

Flag Flutter: Potential Flow Around a Rectangular Plate

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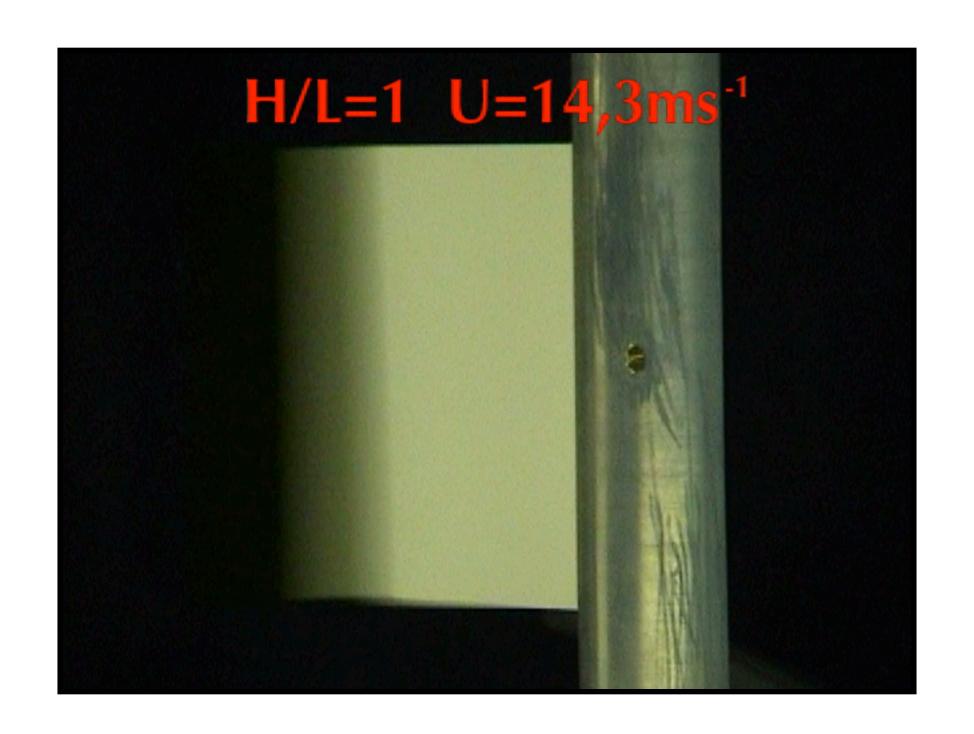
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Collaborators

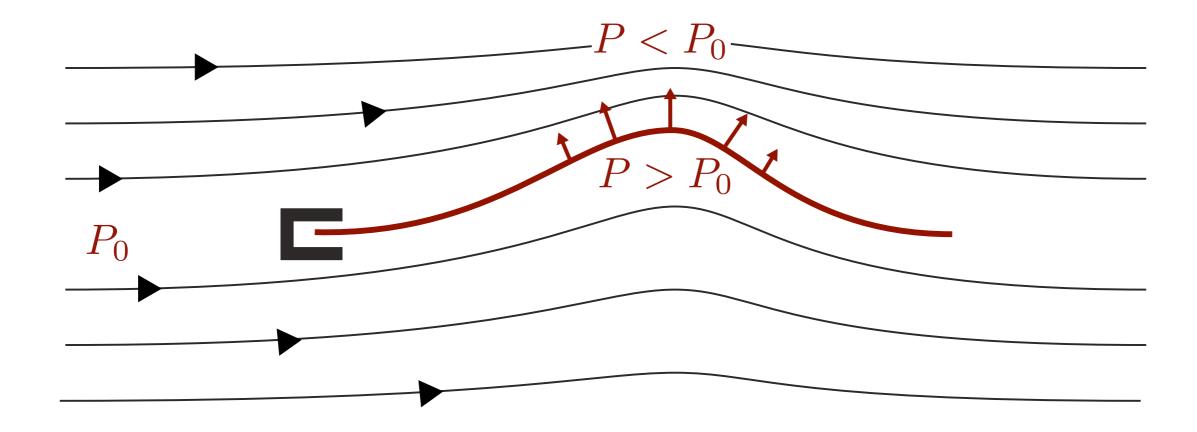
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ANR Contract "DRAPEAU"



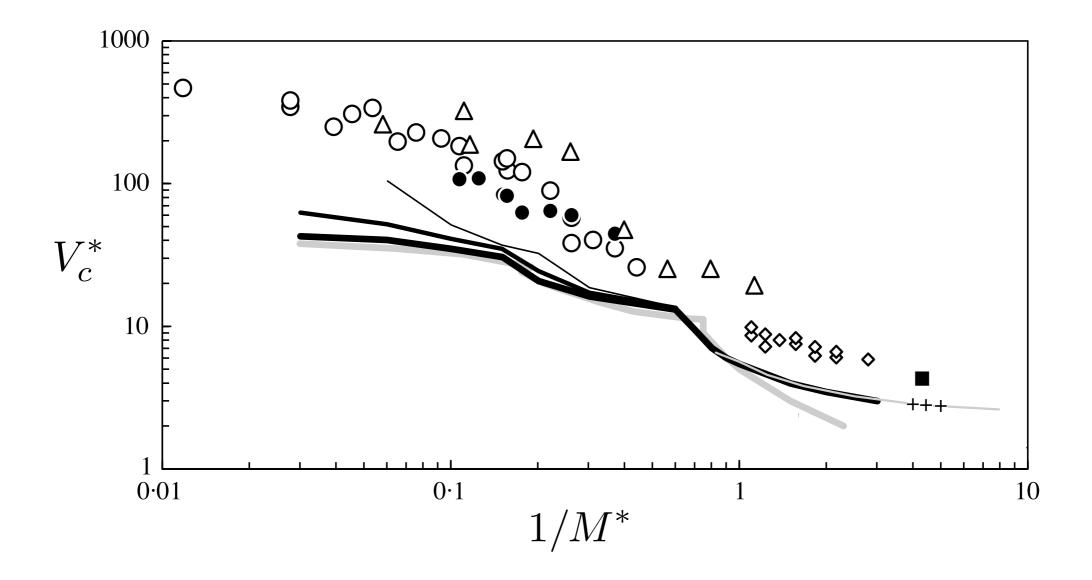
Instability Mechanism



Pressure: destabilising

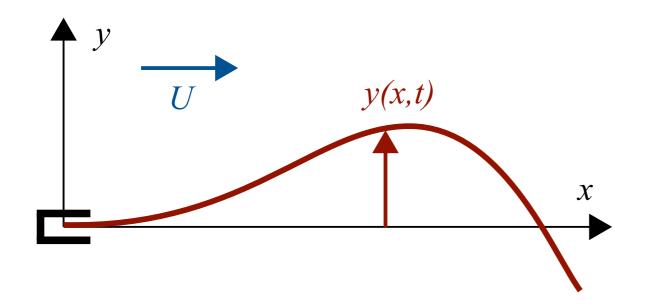
• Elasticity: stabilising

Literature



• [Ref] Watanabe, Isogai, Suzuki & Sugihara, J. Fluids Struct. (2002)

Equation of Motion



EI: flexural rigidity

 ρ : mass per unit area

Linearised Euler-Bernoulli beam equation

$$\rho \partial_t^2 y + EI \partial_x^4 y = \Delta P(x)$$

Galerkin Method

Galerkin–Fourier expansion

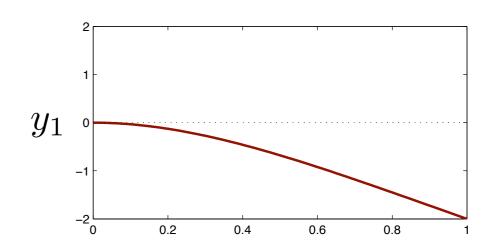
$$y(x,t) = \sum_{n} A_n y_n(x) e^{i\omega t}$$

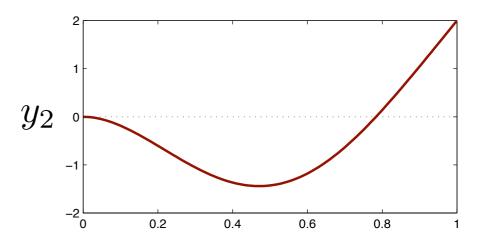
PDE — eigenvalue problem

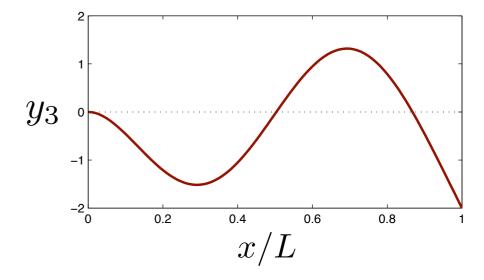
$$\rho \partial_t^2 y + EI \partial_x^4 y = \Delta P(x)$$



$$\left(-\rho\omega^2\underline{I} + EI\underline{K} - \underline{P}(\omega)\right)\underline{A} = 0$$

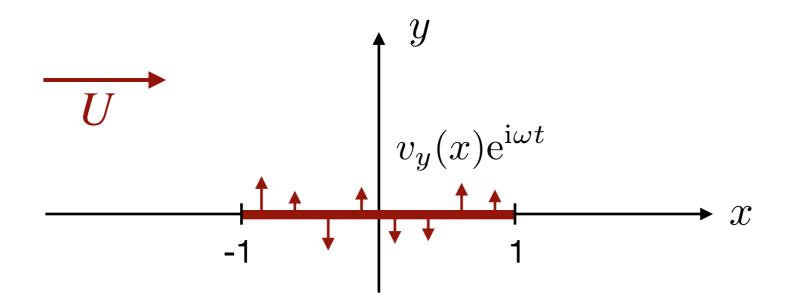






Flow Around the Plate

Potential Flow



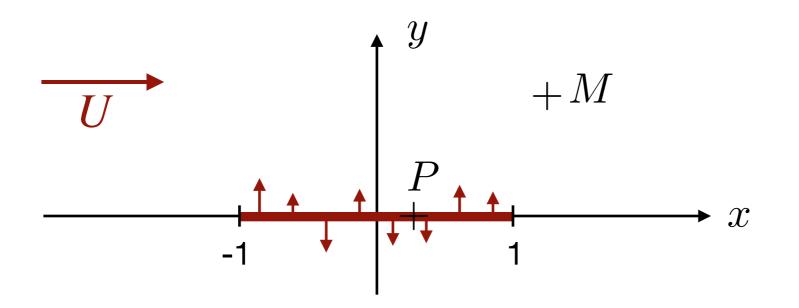
• Perturbation potential: $\phi(x,y)e^{i\omega t}$

$$\begin{cases} \Delta \phi = 0 \\ \partial_y \phi|_{y=0} = v_y = (\partial_t + U \partial_x) y_n \text{ for } x \in [-1 \quad 1] \end{cases}$$

Perturbation pressure (Bernoulli equation)

$$P = (\partial_t + U\partial_x) \phi e^{i\omega t}$$

Inverse Problem



$$\begin{cases} \Delta \phi = 0 \\ \partial_y \phi|_{y=0} = v_y \end{cases}$$

Green's representation theorem

$$\phi(M) = -\int_{P \in S} \delta \phi \, \partial_{y_p} G(|MP|) \mathrm{d}S_p$$

$$G(r) = \frac{1}{2\pi} \ln r \qquad \text{(in 2D)}$$

Inverse problem (Fredholm equation of 1st kind)

$$v_y(x) = \oint_{-1^-}^{1^+} \delta\phi(\xi) \frac{d\xi}{2\pi(x-\xi)^2}$$

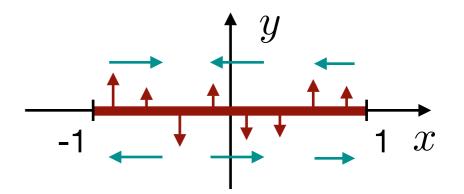
Vorticity Distribution

Symmetry of perturbation potential

$$\phi(-y) = -\phi(y)$$

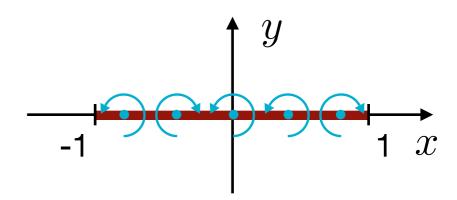
$$v_x = \partial_x \phi$$

$$v_x(-y) = -v_x(y)$$



Vorticity distribution

$$\gamma(x) = -\partial_x \delta \phi$$



Inverse problem for vorticity

$$v_y(x) = \oint_{-1^-}^{1^+} \delta\phi(\xi) \frac{d\xi}{2\pi(x-\xi)^2}$$

$$v_y(x) = \oint_{-1}^1 \gamma(\xi) \frac{\mathrm{d}\xi}{2\pi(x-\xi)}$$

Oscillating Plate

Inversion formula (Söhngen, 1939)

$$v_y(x) = \oint_{-1}^1 \gamma(\xi) \frac{\mathrm{d}\xi}{2\pi(x-\xi)}$$

$$\gamma(x) = -\frac{2}{\pi} \sqrt{\frac{1-x}{1+x}} \oint_{-1}^{1} \sqrt{\frac{1+\xi}{1-\xi}} \frac{v_y(\xi)}{x-\xi} d\xi + \frac{\alpha}{\sqrt{1-x^2}}$$

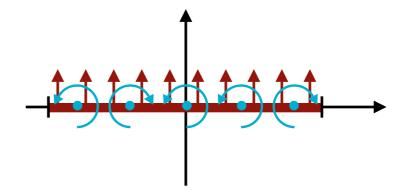
ullet Particular case: $v_y=1$

$$\Delta P(x) = -\gamma(x) - i\omega \int_{-1}^{x} \gamma(\xi) d\xi$$



Flow with no circulation

$$\Delta P(x) = \frac{2x}{\sqrt{1-x^2}} - 2i\omega\sqrt{1-x^2}$$

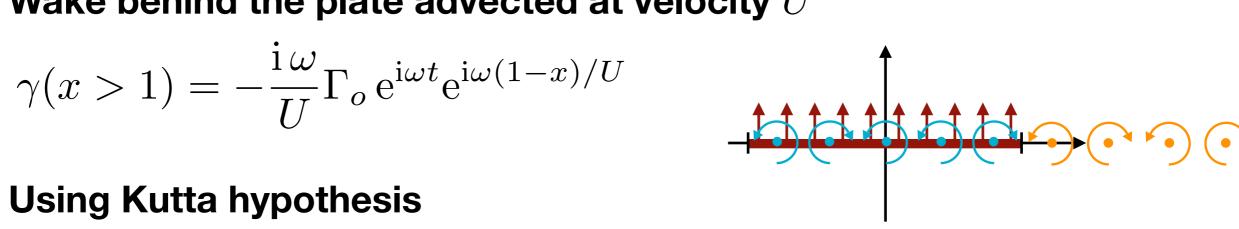


ullet Flow with circulation $\Gamma = \Gamma_o \mathrm{e}^{\mathrm{i}\omega t}$

$$\gamma|_{x=1^+}=-rac{1}{U}rac{\mathrm{d}\Gamma}{\mathrm{d}t}=-rac{\mathrm{i}\,\omega}{U}\Gamma_o\,\mathrm{e}^{\mathrm{i}\omega t}$$
 (Kelvin's circulation theorem)

ullet Wake behind the plate advected at velocity U

$$\gamma(x > 1) = -\frac{\mathrm{i}\,\omega}{U} \Gamma_o \,\mathrm{e}^{\mathrm{i}\omega t} \mathrm{e}^{\mathrm{i}\omega(1-x)/U}$$



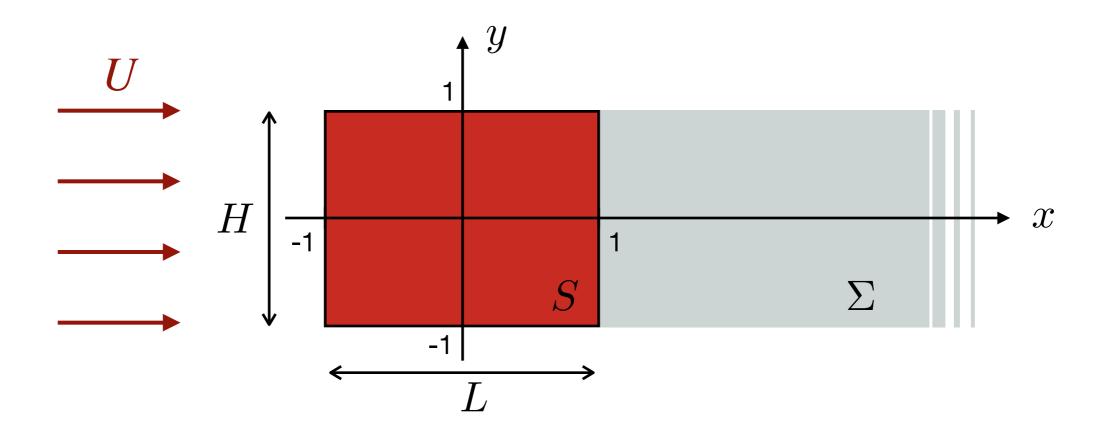
Using Kutta hypothesis

$$\Delta P(x) = 2 C(\omega) \sqrt{\frac{1-x}{1+x}} - 2 i \omega \sqrt{1-x^2}$$

$$C(\omega) = \frac{H_1^{(2)}(\omega)}{H_1^{(2)}(\omega) + \mathrm{i} H_0^{(2)}(\omega)} \tag{Theodorsen function}$$

3D Flow?

3D Problem



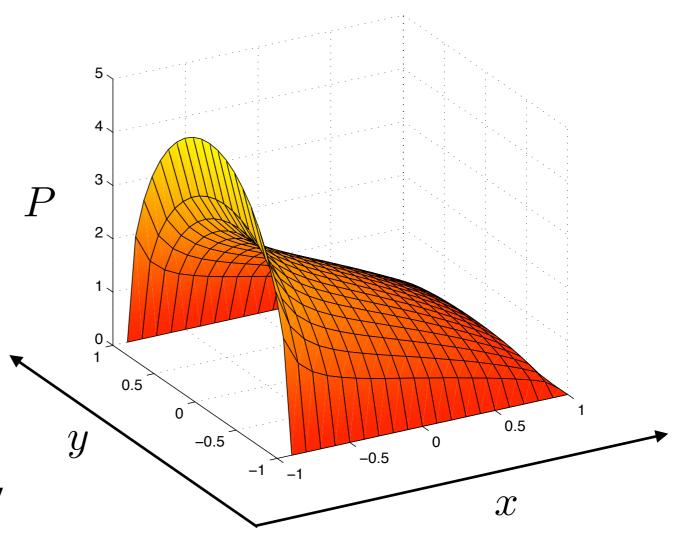
Inverse problem (Fredholm equation of 1st kind)

$$v_z(x,y) = \frac{A}{4\pi} \iint_{S+\Sigma} \frac{\gamma_y(\xi,\eta) \,d\xi d\eta}{(y-\eta)^2} \left(1 + \frac{A(x-\xi)}{[A^2(x-\xi)^2 + (y-\eta)^2]^{1/2}} \right)$$

$$A = L/H$$

• Expansion in powers of A or A^{-1} (lifting-line theory)

Singularities in the Pressure Field



Leading-edge singularity

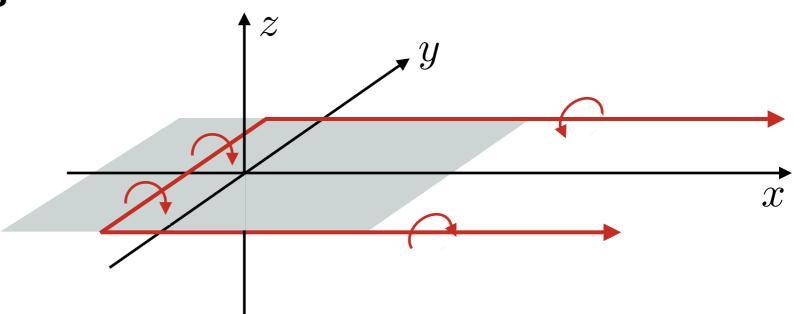
$$P \sim \frac{1}{\sqrt{n}}$$

Trailing-edge and side edges

$$P \sim \sqrt{n}$$

Vorticity Distribution

Vorticity lines

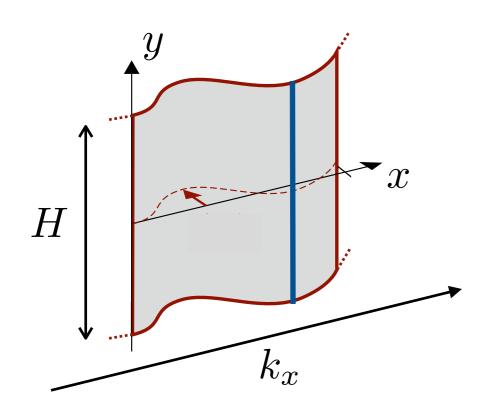


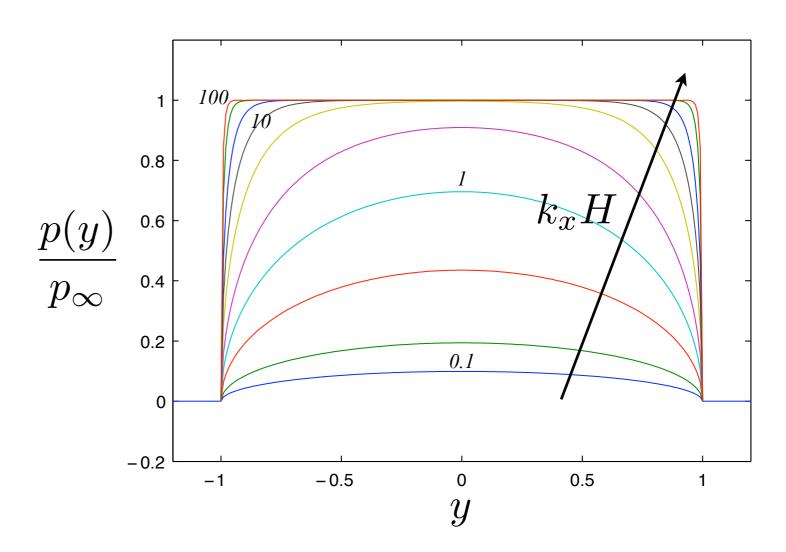
Vortex lattice method

$$v_i = \sum_j K_{ij} \Gamma_j$$
 x_j

[Ref] Tang, Yamamoto & Dowell, J. Fluids Struct. (2003)

Fourier Space





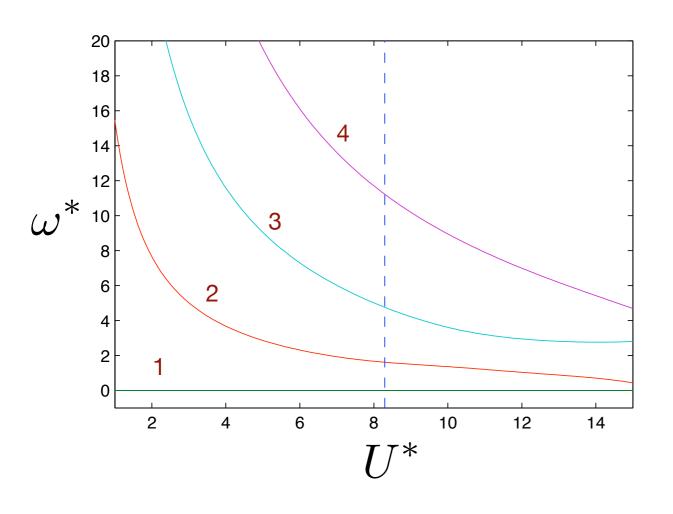
- ullet Assuming $v_z\sim {
 m e}^{{
 m i}k_xx}$
- ullet The pressure along $y \hspace{0.1cm}$ can be determined

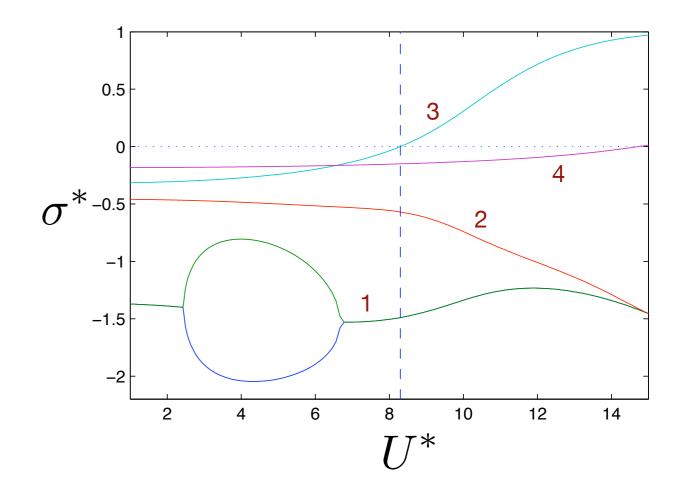
$$\langle p(y) \rangle = p_{\infty} \left(1 - \frac{1}{k_x H} + \mathcal{O}\left(e^{-k_x H}\right) \right)$$

Results of the Stability Analysis

Flutter Modes

• Parameters: $H^* = \infty$ $M^* = 2.5$

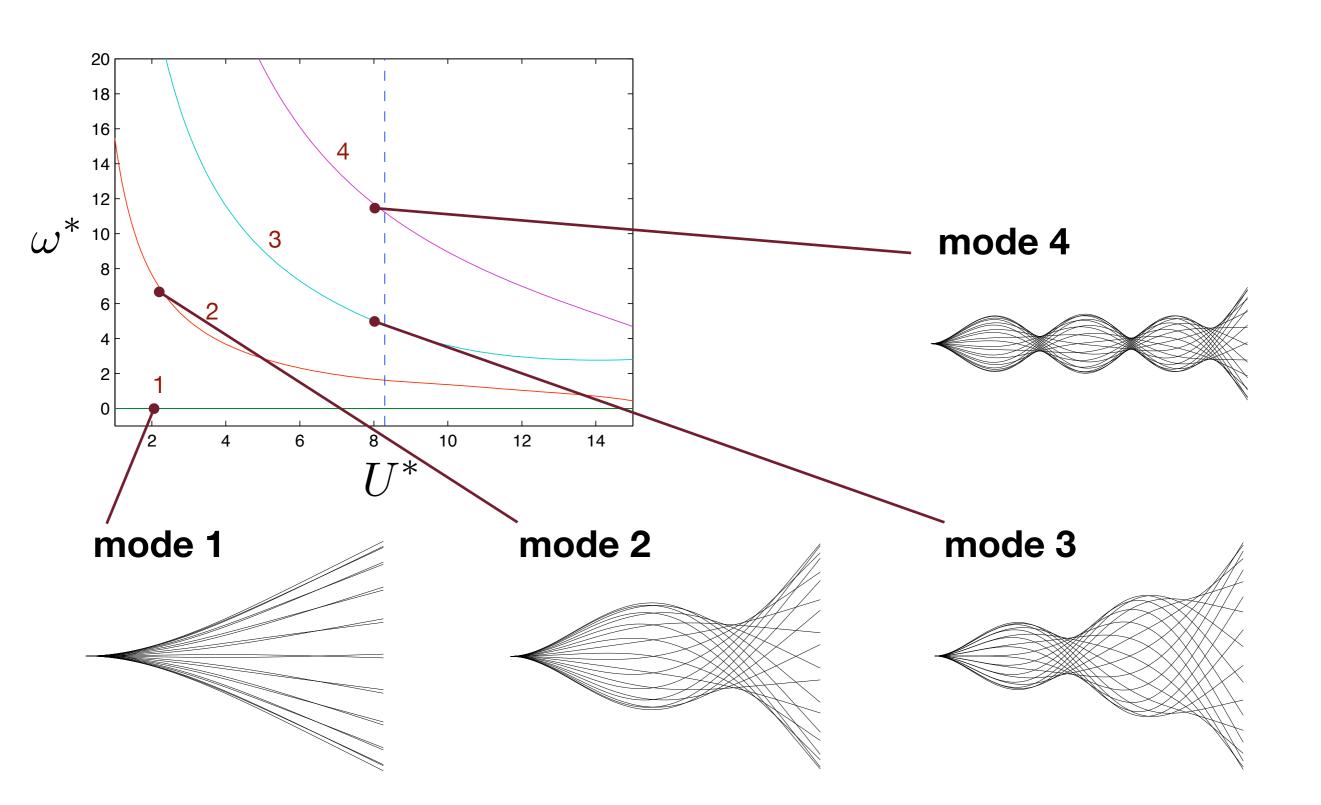




Dimensionless numbers

$$U^* = \sqrt{\frac{\rho}{EI}}LU$$
 $M^* = \frac{\rho_{air}L}{\rho}$ $H^* = \frac{H}{L}$ $\omega^* = \frac{L\omega}{U}$

Shape of Flutter Modes



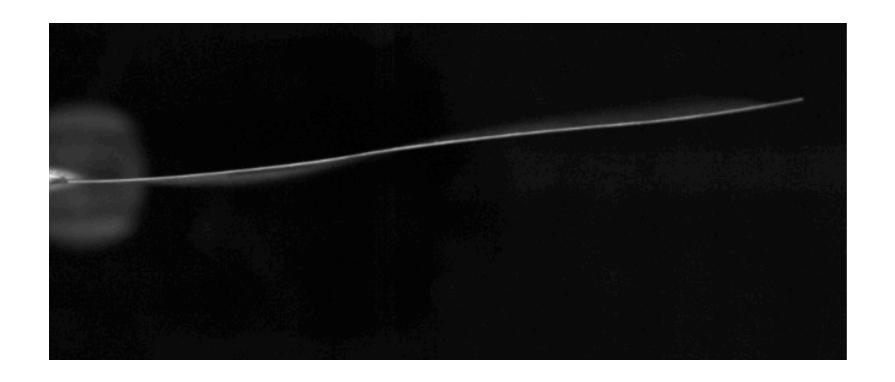
Unstable modes (50x slower)



mode 2

$$M^* = 0.74$$

$$L = 8 \,\mathrm{cm}$$

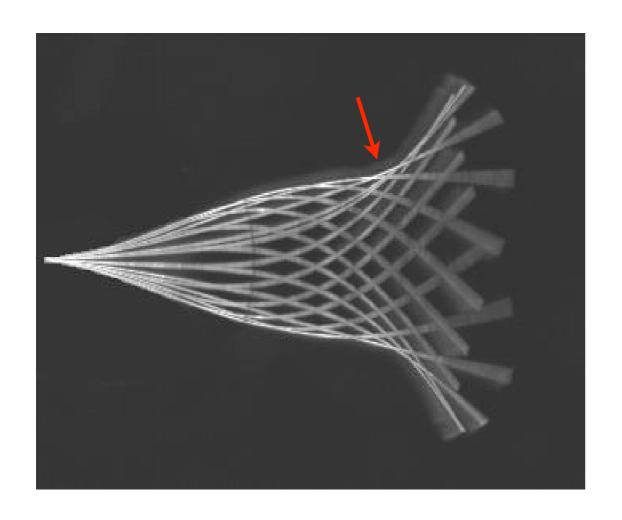


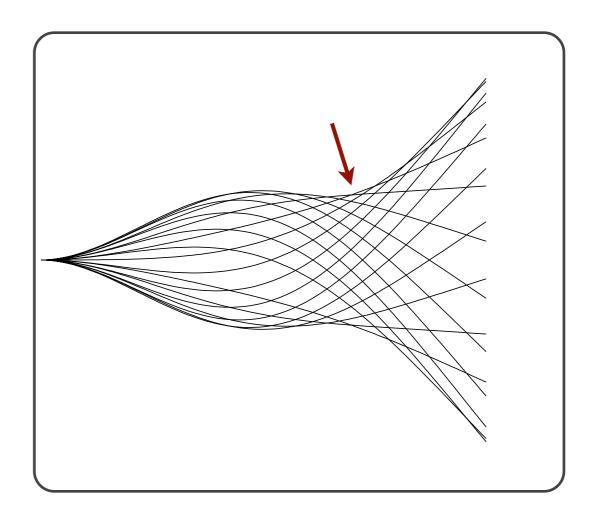
mode 3

$$M^* = 1.94$$

$$L = 21 \, \mathrm{cm}$$

Mode 2

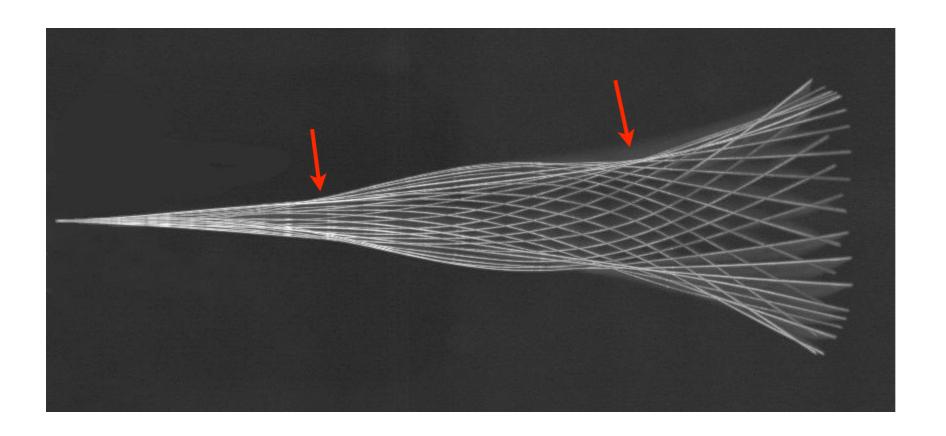




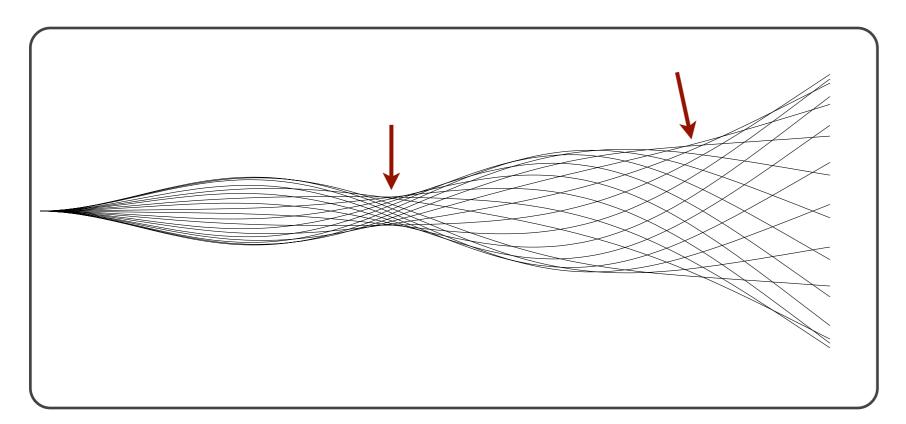
$$\omega^* = 1.82$$

$$\omega^* = 2.01$$

Mode 3

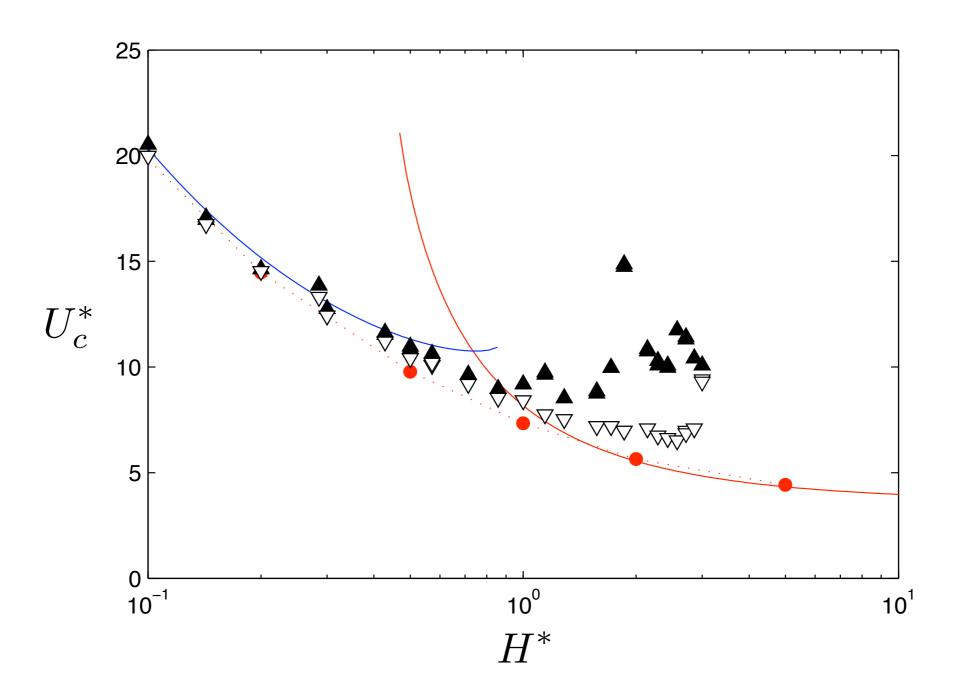


$$\omega^* = 2.5$$



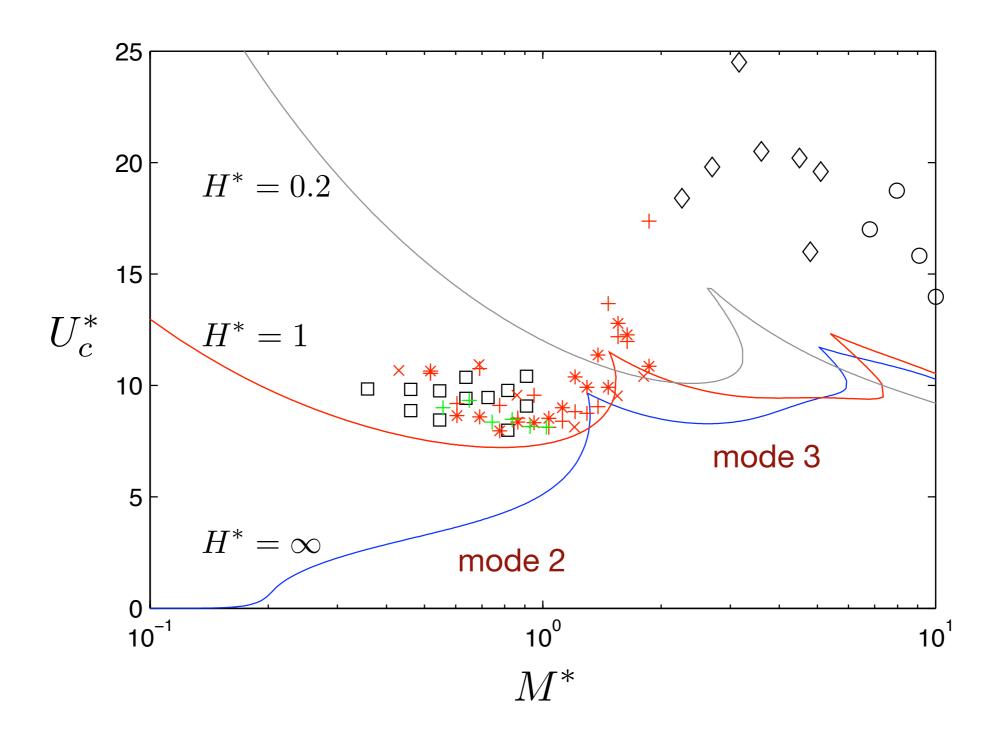
$$\omega^* = 3.5$$

Stability curve



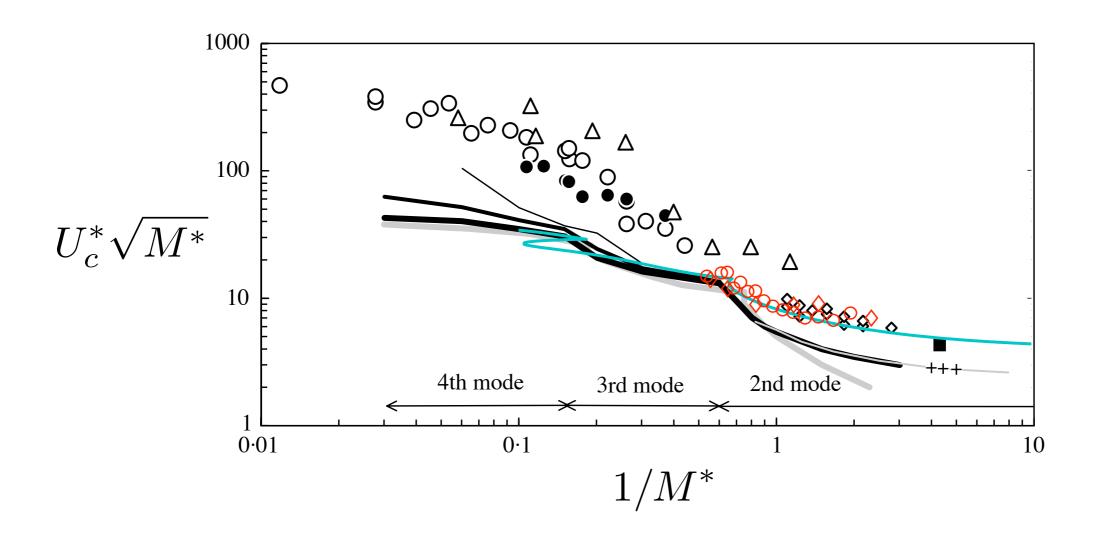
ullet Mass ratio: $M^{st}=0.6$

Stability curve



 $_{\bullet}$ Experiments for $H^{\ast}=1~$ and $H^{\ast}=0.25$

Conclusion



- Flow inherently singular
- Good agreement for 1D mode + 3D flow
- Importance of 3D effects