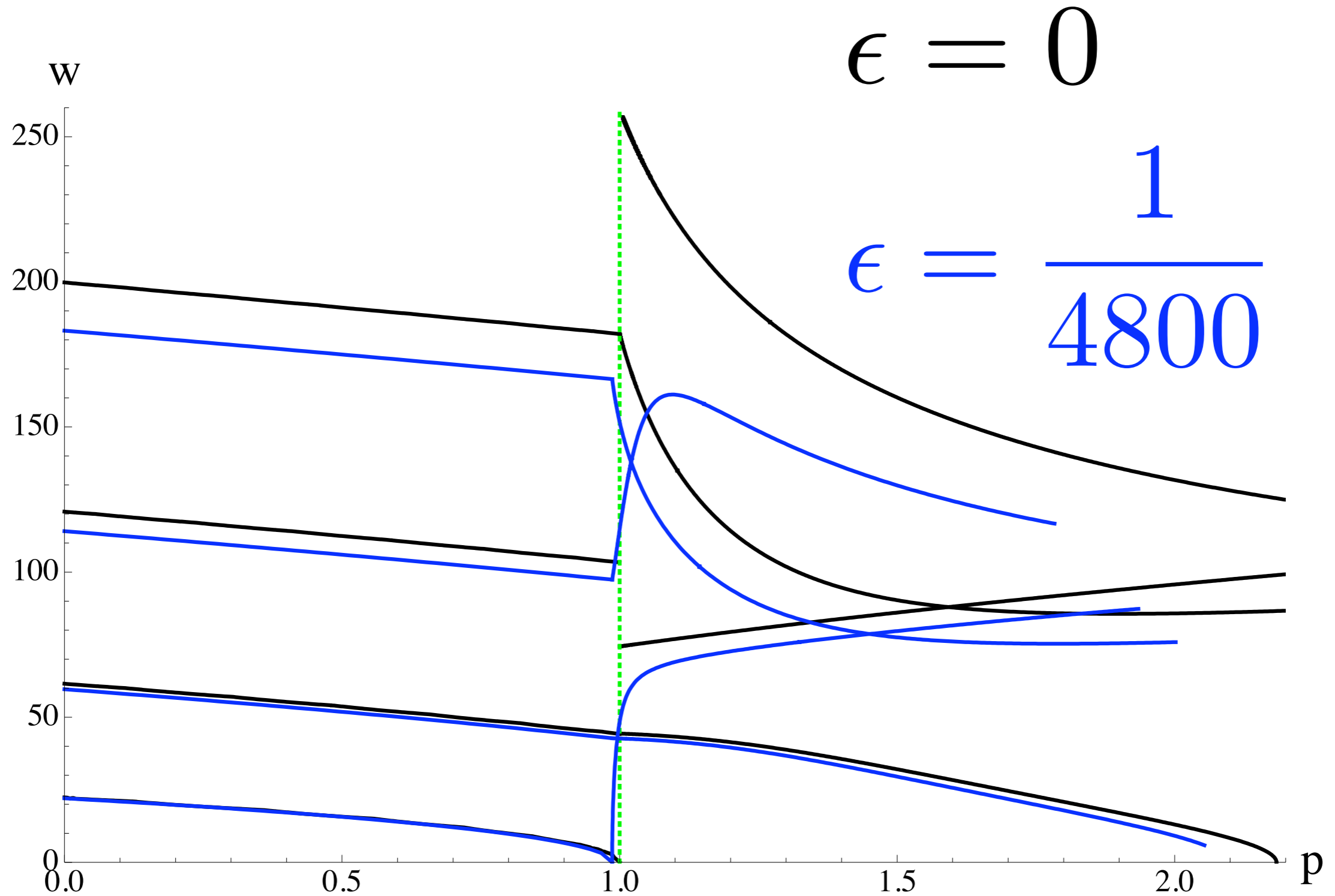
with shear, extension, and rotational inertia



Vibrations

$\epsilon > 0$ Timoshenko

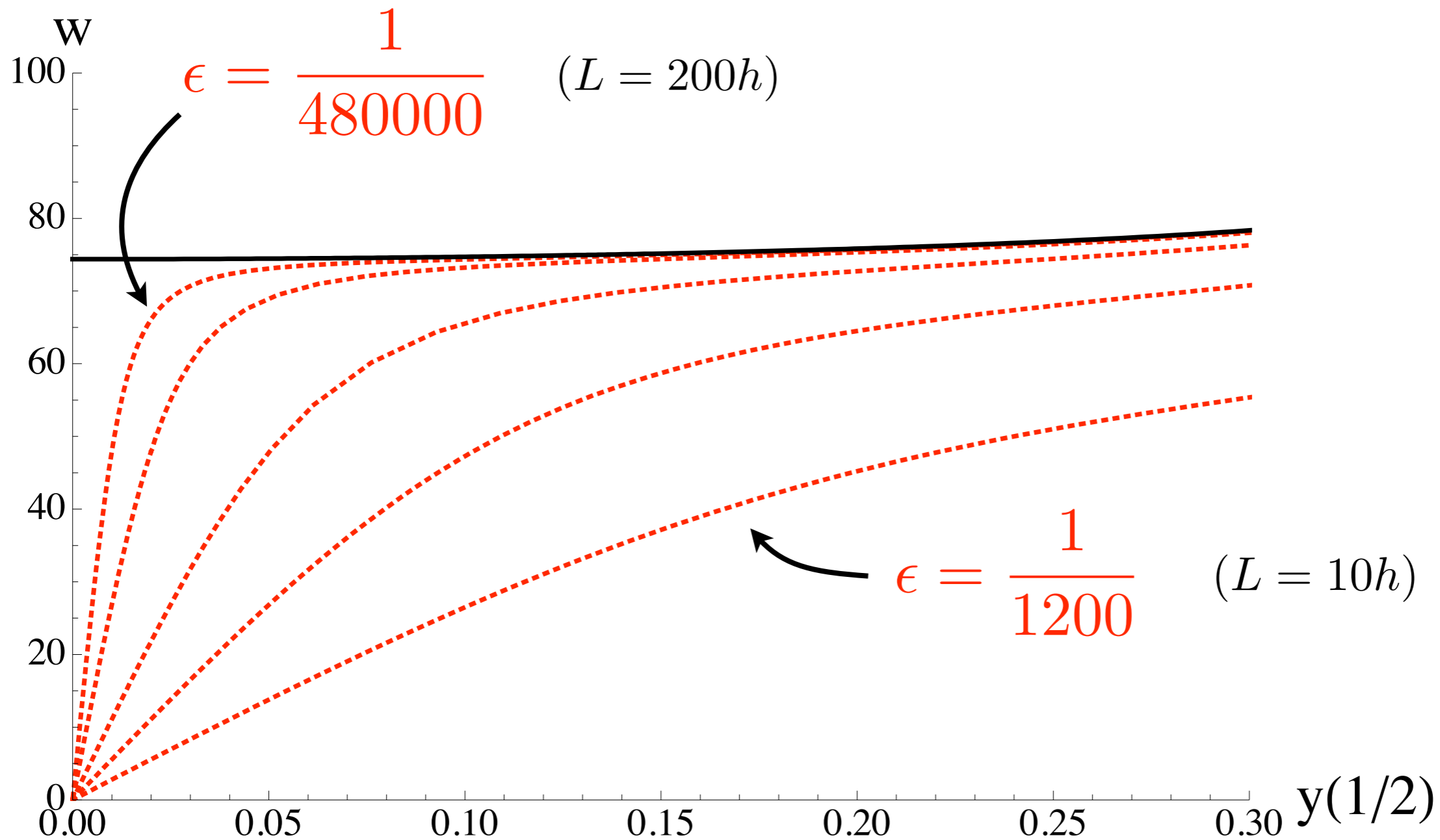
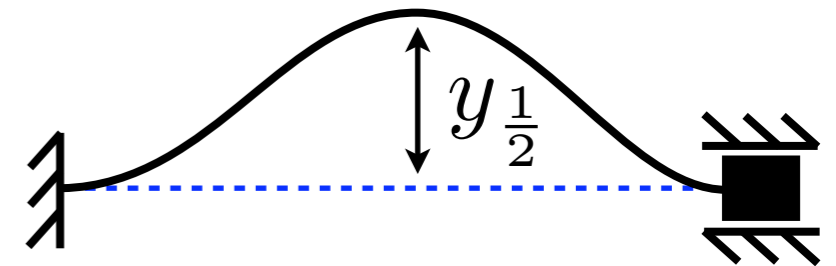
with shear, extension, and rotational inertia



Vibrations

$\epsilon > 0$ Timoshenko

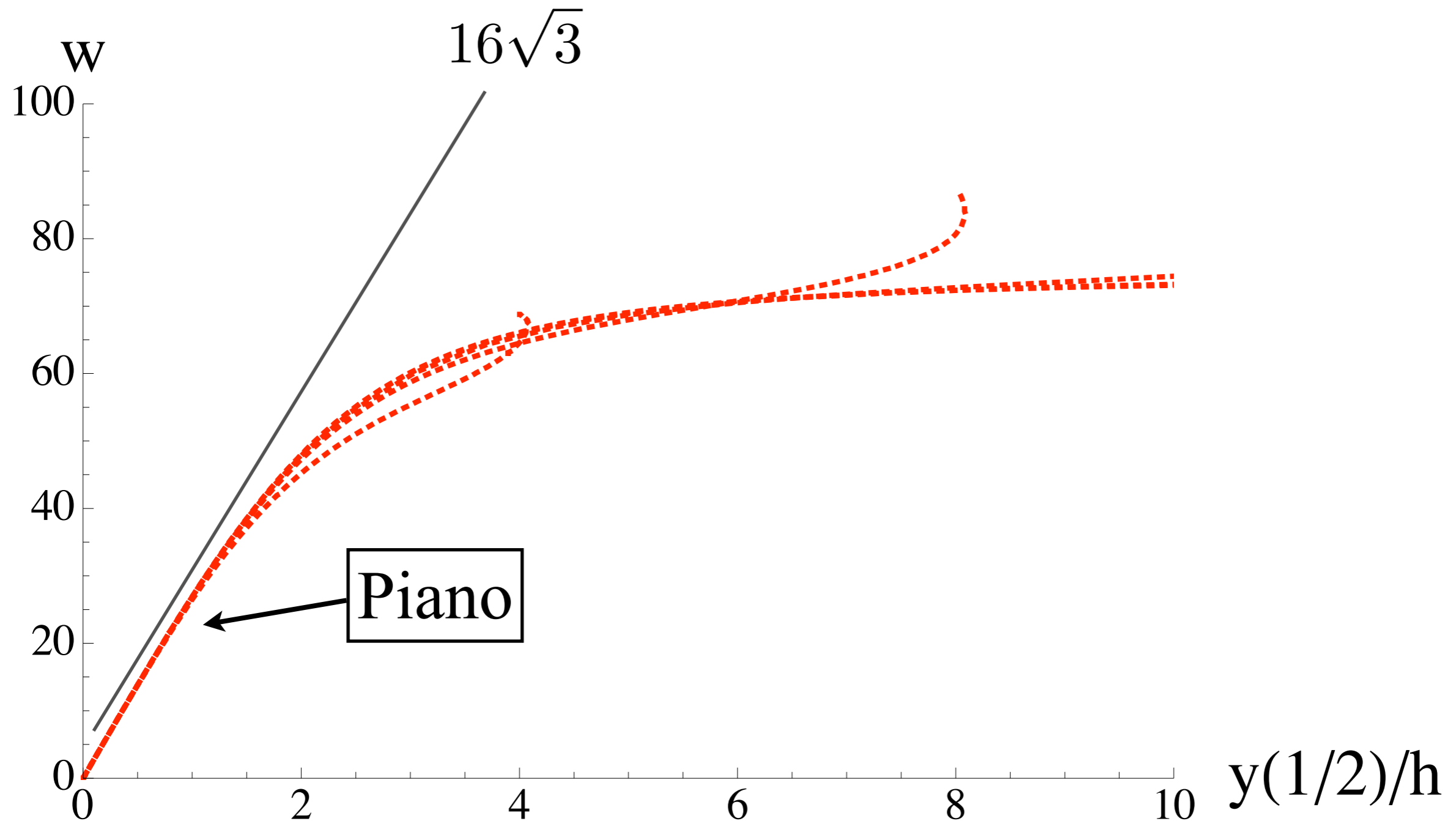
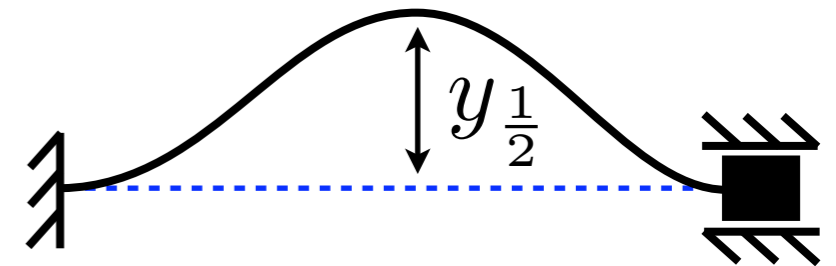
with shear, extension, and rotational inertia



Vibrations

with shear, extension, and rotational inertia

$\epsilon > 0$ Timoshenko

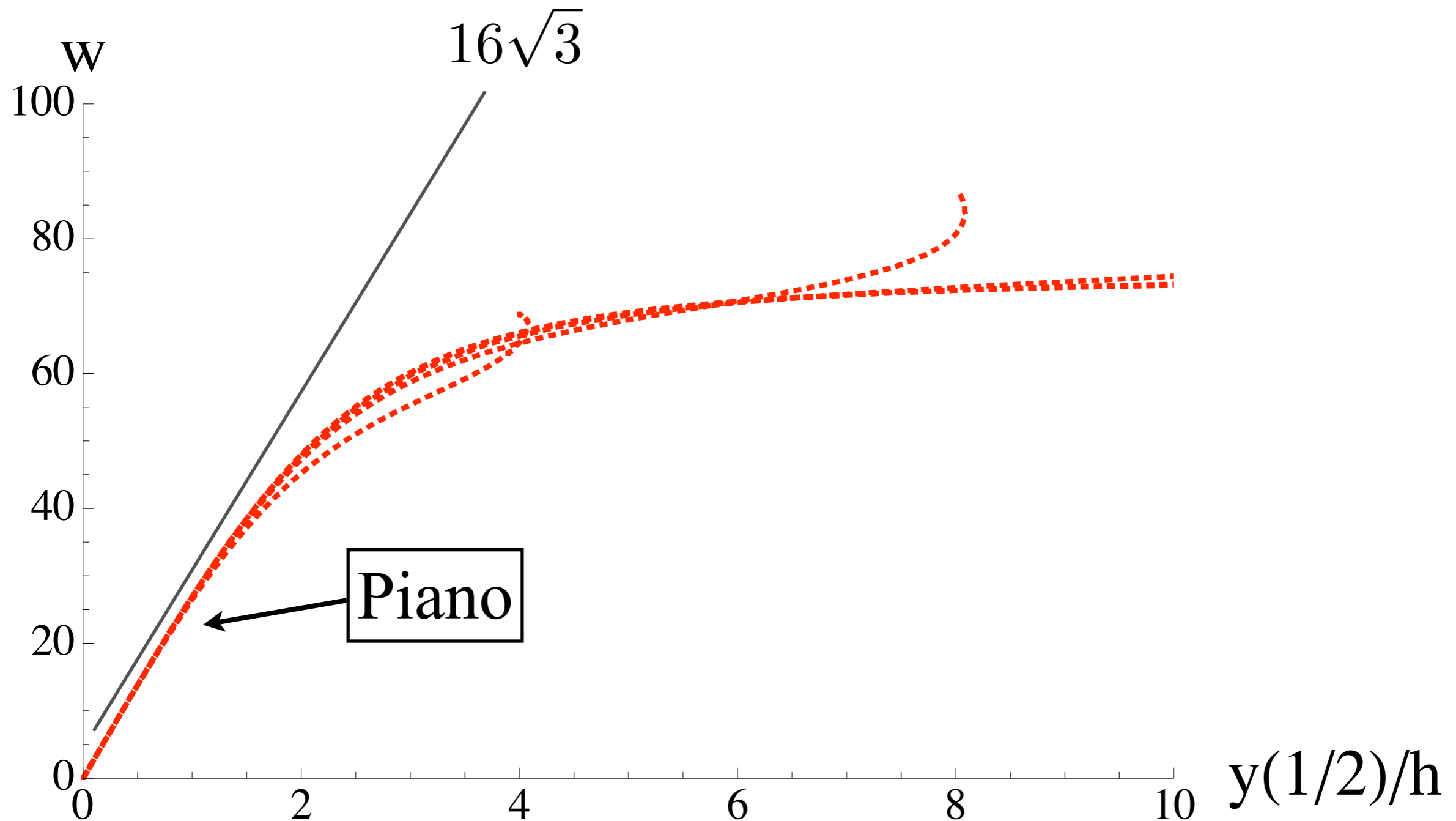


Vibrations

$\epsilon > 0$ Timoshenko

with shear, extension, and rotational inertia

$$\omega = \Omega L^2 \left(\frac{\rho A}{EI} \right)^{1/2}$$



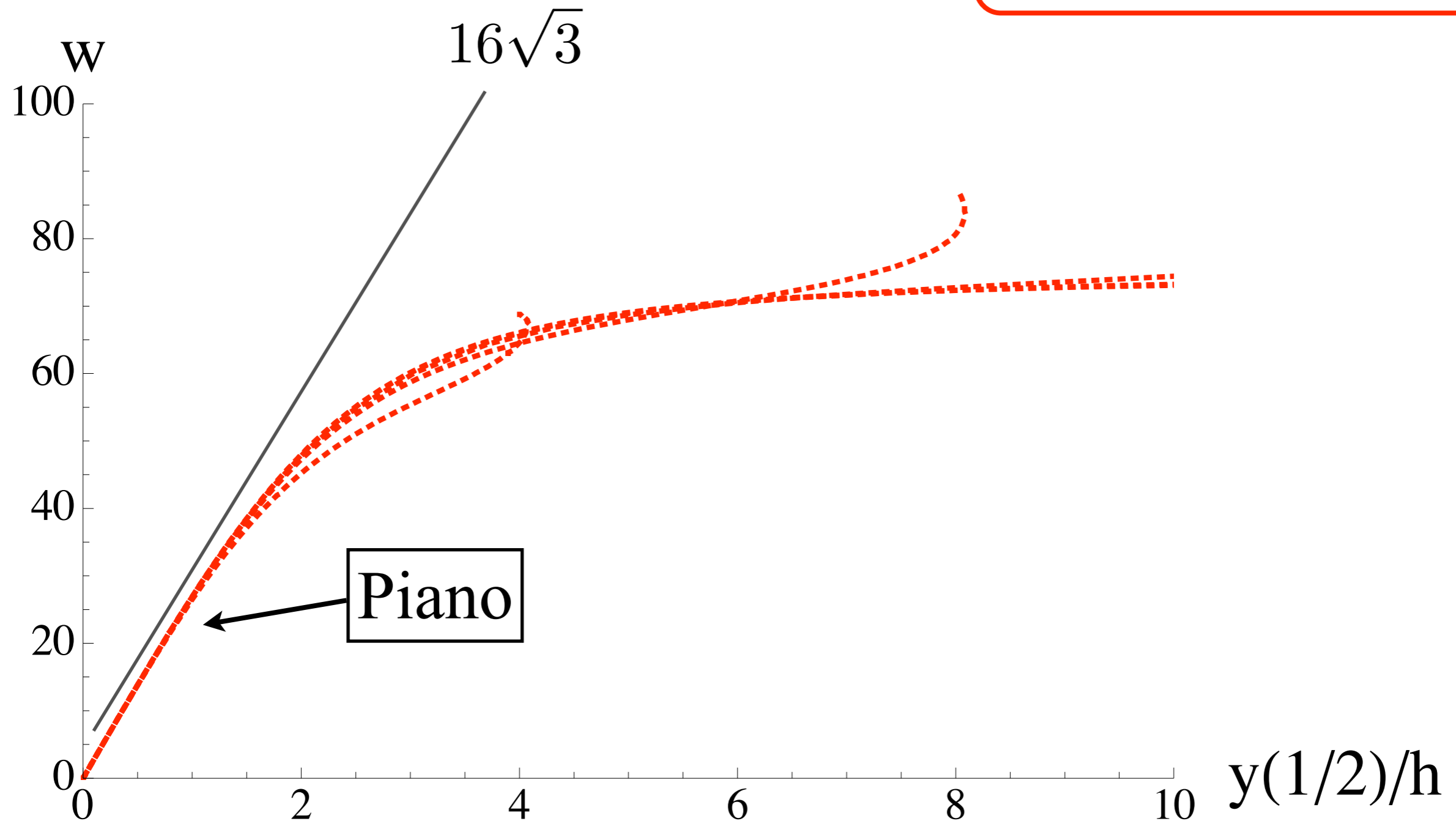
Vibrations

$\epsilon > 0$ Timoshenko

with shear, extension, and rotational inertia

$$\omega = \Omega L^2 \left(\frac{\rho A}{EI} \right)^{1/2}$$

$$\Omega = 8 \frac{Y_{1/2}}{L^2} \left(\frac{E}{\rho} \right)^{1/2}$$



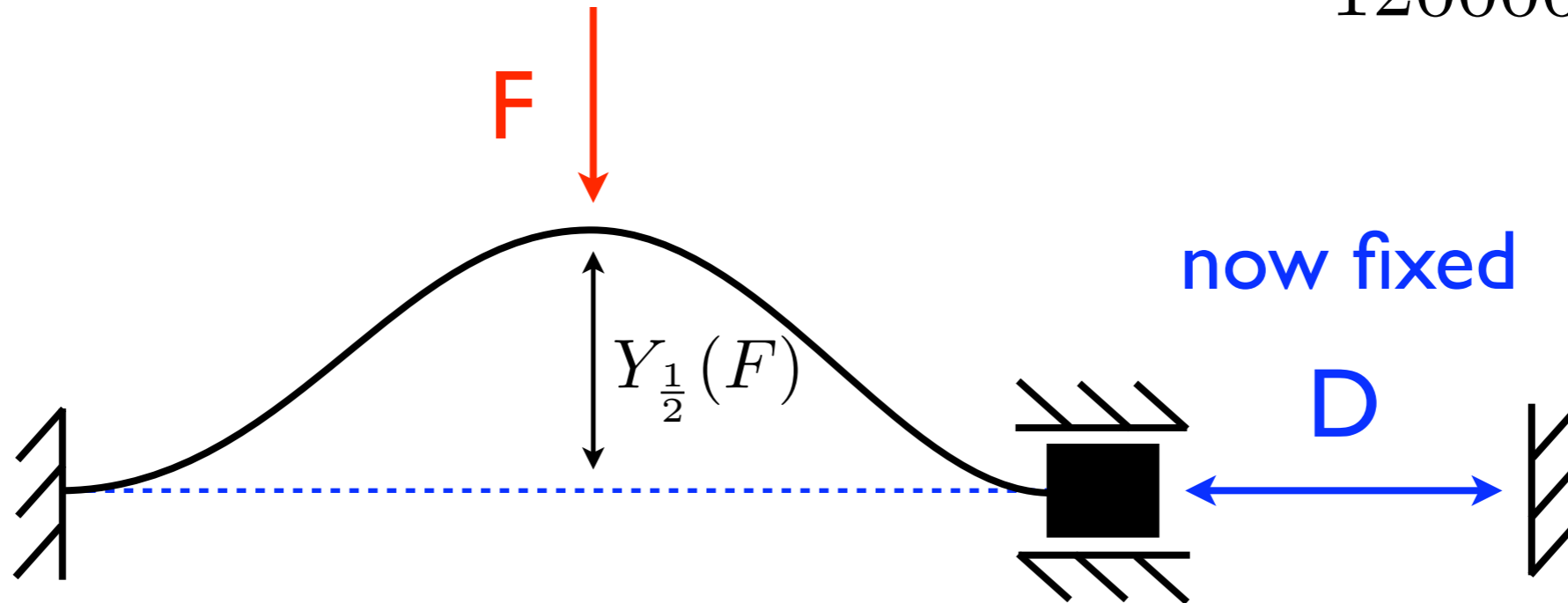
Vibrations

$\epsilon > 0$ Timoshenko

$$L = 100 h$$

$$\epsilon = \frac{1}{120000}$$

force is increased



we start with $Y_{\frac{1}{2}}(F = 0) = 1.0 h$

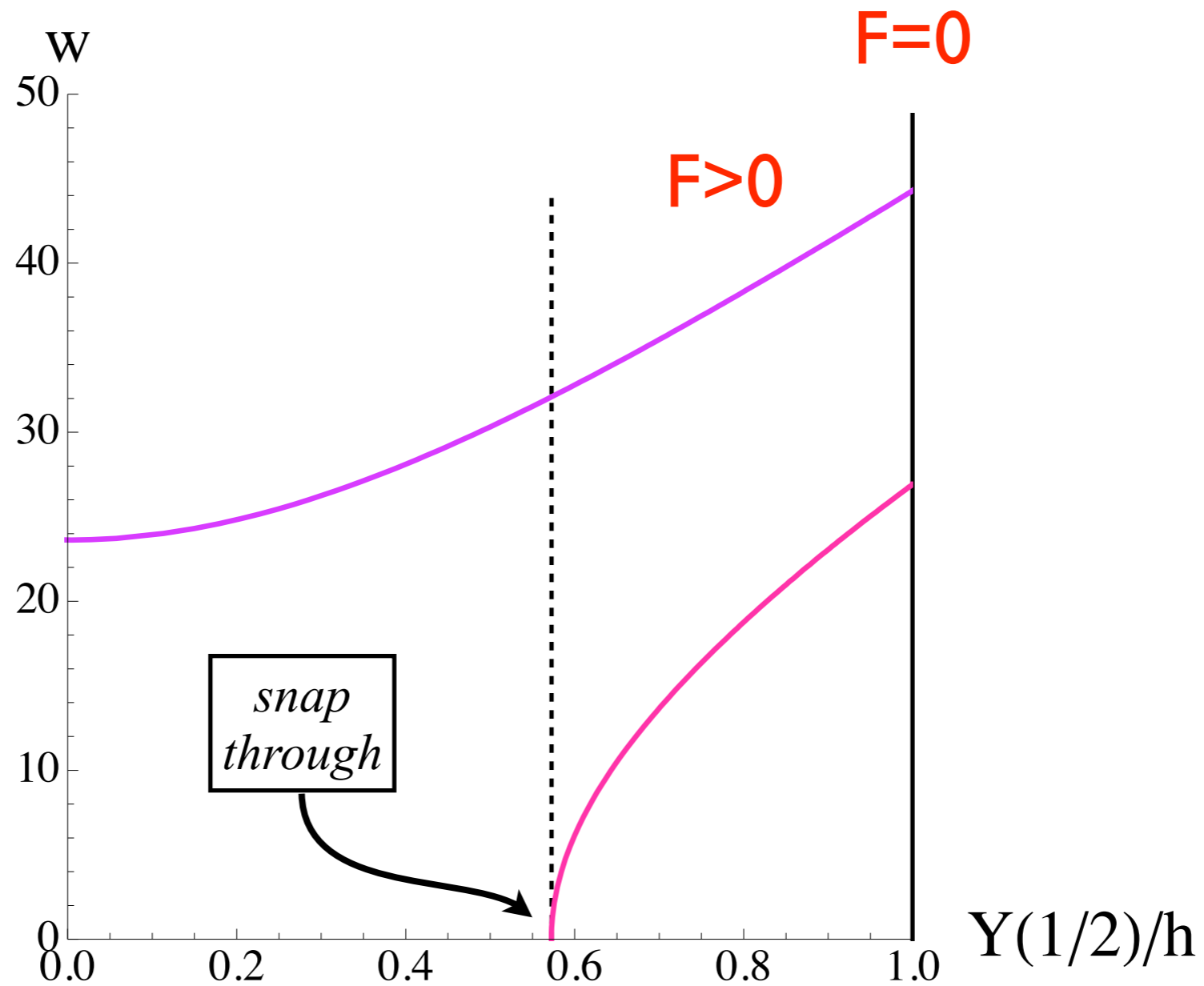
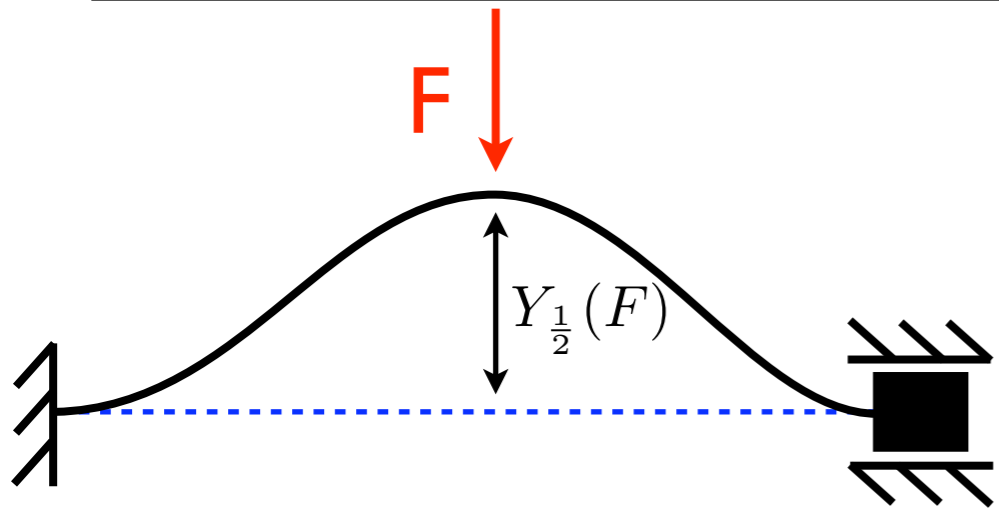
we look at $\omega = \omega(F)$

Vibrations

$\epsilon > 0$ Timoshenko

$$L = 100 h$$

$$\epsilon = \frac{1}{120000}$$



Perspectives

- *modes* -> *forced oscillations*
(*coupling with strings*)
- *beam* -> *plate*



Thank you