

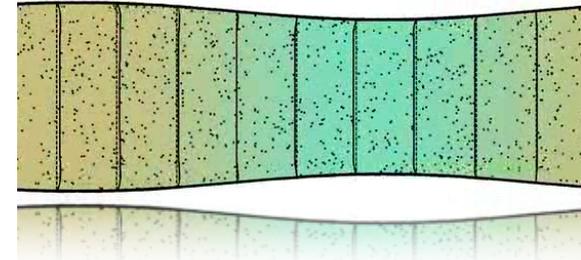
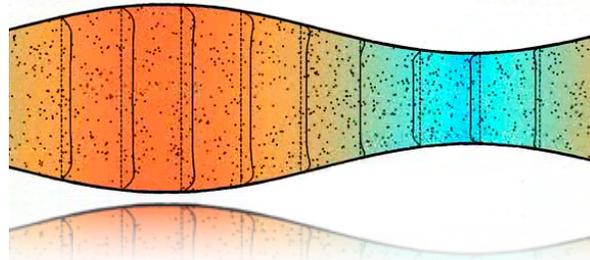
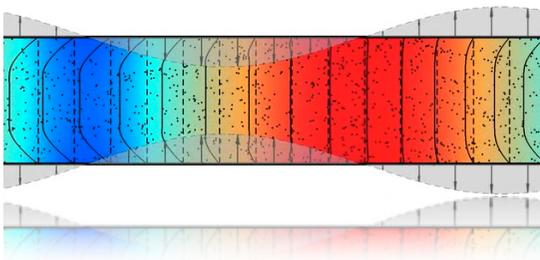
# Pumping from the walls

“La distinction des mécanismes”

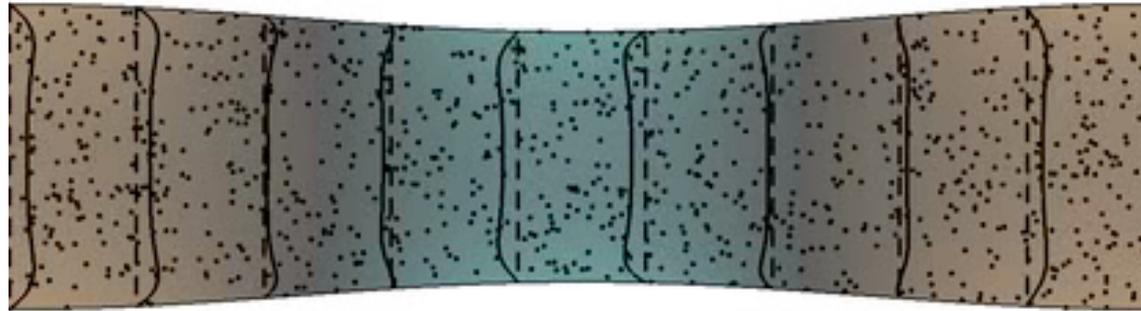
Jérôme Hoëffner

Koji Fukagata

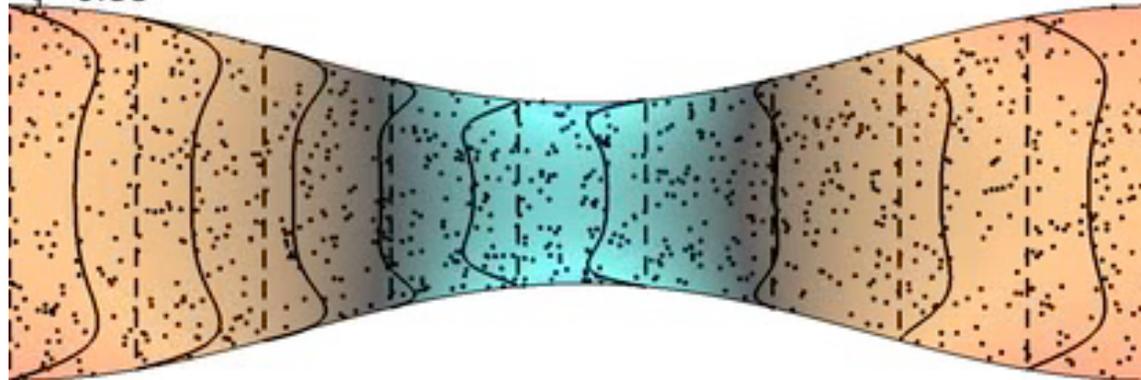
Keio University, Yokohama, Japan.  
Institut Jean le Rond D'Alembert, Paris, France.



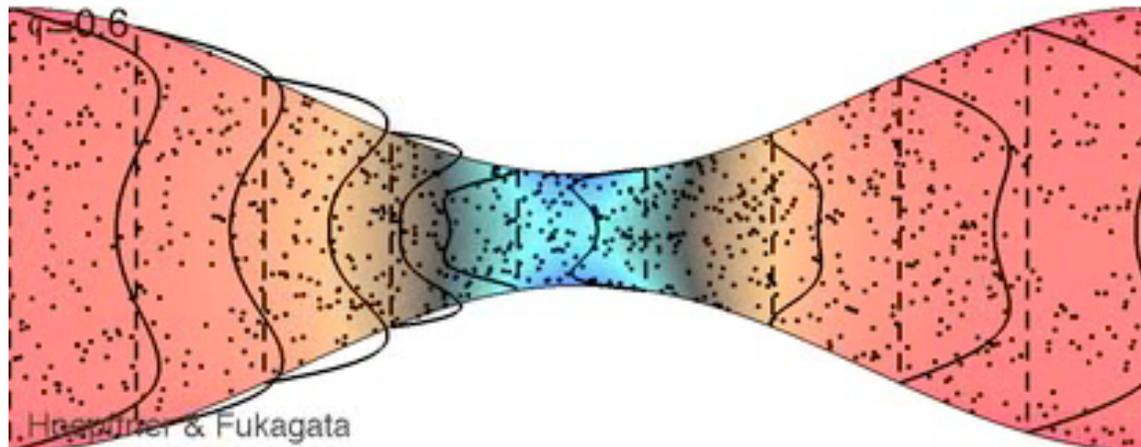
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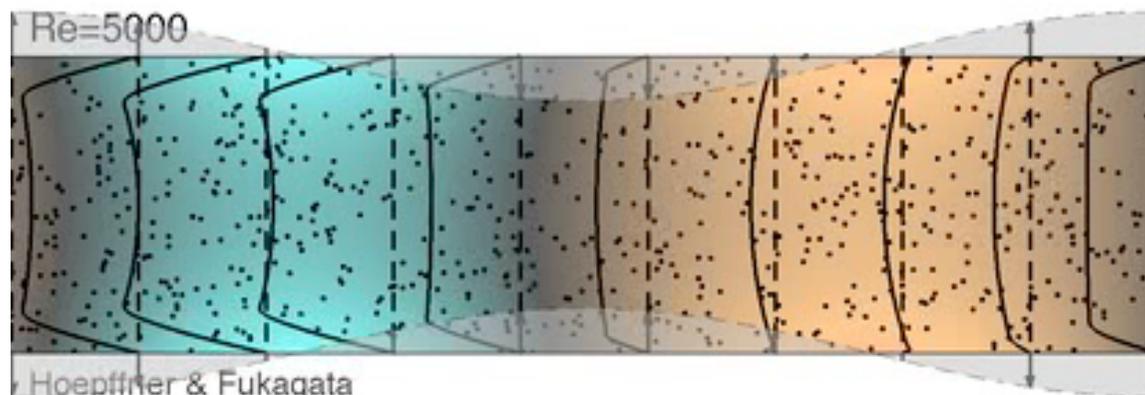
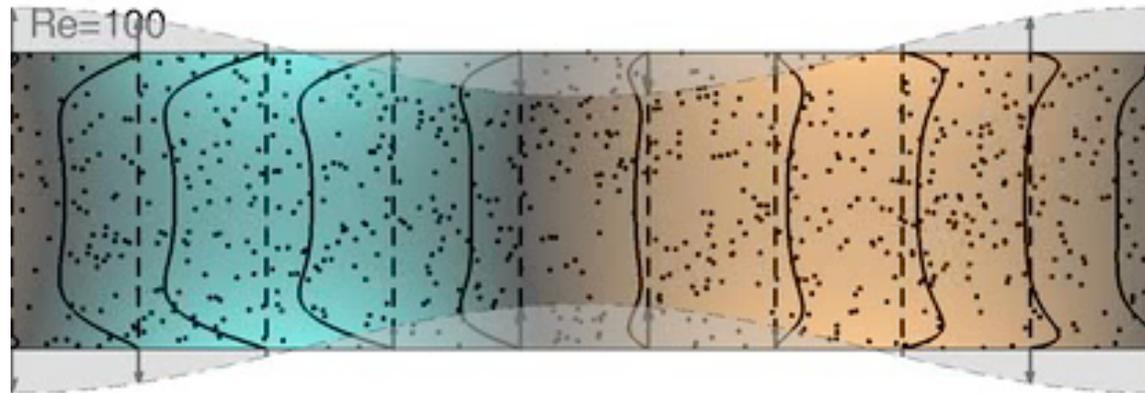
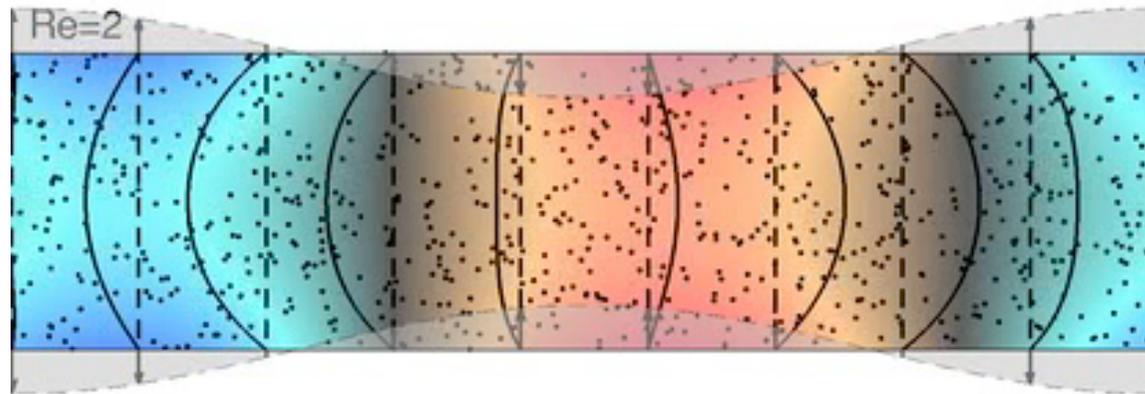
$q=0.35$



$q=0.6$

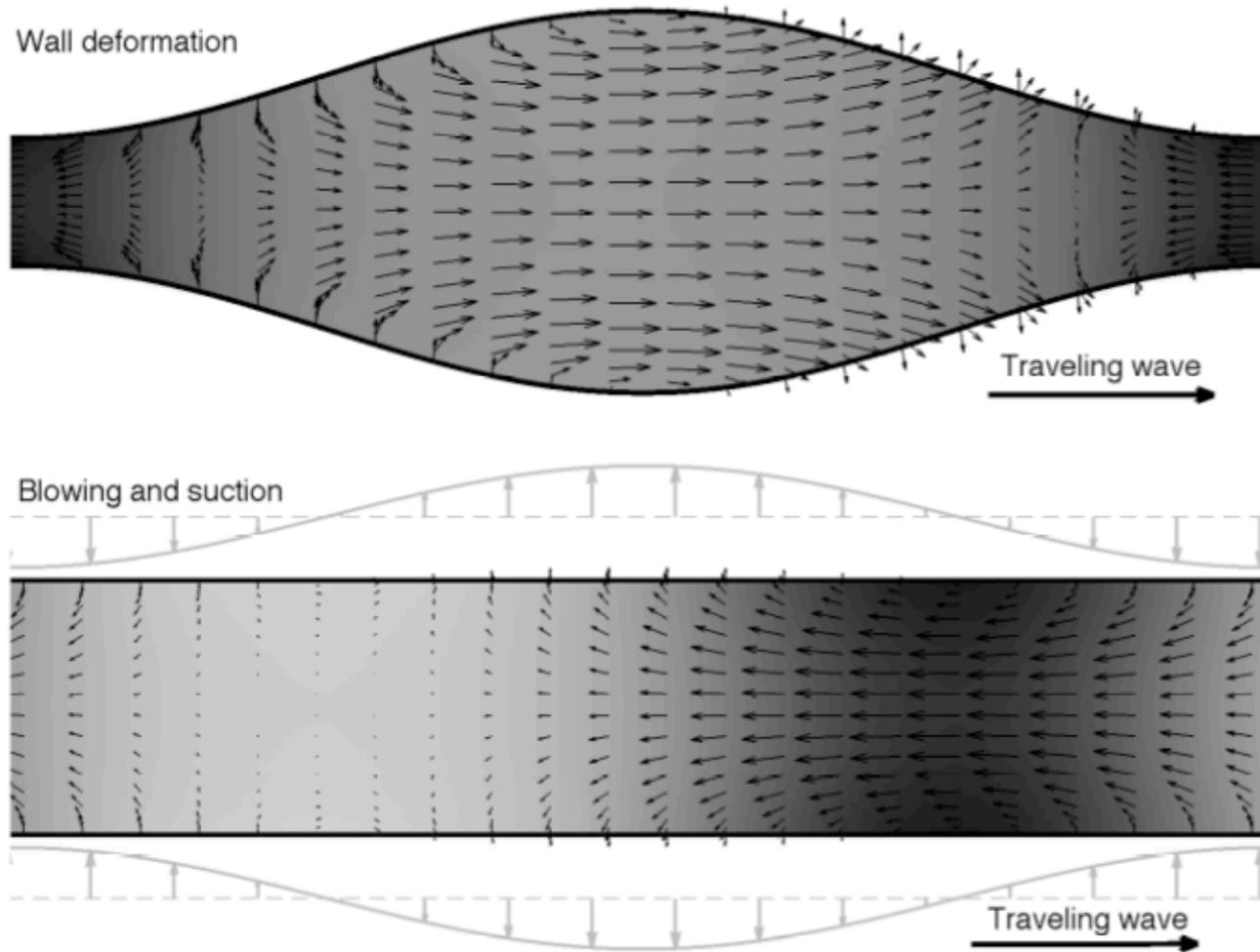


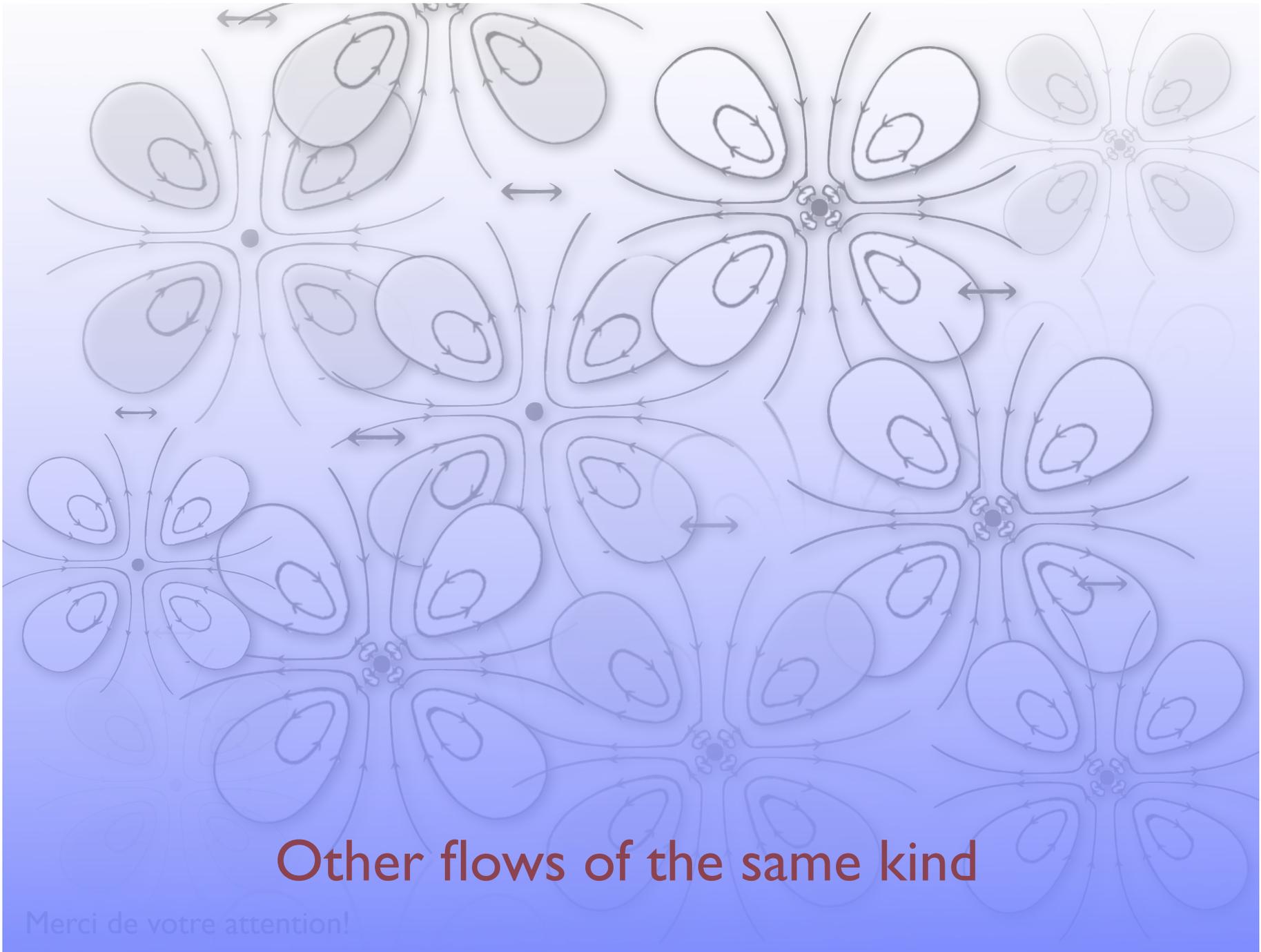
Hosono & Fukagata



Hoepffner & Fukagata

# Pumping in two directions

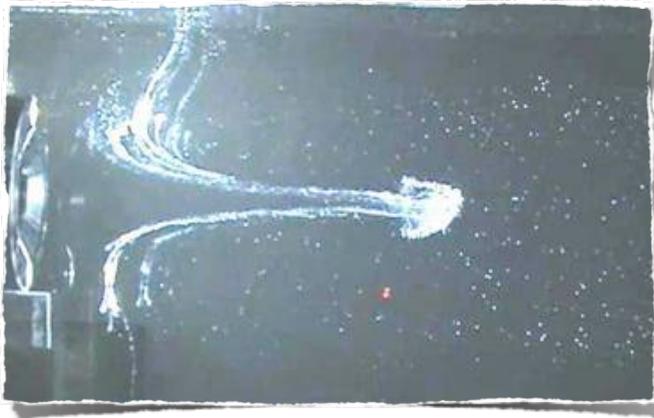




**Other flows of the same kind**

Merci de votre attention!

# “Quartz wind”

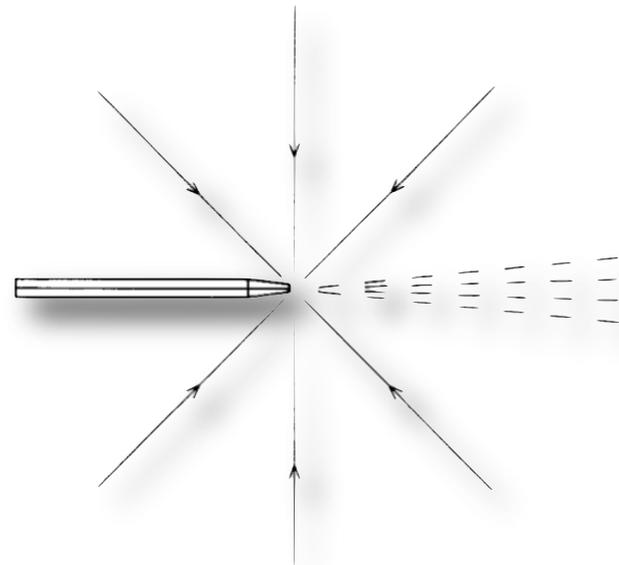


Jet generated by an ultrasonic beam

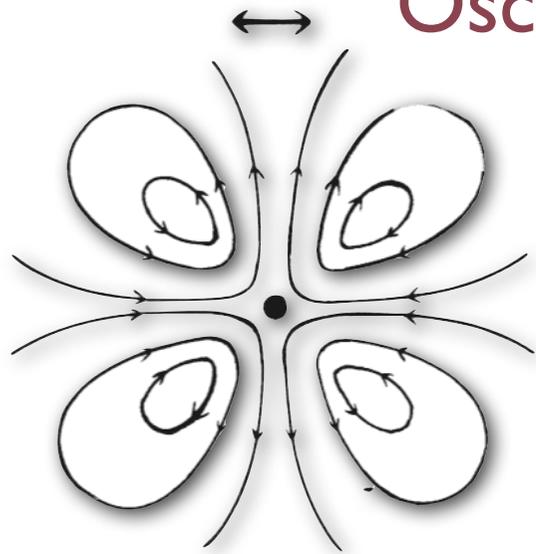
[http://www.lmfa.ec-lyon.fr/perso/Valery.Botton/acoustic\\_streaming\\_bis.html](http://www.lmfa.ec-lyon.fr/perso/Valery.Botton/acoustic_streaming_bis.html)

*Journal of Sound and Vibration* (1978) **61**(3), 391–418  
ACOUSTIC STREAMING†

SIR JAMES LIGHTHILL



## Oscillating cylinder

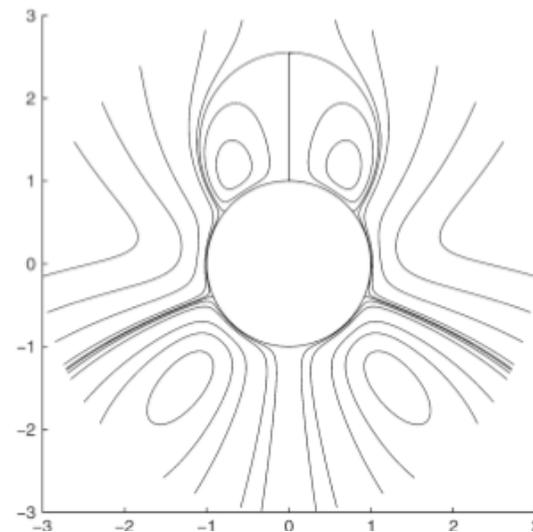


### Flow near an oscillating cylinder in dilute viscoelastic fluid

THERE are many natural phenomena in which nonlinear interactions of time-dependent inputs give rise to steady—that is, time-independent outputs. One of these is a steady streaming belonging to a class of secondary flows sometimes called acoustic streaming. It occurs when a circular cylinder oscillates normal to its axis in an unbounded Newtonian fluid<sup>1-3</sup>. We report here on the steady secondary flow induced when a long thin cylinder oscillates as described in a viscoelastic liquid. We found that the direction of steady streaming is opposite to that found for the bulk of fluid when the experiment is performed with a Newtonian fluid.

CHINGFENG CHANG  
W. R. SCHOWALTER

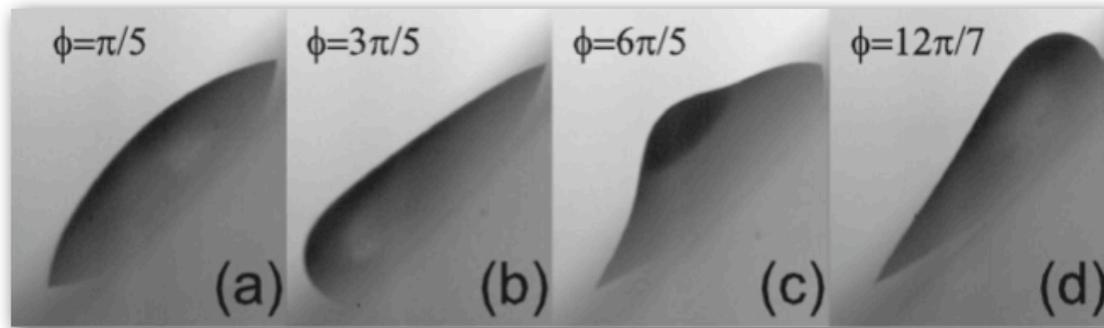
*Nature* Vol. 252 December 20/27 1974



STEADY STREAMING AROUND A SPHERICAL  
DROP DISPLACED FROM THE VELOCITY  
ANTINODE IN AN ACOUSTIC LEVITATION FIELD

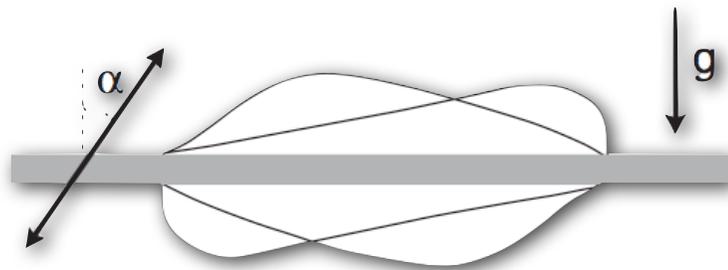
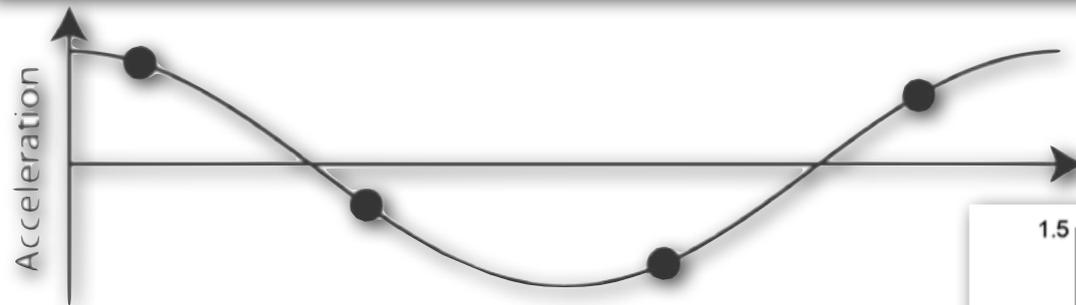
by A. Y. REDNIKOV, HONG ZHAO, S. S. SADHAL†

# Climbing drop

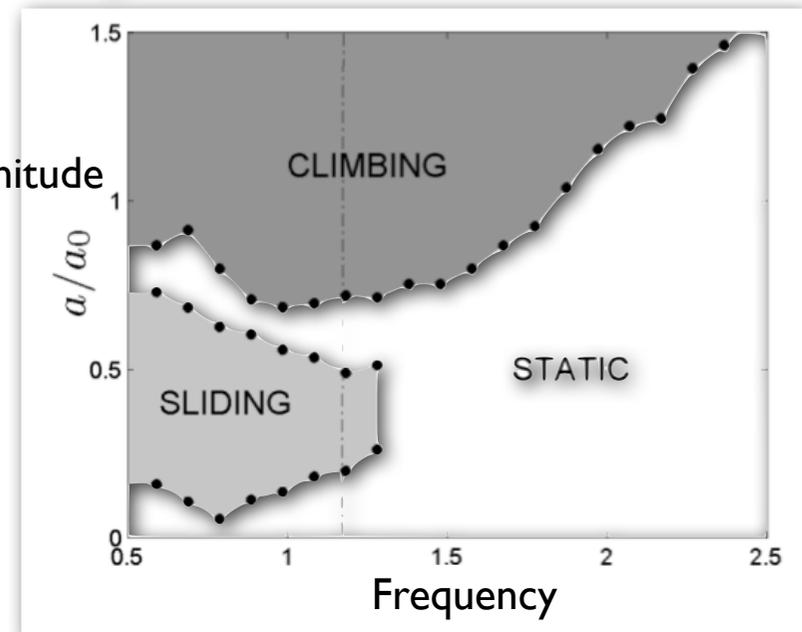


## Vibration-Induced Climbing of Drops

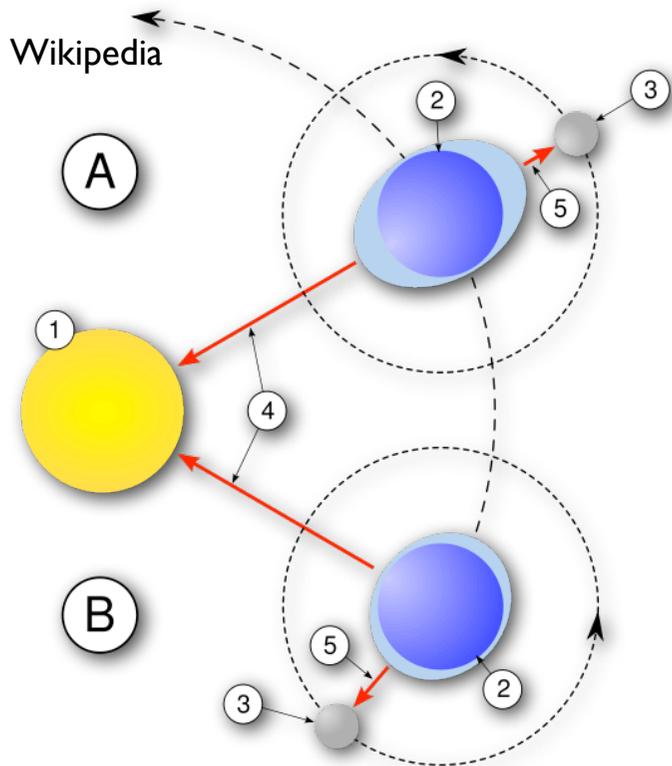
P. Brunet,<sup>\*</sup> J. Eggers, and R. D. Deegan  
PRL **99**, 144501 (2007)



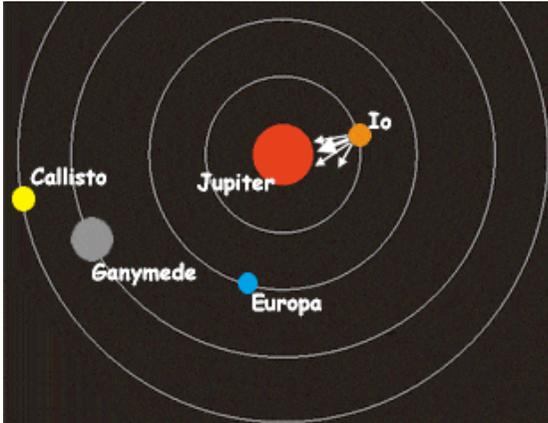
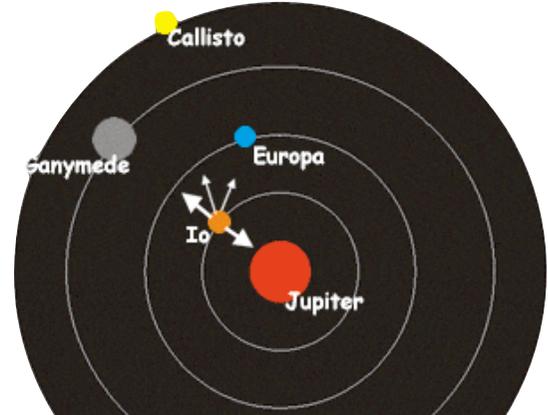
Magnitude



# Tidal forcing

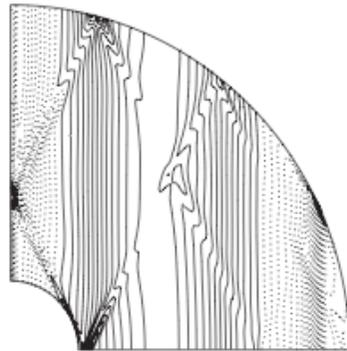


Io: satellite of Jupiter



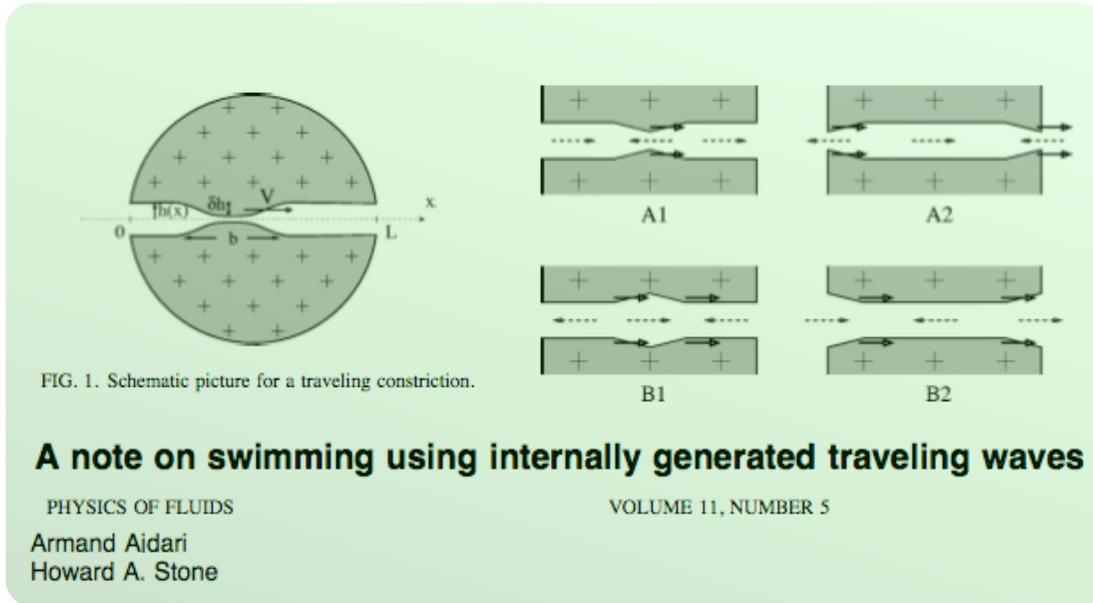
## Zonal Wind Driven by Inertial Modes

A. Tilgner  
 PRL 99, 194501 (2007)



[http://spaceplace.nasa.gov/en/kids/gll\\_io\\_fact.shtml](http://spaceplace.nasa.gov/en/kids/gll_io_fact.shtml)

# Life



## Microorganisms

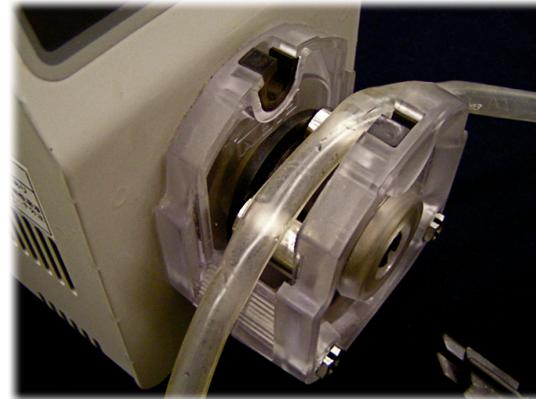
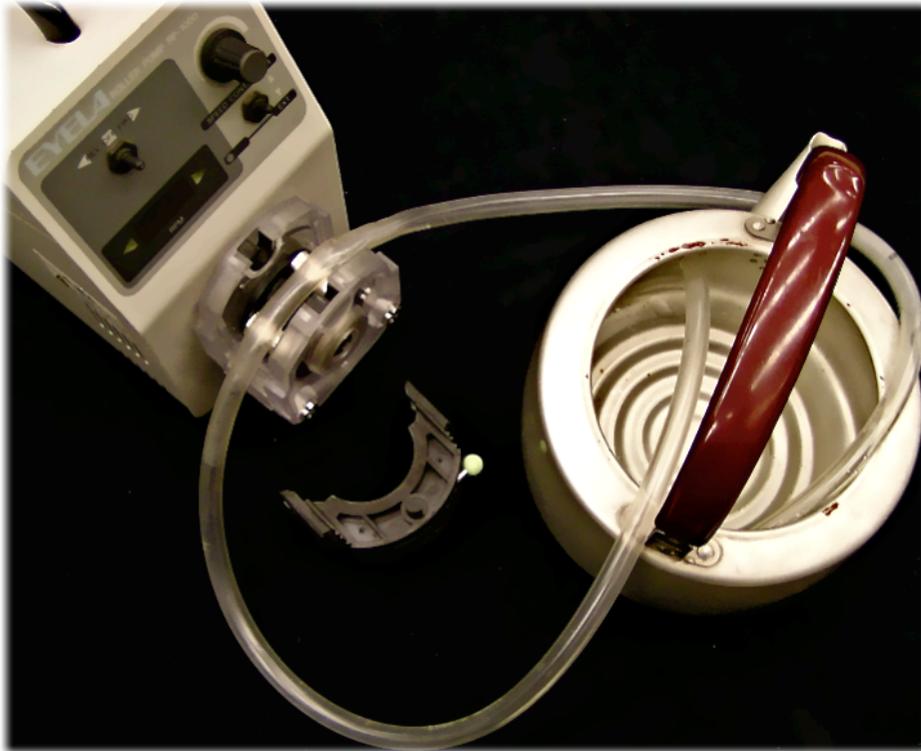
Image: youtube

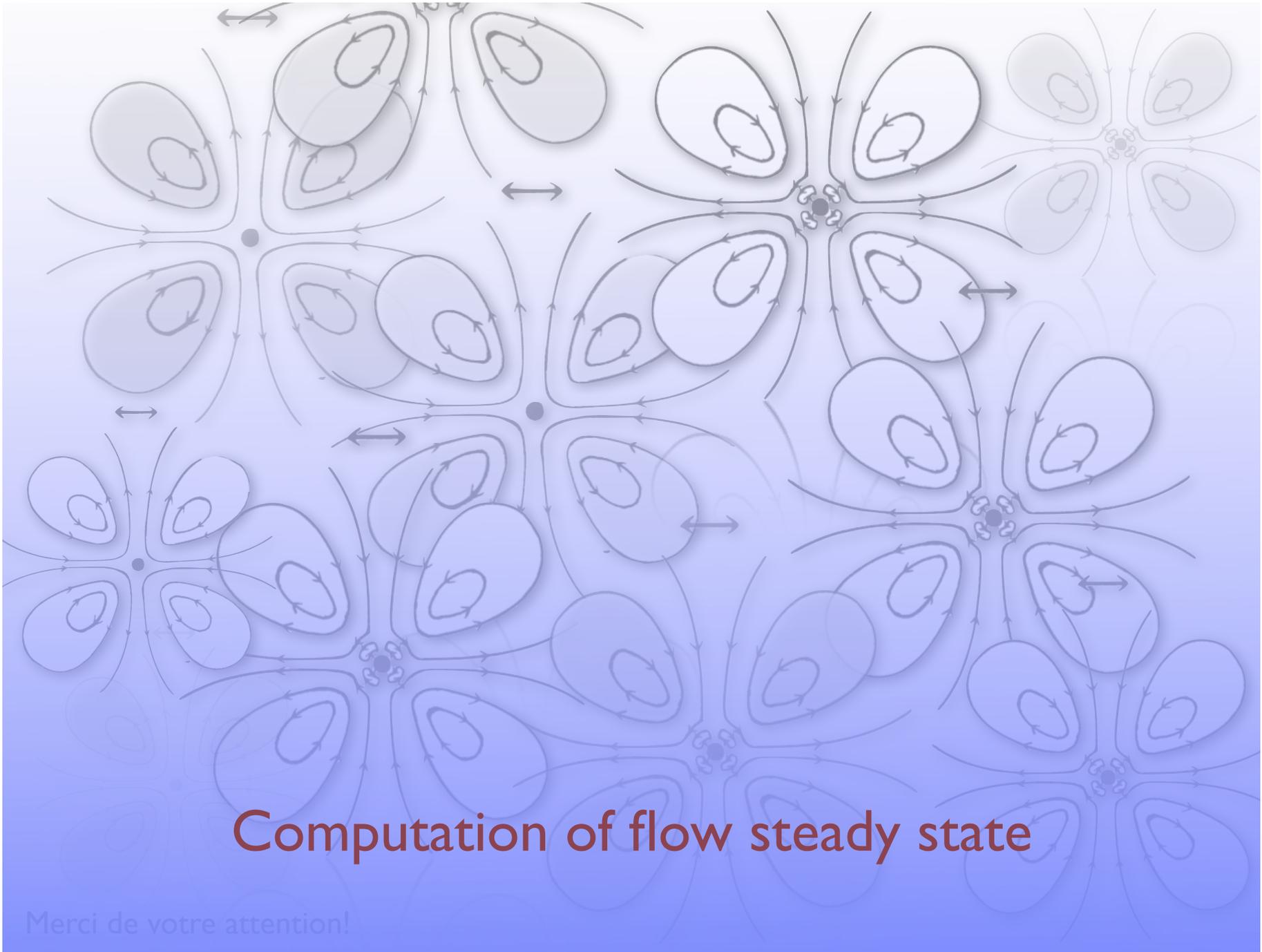


## Gastro-intestinal tract

Pumping in the gastro-intestinal tract, and in the ureter

# Industry





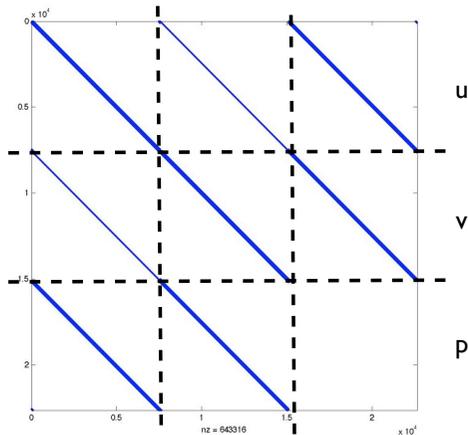
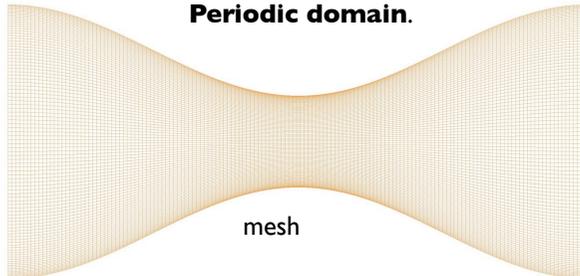
## Computation of flow steady state

Merci de votre attention!

# Tools: steady state

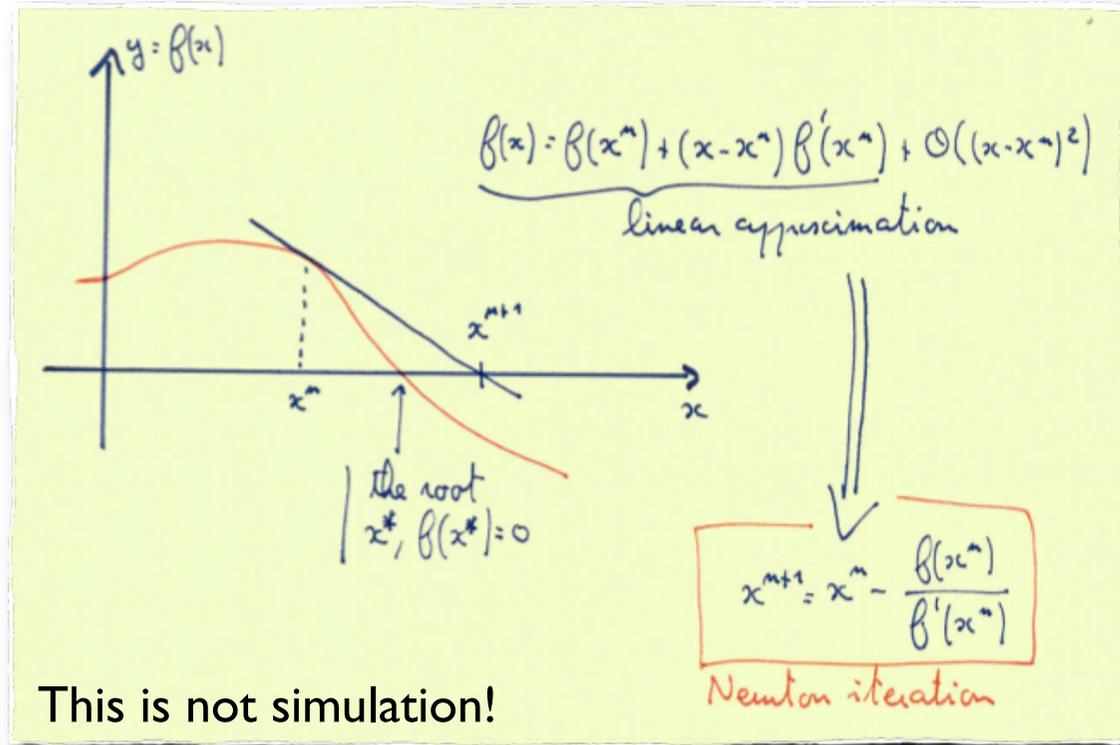
## 1) Nonlinear: steady state

Flow is steady in frame travelling with the wave.  
**Periodic domain.**

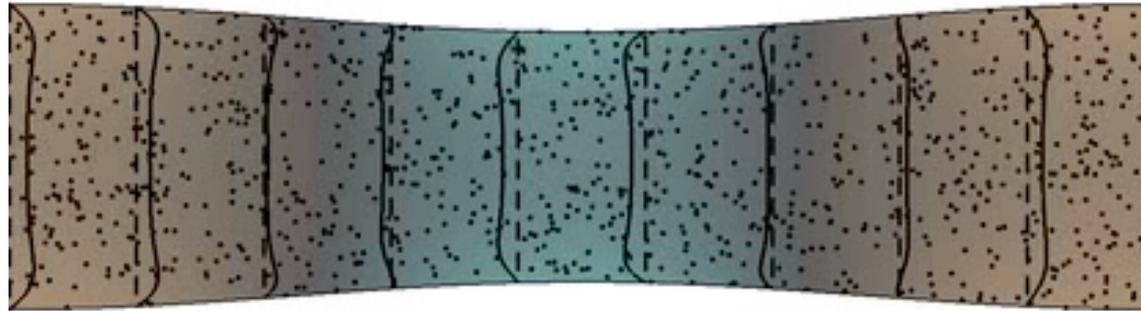


Compute steady state using Newton iterations.  
Here the structure of the **Jacobian**

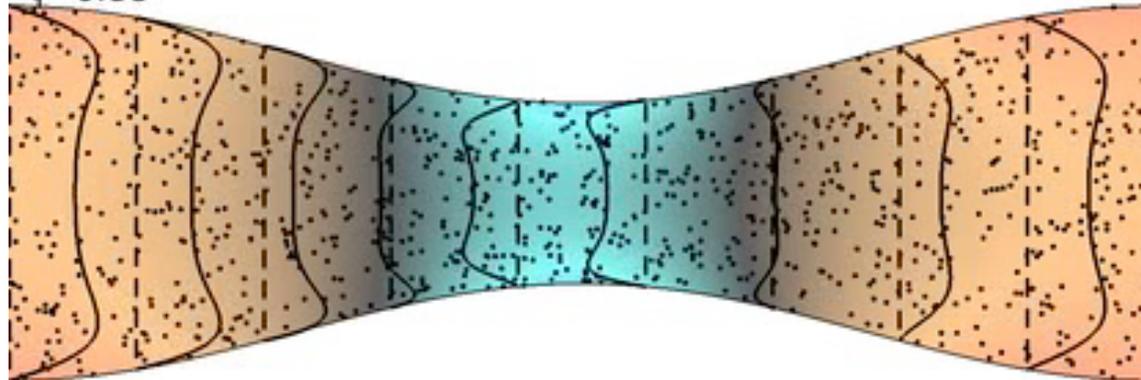
Finite difference/finite difference.  
**Sparse** matrices



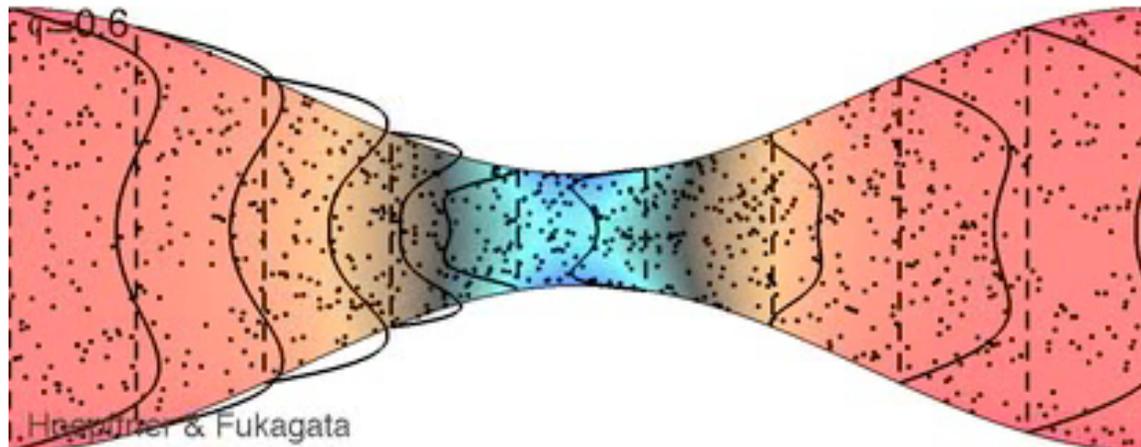
$q=0.1$



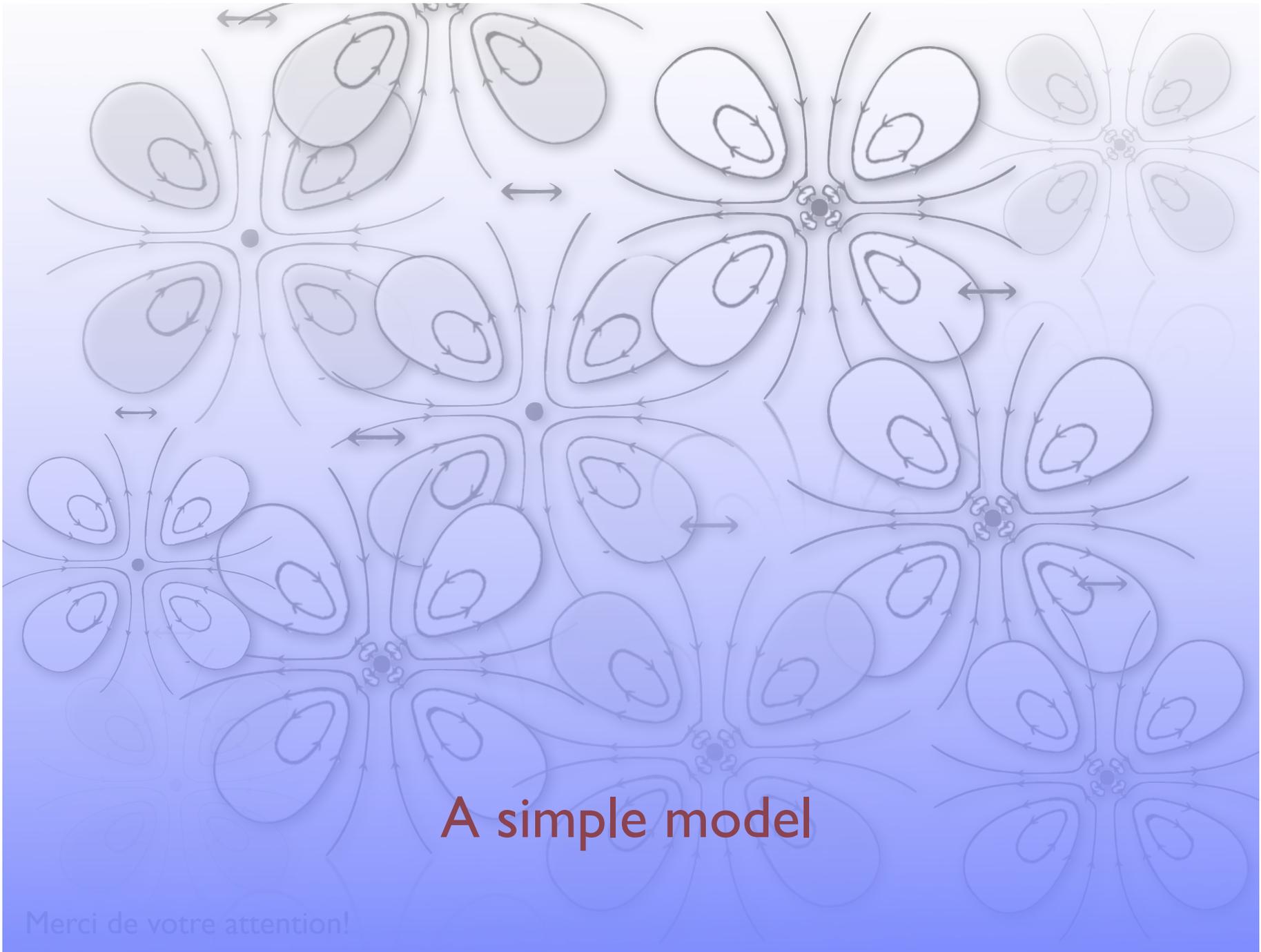
$q=0.35$



$q=0.6$



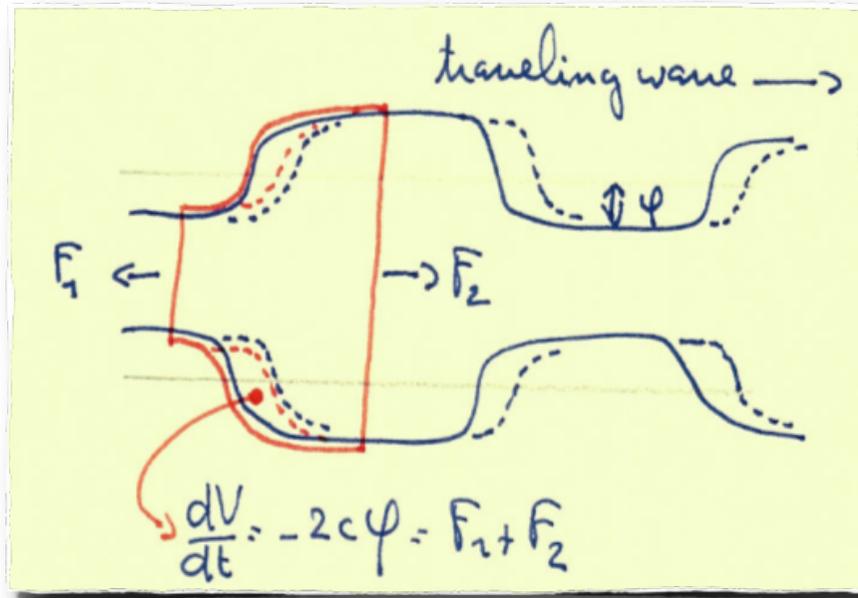
Hosokawa & Fukagata



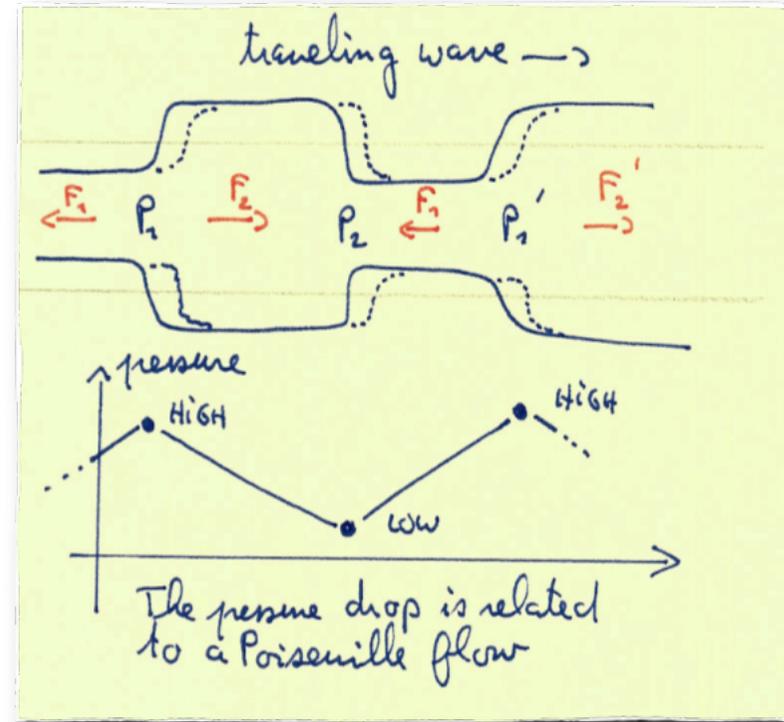
## A simple model

Merci de votre attention!

# Conservation model



Mass conservation

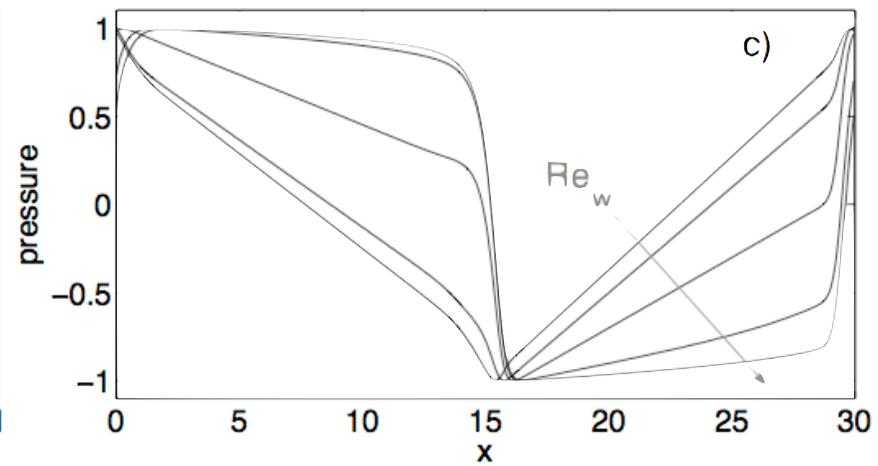
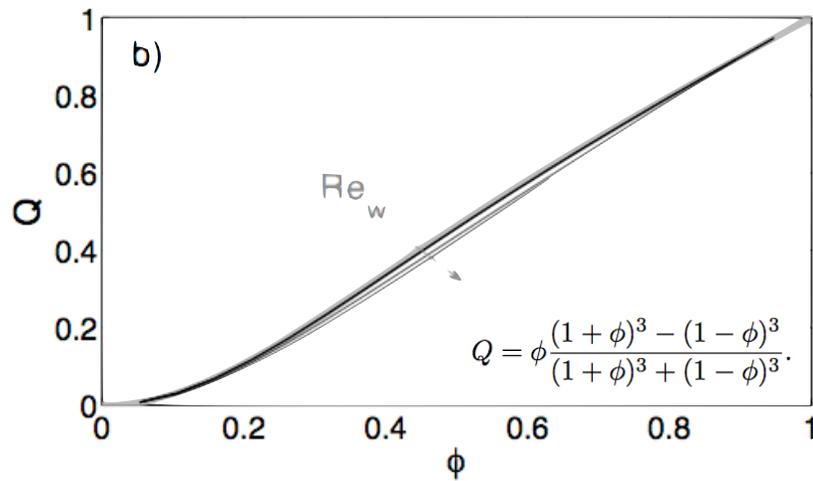
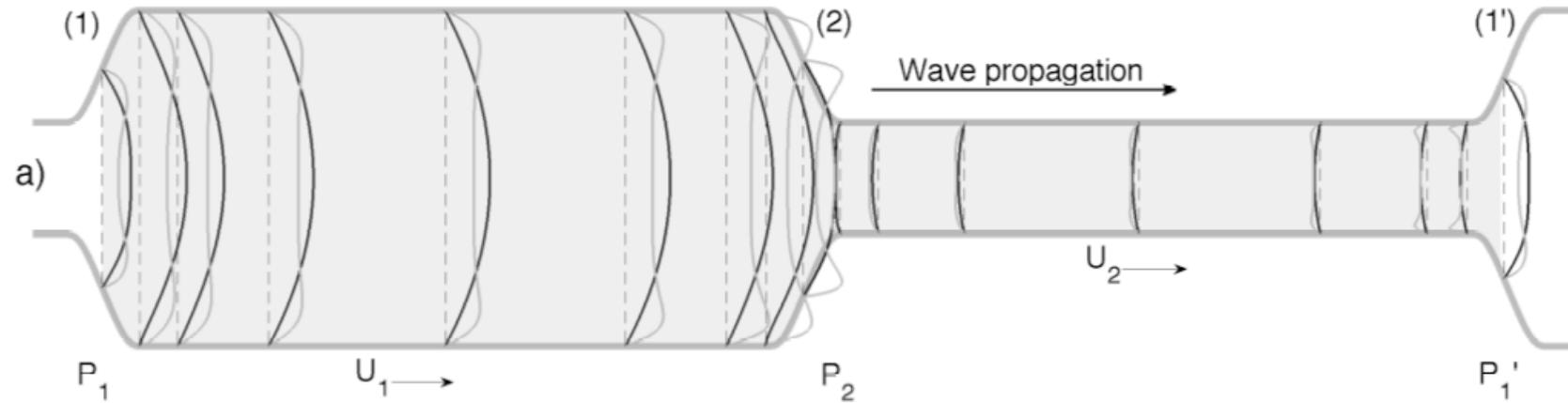


Pressure drop

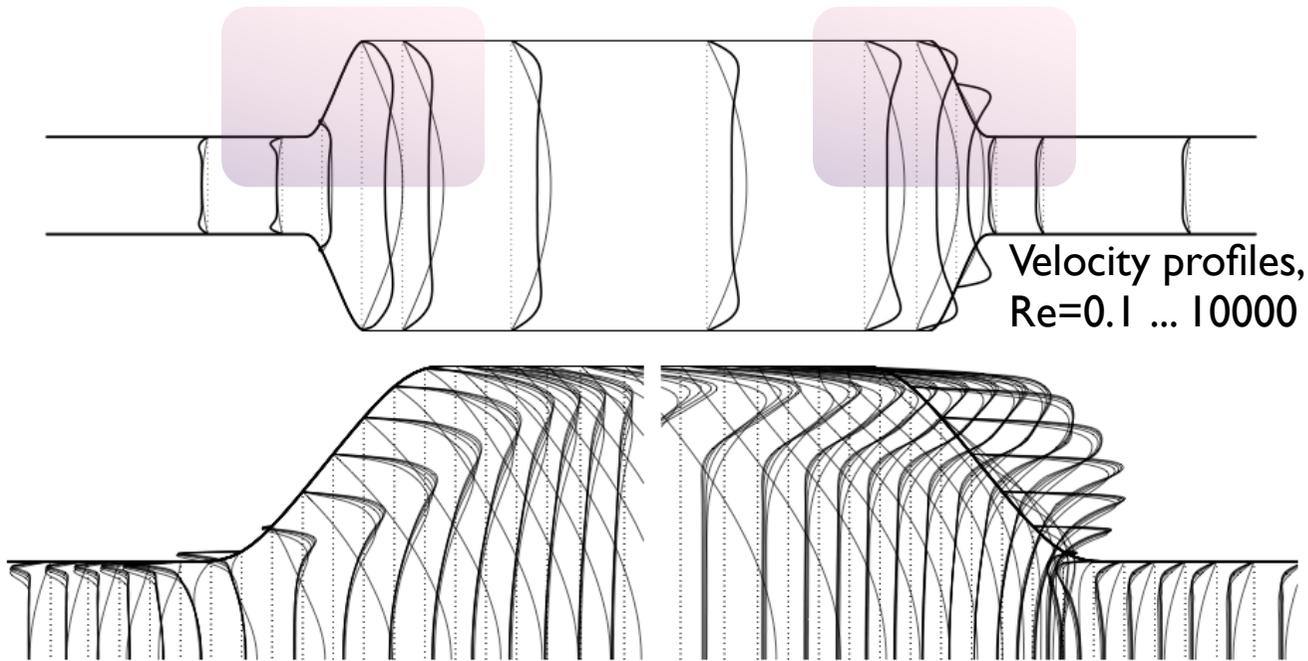
$$Q = \phi \frac{(1 + \phi)^3 - (1 - \phi)^3}{(1 + \phi)^3 + (1 - \phi)^3}$$

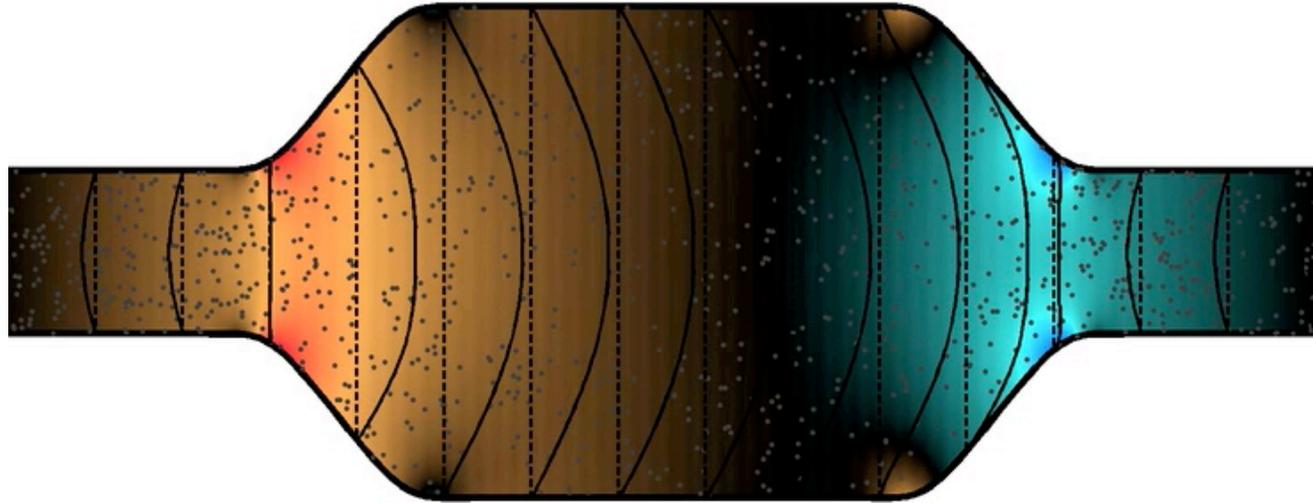
No Reynolds number effect!

# Pumping mechanism

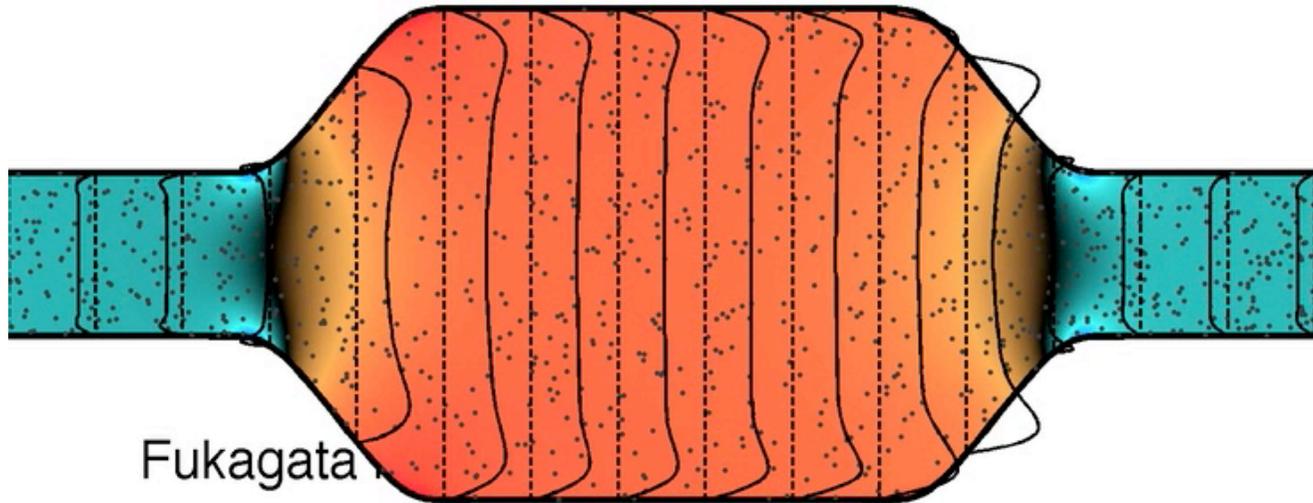


# Velocity profiles



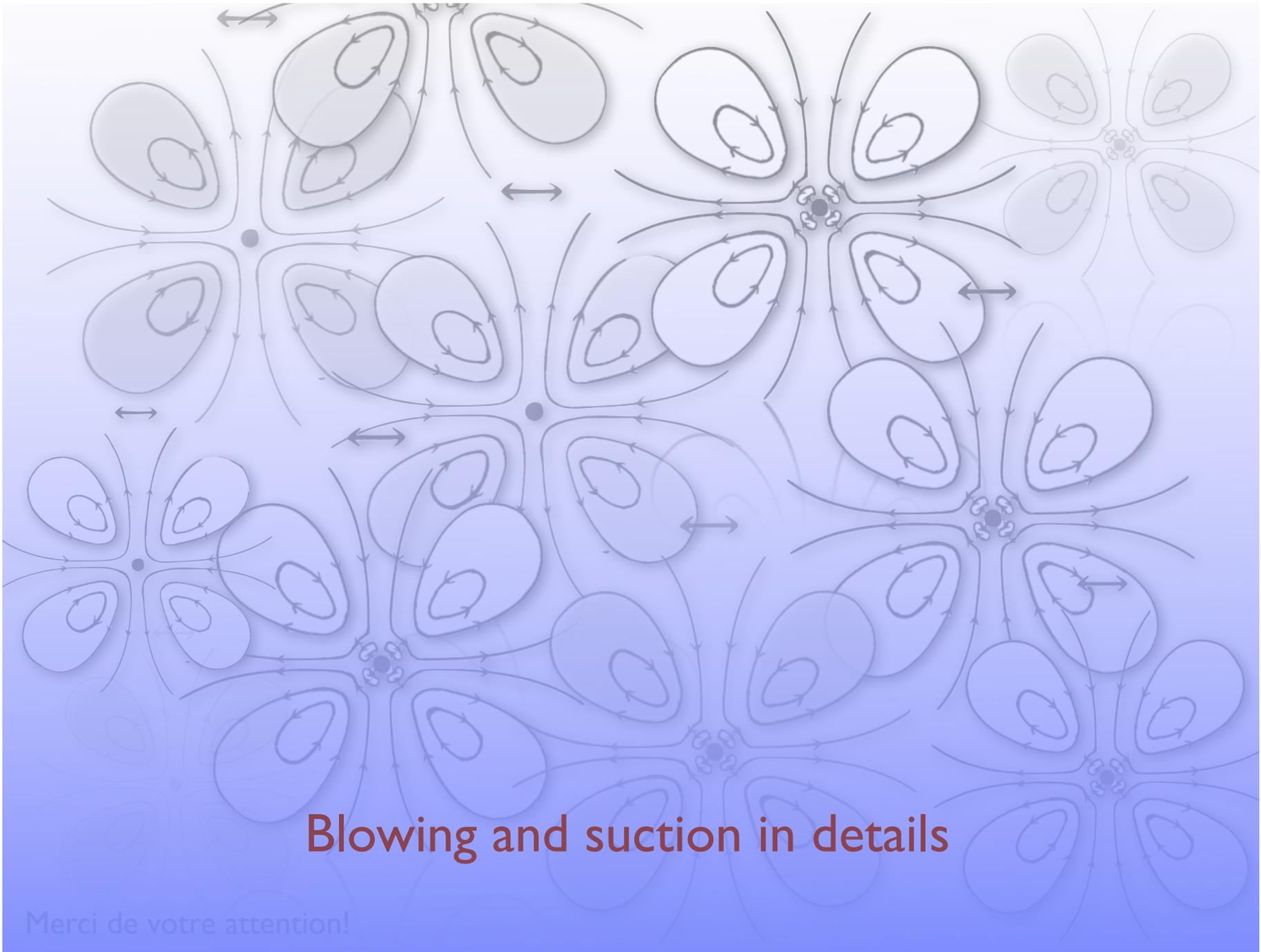


Re=1



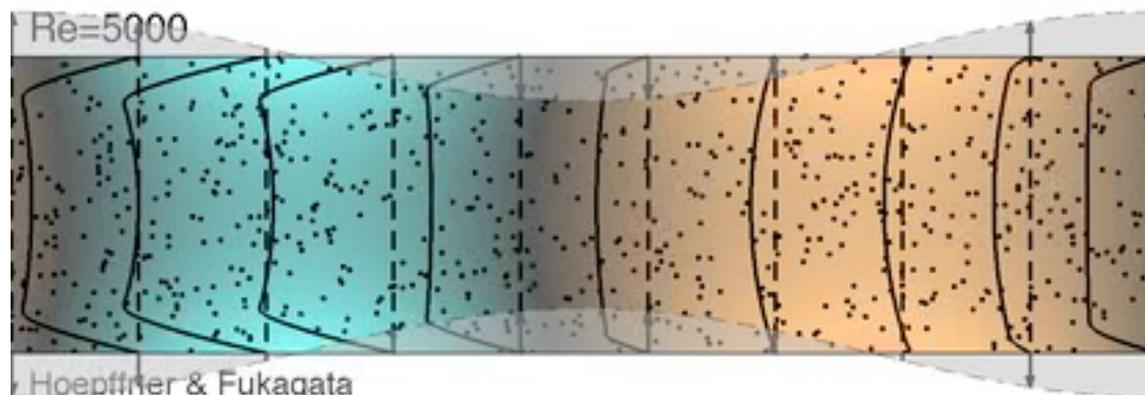
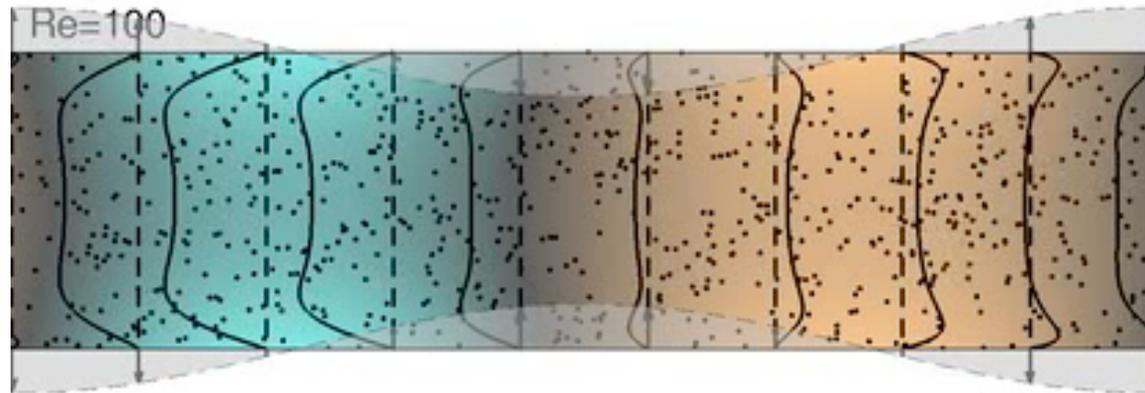
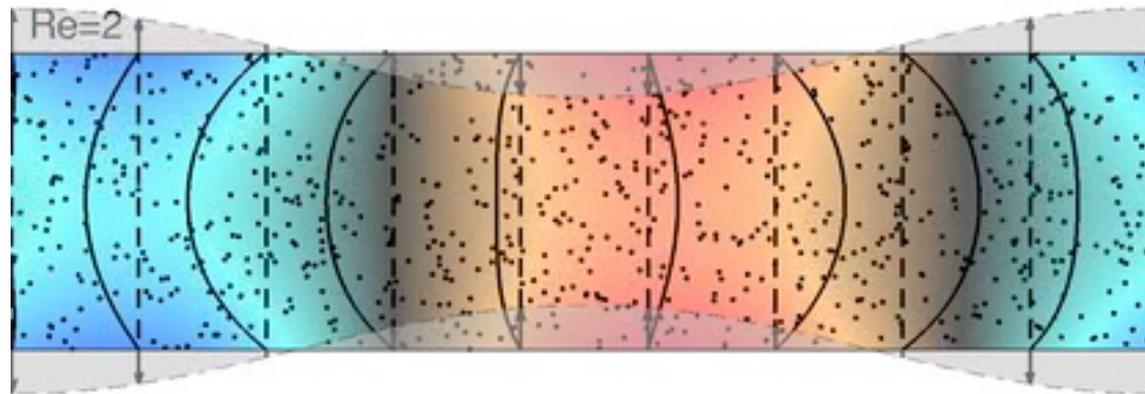
Re=1000

Fukagata



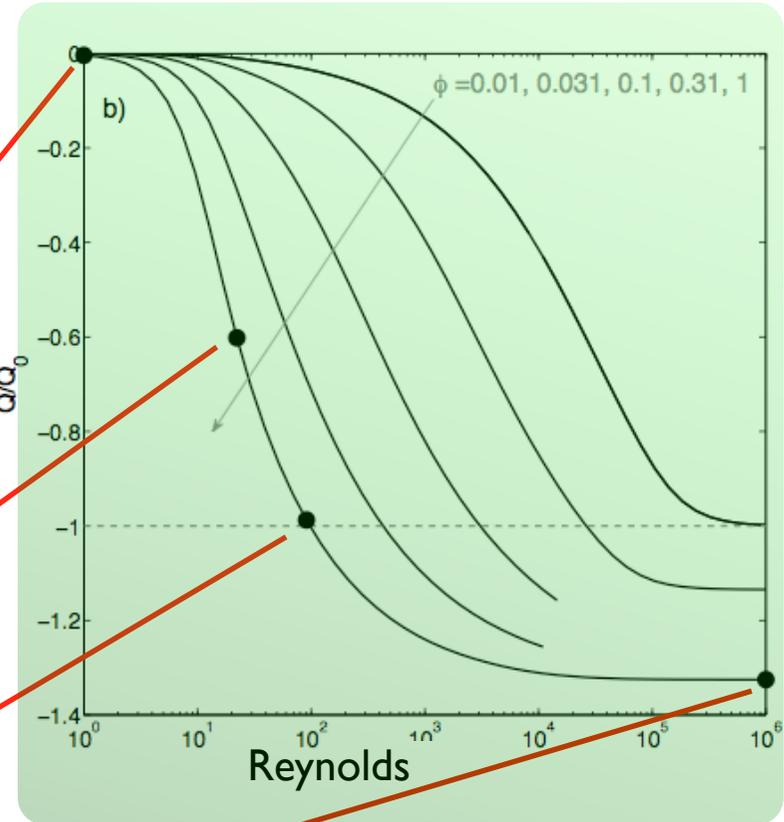
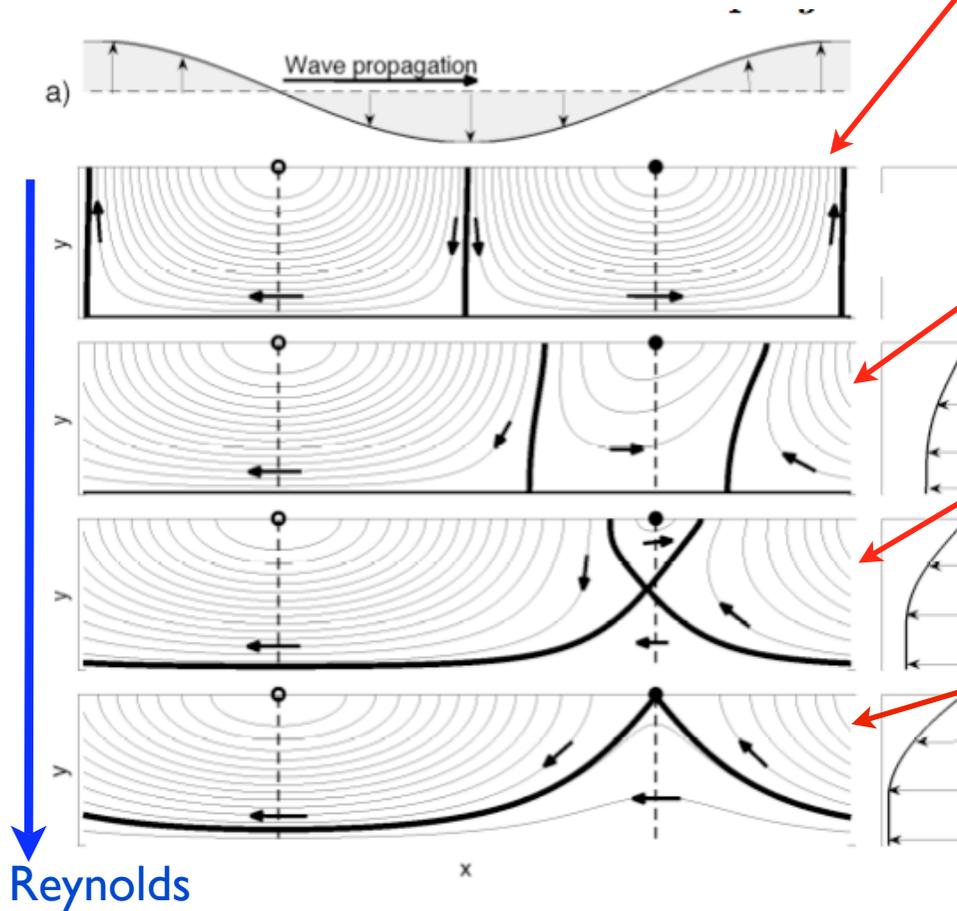
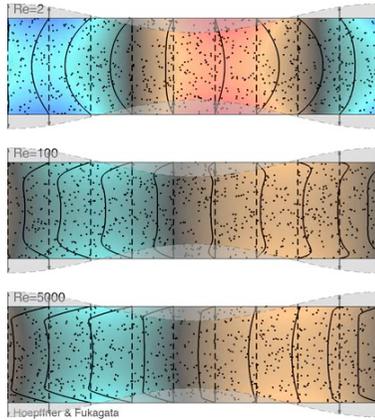
## Blowing and suction in details

Merci de votre attention!

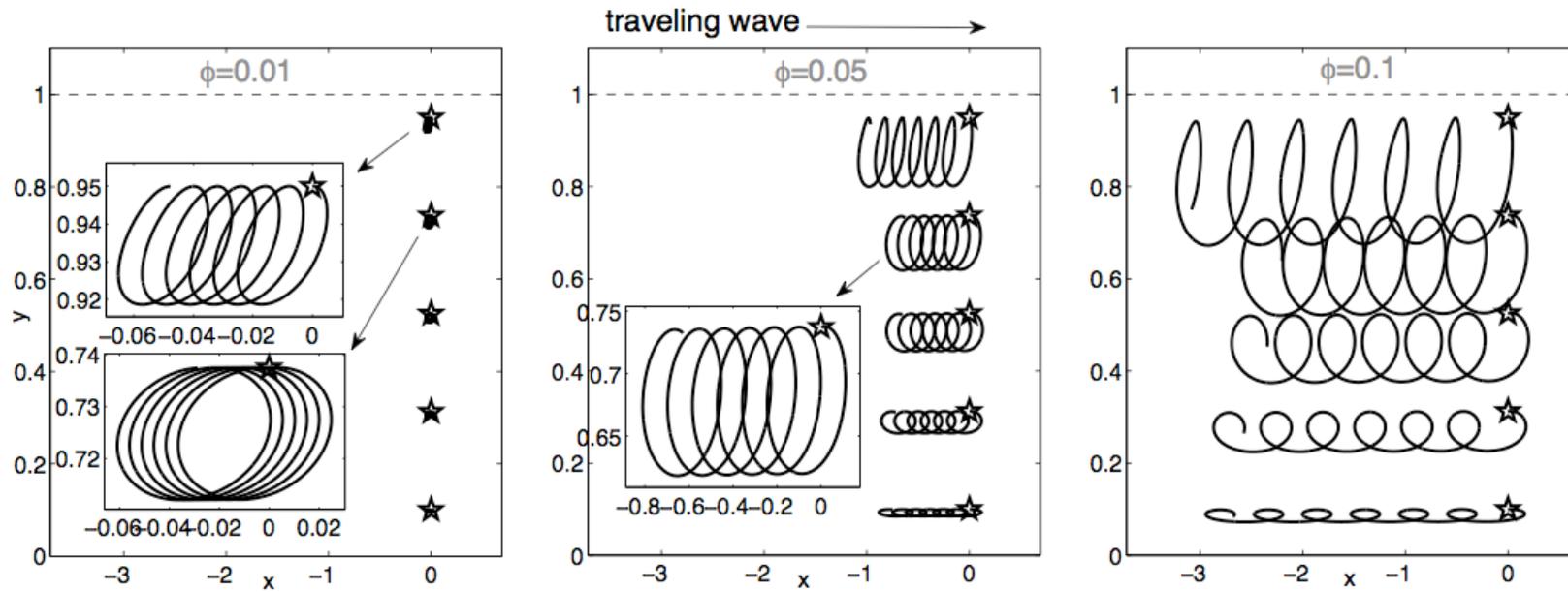


Hoepffner & Fukagata

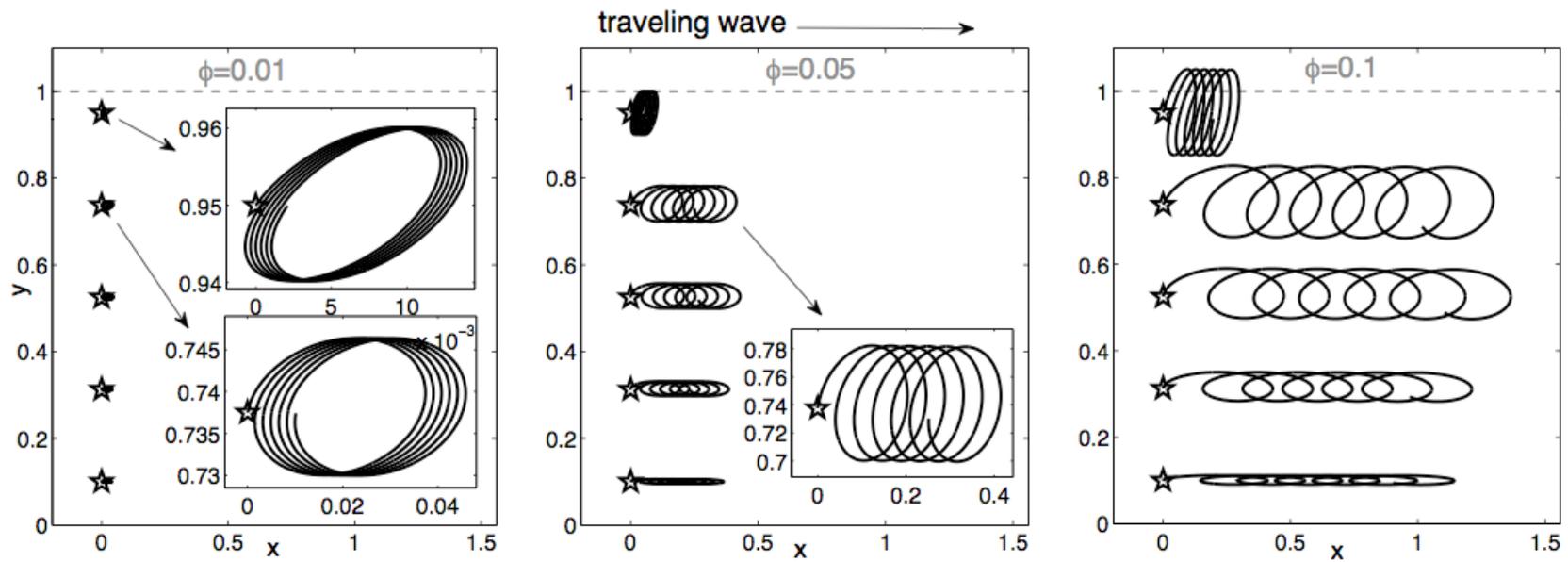
# Flow structure

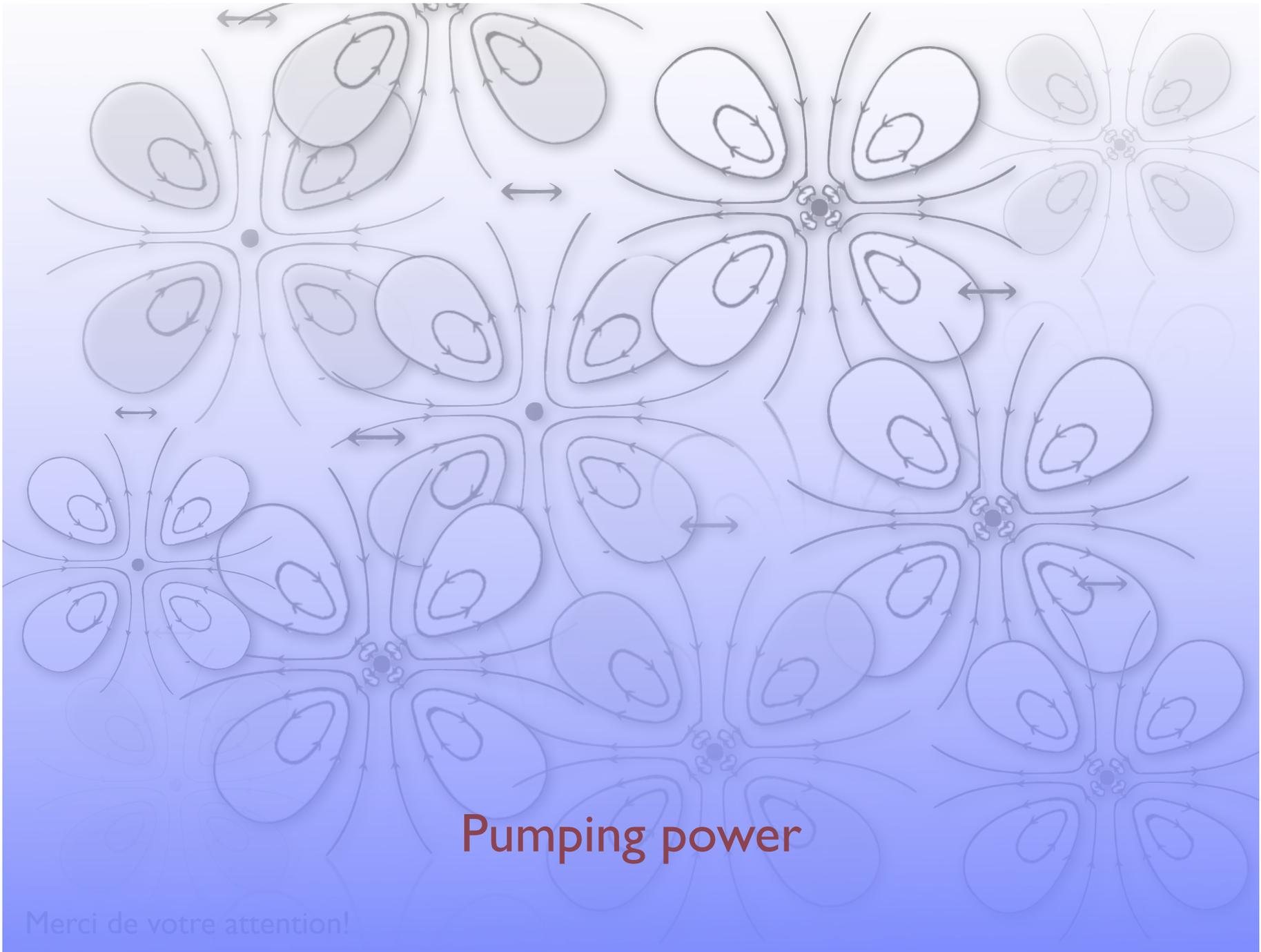


# Blowing&suction: trajectories



# Peristalsis: trajectories

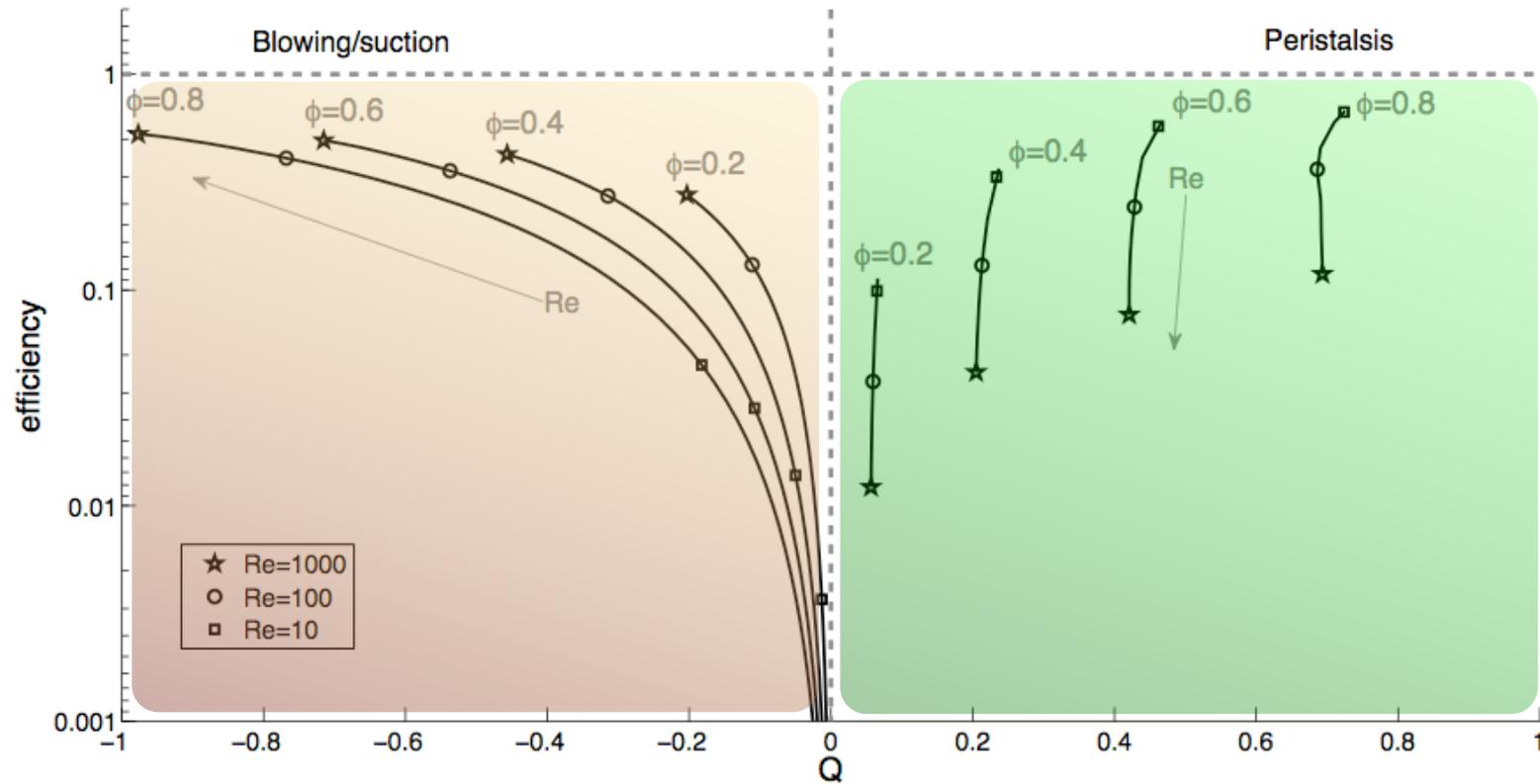




## Pumping power

Merci de votre attention!

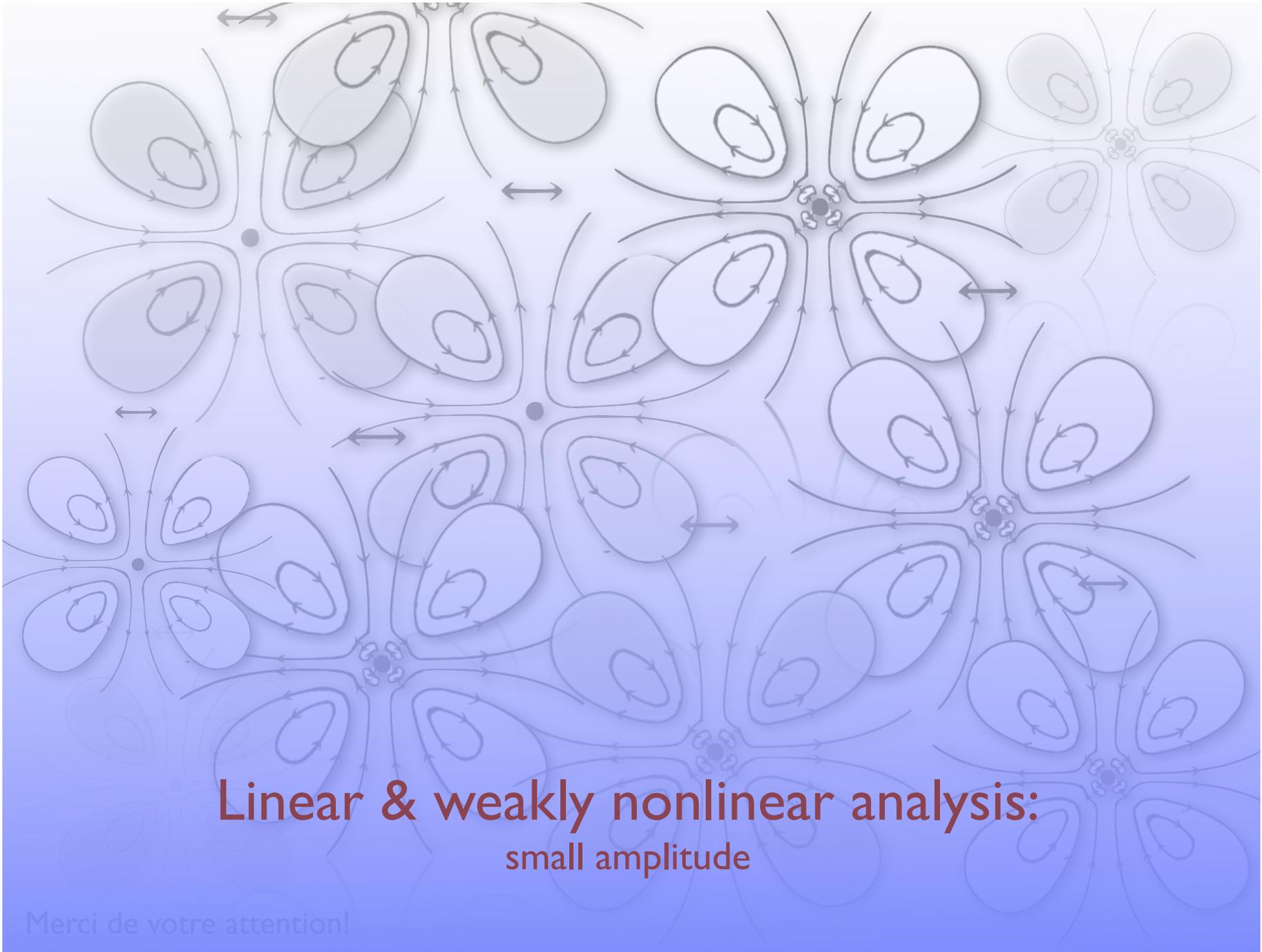
# Energy



Pumping flux  $Q$

Efficiency:

$\frac{\text{energy for this flux by Poiseuille}}{\text{energy spent here}}$



**Linear & weakly nonlinear analysis:**  
small amplitude

Merci de votre attention!

# Weakly nonlinear

## Weakly nonlinear: Fourier modes

Taylor expansion in (small) boundary condition amplitude:  $u = u^{(0)} + \epsilon u^{(1)} + \epsilon^2 u^{(2)}/2 + \dots$

Order 0: Poiseuille  
Order 1: Orr-Sommerfeld  
Order 2:  $\longrightarrow$

$$\frac{1}{Re} u_{yy}^{(2,0)} = \overline{v^{(1)} u_y^{(1)}}$$

Order 2 solution, Fourier mode 0:  
**pumping profile**

Average in time and in space

Order 0: Poiseuille - steady - due to the external pressure gradient

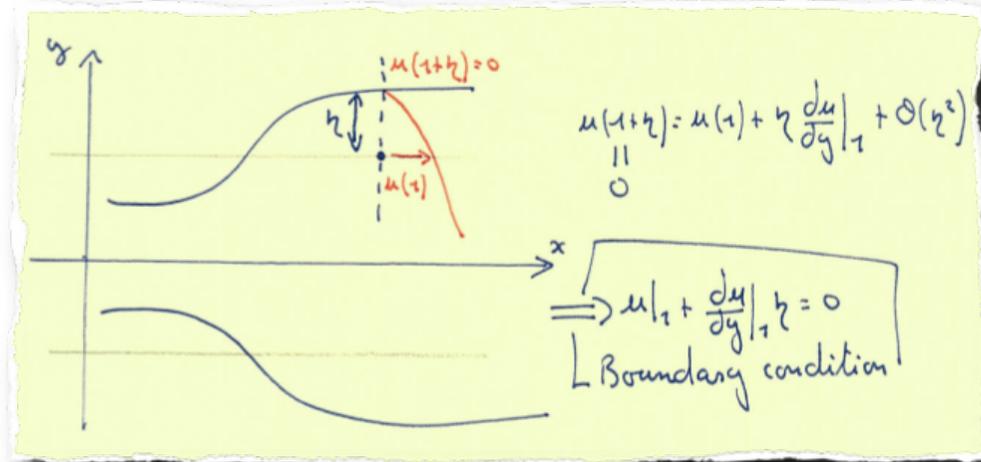
Order 1: Orr-sommerfeld, oscillatory, no mean motion

Order 2: Forced by the order one: pumping

# Boundary conditions

Peristalsis: No slip condition at the displaced wall

$$u|_{1+\varepsilon\eta} = 0 = u|_1 + \varepsilon\eta u_{y|_1} + \frac{(\varepsilon\eta)^2}{2} u_{yy|_1}$$



$$u|_{1+\eta} = 0 = [U + \varepsilon u' + \frac{\varepsilon^2}{2} u'']|_1 + \varepsilon\eta [U_y + \varepsilon u'_y + \frac{\varepsilon^2}{2} u''_y]|_1 + \frac{(\varepsilon\eta)^2}{2} [U_{yy} + \varepsilon u'_{yy} + \frac{\varepsilon^2}{2} u''_{yy}]|_1.$$

Double expansion:  
different order are mixed

Peristalsis:

$$u'|_{\pm 1} = 0, \quad v'|_{\pm 1} = \eta t,$$

Blowing/suction:

$$u'|_{\pm 1} = 0, \quad v'|_{\pm 1} = \eta,$$

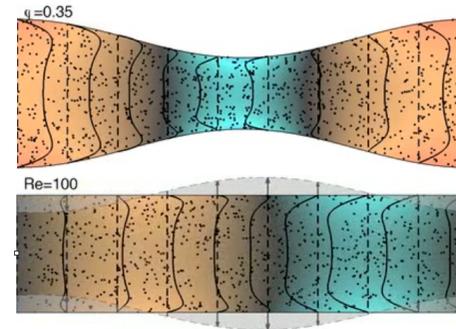
$$\overline{u''}|_{\pm 1} = \eta \overline{u'_y}|_{y=\pm 1},$$

$$\overline{u''}|_{\pm 1} = 0.$$

Slip condition for peristalsis!

# Main conclusions

- 1) oscillatory forcing can lead to mean drift - many examples
- 2) Similar flows, opposite flux



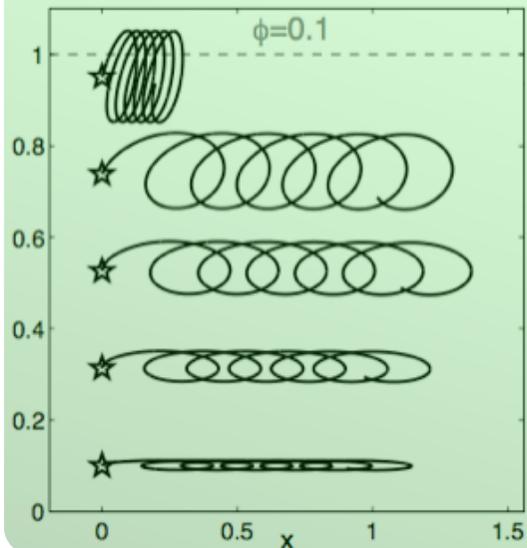
Levels of modelisation:

- **Steady state solution:**  
extract data from the equations
- **Conservation laws:**  
allowed to change the geometry to enlight the physics
- **Low amplitude perturbations**  
(asymptotics): extract structure from the equations

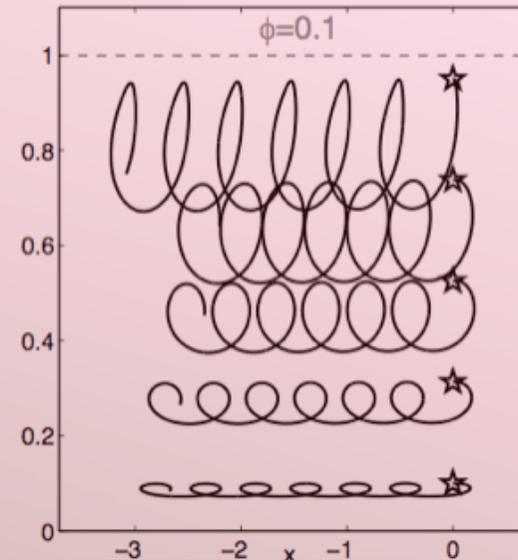
# A general feature

- 1) particles are entrained into a **circular motion** by the wall actuation
- 2) the pumping direction originates from a different viscous damping during the **backward and forward** motion of fluid particles along this circular trajectories.

For **peristalsis**, the particles' backward motion takes place in the constricted section of the channel, where viscosity slows down the flow.



For **blowing and suction**, the particles' forward motion takes place close to the walls, where viscosity slows down the flow.

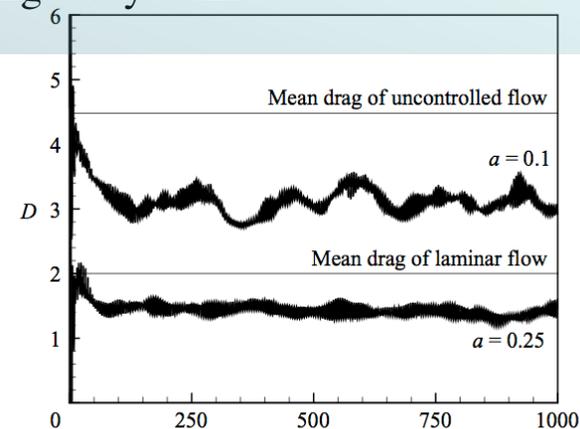


## Sustained sub-laminar drag in a fully developed channel flow

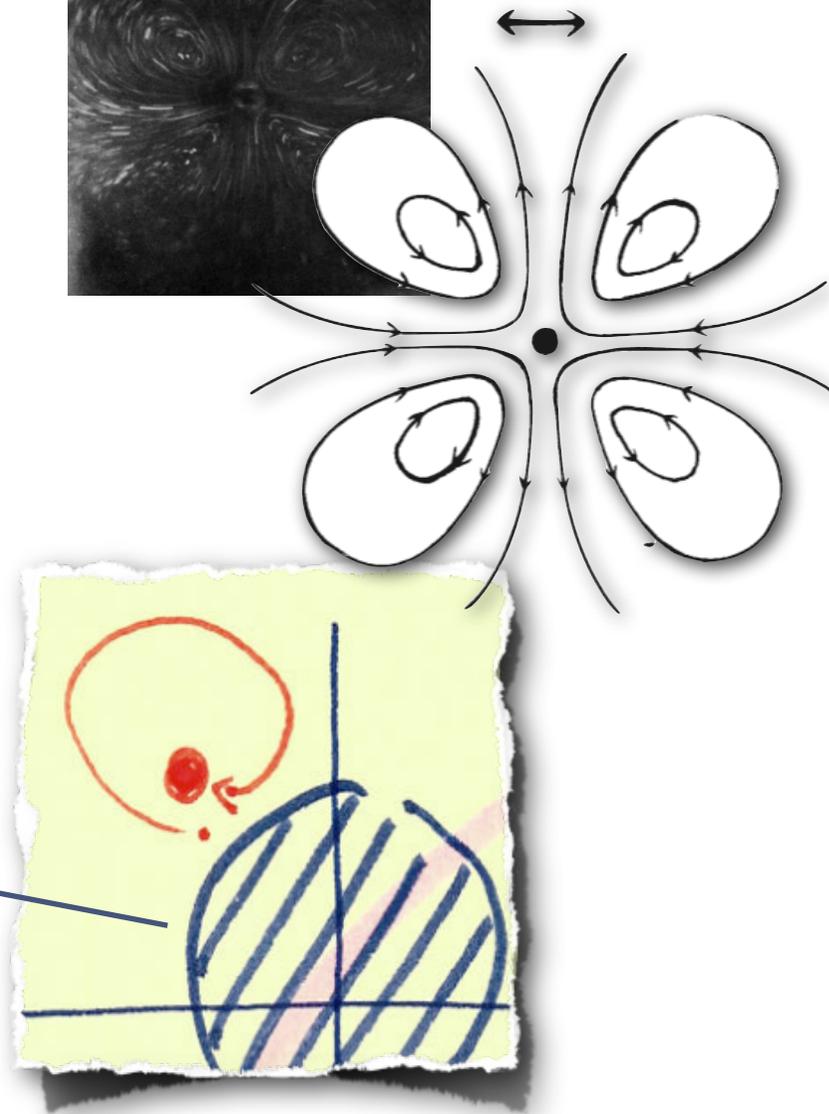
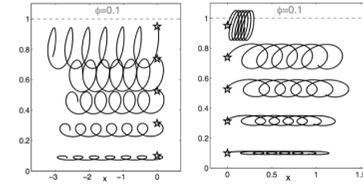
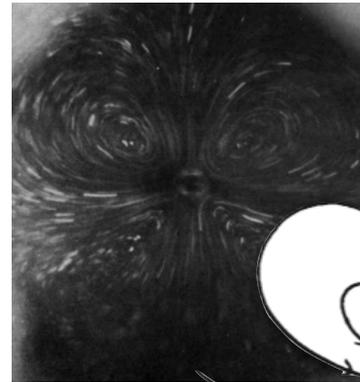
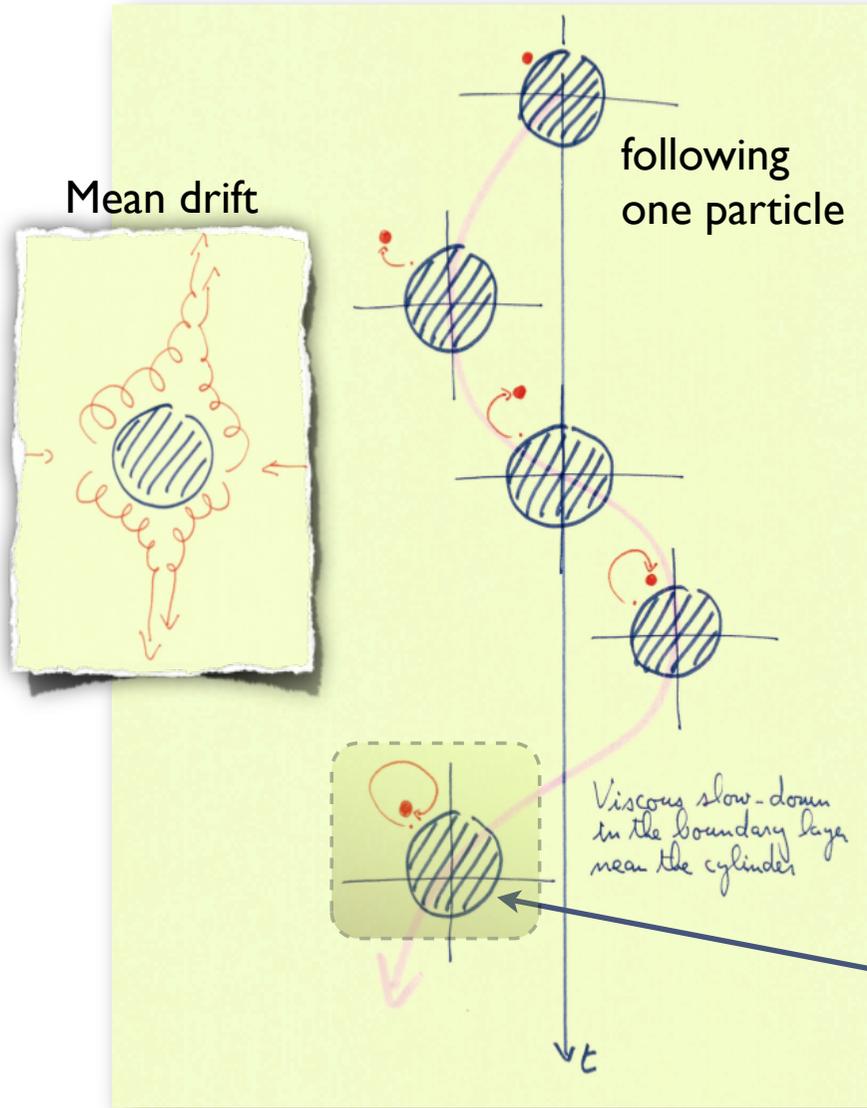
By TAE GEE MIN, SUNG MOON KANG, JASON L. SPEYER  
AND JOHN KIM

Finally, the current control scheme, consisting of surface blowing and suction in the form of travelling waves, is mathematically simple [...], yet it may not be straightforward to implement in real flows. [...]

However, a moving surface with wavy motion would produce a similar effect, since wavy walls with small amplitudes can be approximated by surface blowing and suction. We plan to perform simulations over moving wavy walls.

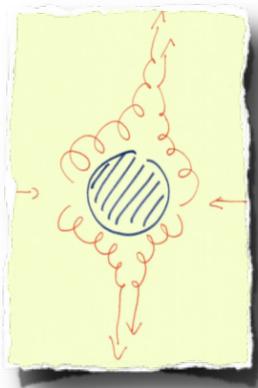
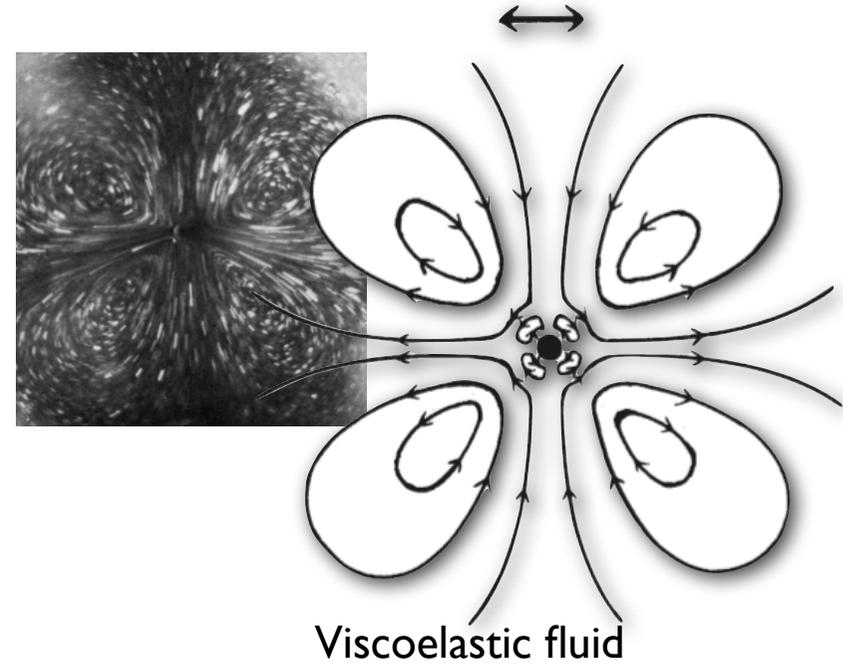
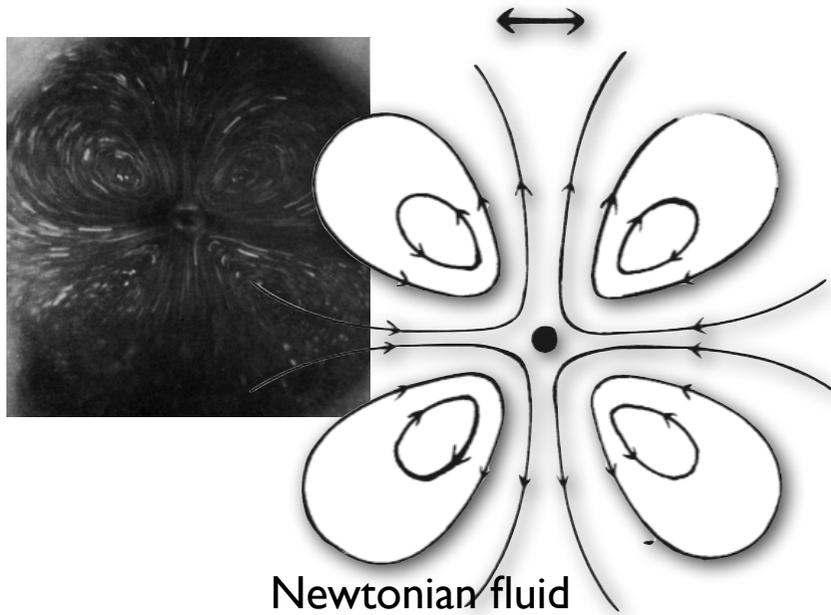


# Oscillating cylinder

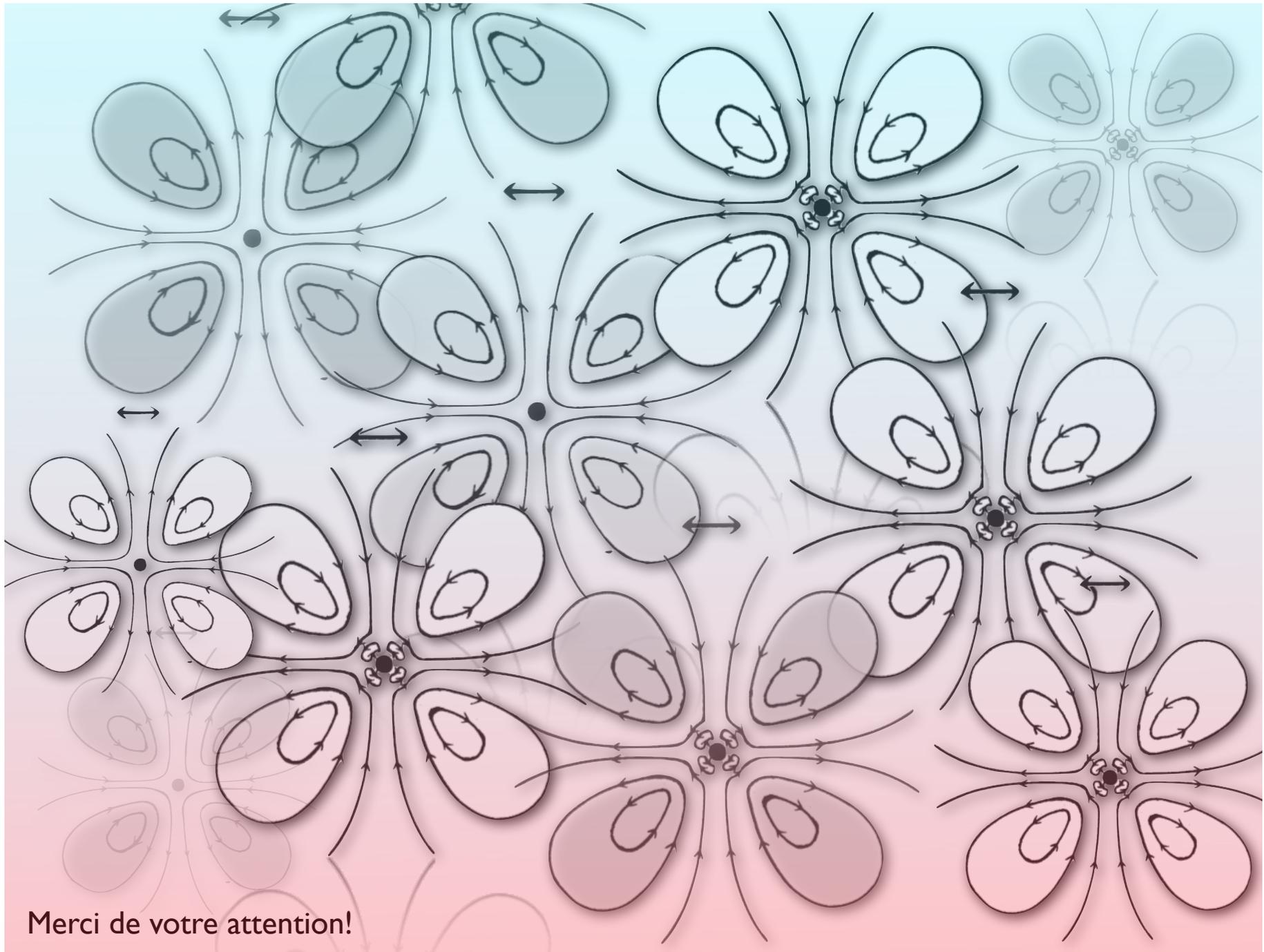


# Oscillating cylinder

Flow near an oscillating cylinder  
in dilute viscoelastic fluid



**Viscoelastic:**  
flux in direction opposed as to Newtonian



Merci de votre attention!